

Hints: Paper by Peters (1964)

Physical Review, vol. 136, page 1224.

In short, using the weak-field approximation for general relativity and after pages of derivation, we have the following equations that govern the time evolution of semi-major axis 'a' and the eccentricity 'e' for the orbit of two point masses emitting gravitational waves:

$$\left\{ \begin{array}{l} \frac{da}{dt} = -\frac{64}{5} \cdot \frac{G^3 M_1 M_2 (M_1 + M_2)}{c^5 a^3 (1-e^2)^{7/2}} \left(1 + \frac{73}{24} e^2 + \frac{37}{96} e^4 \right) \quad \text{--- (1)} \\ \frac{de}{dt} = -\frac{304}{15} \cdot \frac{G^3 M_1 M_2 (M_1 + M_2)}{c^5 a^4 (1-e^2)^{5/2}} \cdot e \cdot \left(1 + \frac{121}{304} e^2 \right) \quad \text{--- (2)} \end{array} \right.$$

$$\frac{\textcircled{1}}{\textcircled{2}} : \quad \frac{da}{de} = \frac{12}{19} \cdot \frac{a}{e} \cdot \frac{\left(1 + \frac{73}{24} e^2 + \frac{37}{96} e^4 \right)}{\left(1 - e^2 \right) \left(1 + \frac{121}{304} e^2 \right)} \quad \text{--- (3)}$$

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(3) can be integrated to find $a(e)$,

$$a(e) = \frac{C_0 e^{12/19}}{(1-e^2)} \left(1 + \frac{121}{304} e^2\right)^{870/2299}$$

where $C_0 = \frac{a_0 (1-e_0^2)}{e_0^{12/19}} \cdot \left(1 + \frac{121}{304} e_0^2\right)^{-870/2299}$ is a constant fixed by the initial a_0 and e_0 .

By setting $\beta = \frac{64}{5} \cdot \frac{G^3 M_1 M_2 (M_1 + M_2)}{c^5}$, we can rewrite (2),

$$\frac{de}{dt} = -\frac{19}{12} \cdot \frac{\beta}{C_0^4} \cdot \frac{e^{-29/19} (1-e^2)^{3/2}}{\left[1 + \frac{121}{304} e^2\right]^{1181/2299}} \quad \text{--- (4)}$$

Since $e \rightarrow 0$ as $a \rightarrow 0$, we can find the inspiral time t_{insp} by integrating (4)

$$t_{\text{insp}} = \frac{12}{19} \cdot \frac{C_0^4}{\beta} \cdot \int_0^{e_0} \frac{e^{29/19} \left(1 + \frac{121}{304} e^2\right)^{1181/2299}}{(1-e^2)^{3/2}} de,$$

which can be evaluated numerically.

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Assuming that the black hole (BH) binary system was formed at redshift $z_0 = 20.0$;

LIGO observation found that the merger occurred at $z_{\text{merger}} = 0.088 \sim 0.1$.

By picking a reasonable initial eccentricity e_0 , we can solve for a_0 numerically:

$$\underbrace{t(z_{\text{merger}})}_{\text{known}} - \underbrace{t(z_0)}_{\text{assumed}} = t_{\text{imp}}(\underbrace{a_0}_{\text{find.}}, \underbrace{e_0}_{\text{picked}})$$

Note: If we set $e_0 = 0$, we can write

$$t_{\text{imp}} = a_0^4 / 4\beta.$$

If we consider e_0 to be small,

$$t_{\text{imp}} \simeq \frac{12}{19} \cdot \frac{c_0^4}{\beta} \int_0^{e_0} e^{29/19} de = \frac{c_0^4}{43} \cdot e_0^{48/19}$$

- Both can be solved without numerical integration.