

Pulsars and Magnetars (cont'd)

March 20, 2008

• Rotational Energy

$$E_{\text{rot}} = \frac{1}{2} I \omega^2, \quad \omega \hat{=} \text{angular velocity}$$
$$= \frac{2\pi}{P}$$

$P \leftarrow \text{period}$

$$I \hat{=} \text{moment of inertia} = \int_0^R r^2 dm$$

for solid sphere:



$$I = \frac{2}{5} MR^2$$

→ for typical NS parameters

$$\left(\begin{array}{l} R \approx 10 \text{ km} \\ M \approx 1 m_{\odot} \end{array} \right) \Rightarrow I = 10^{38} \text{ kg m}^2$$

→ know from observations that pulsars slow down!

$$\text{- e.g. Crab Pulsar: } \omega = \frac{2\pi}{P} = \frac{2\pi}{33 \text{ ms}}$$
$$= 190 \text{ s}^{-1}$$

observed: $\frac{dP}{dt} = \frac{1 \text{ ms}}{90 \text{ yr}}$

$$\frac{d\omega}{dt} = \frac{-2\pi}{P^2} \frac{dP}{dt} = -2 \times 10^{-9} \text{ s}^{-2}$$

$$\frac{dE_{\text{rot}}}{dt} = I\omega \frac{d\omega}{dt} = -4 \times 10^{31} \text{ W}$$

Estimate pulsar age τ :

$$P \approx 2 \frac{dP}{dt} \tau \Rightarrow \boxed{\tau \equiv \frac{P}{2\dot{P}}}, \quad \dot{P} = \frac{dP}{dt}$$

("spin-down age")

- $\tau \sim 1,500 \text{ yr}$, for the Crab pulsar.
(true age $\sim 950 \text{ yr}$)
- Compare with radiation output from Crab. $\rightarrow L_{\text{Crab}} = \frac{dE}{dt} = 5 \times 10^{31} \text{ W}$.

Big Q: How is rotation converted into radiation?

• Radiation mechanism

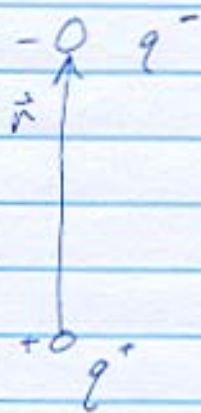
- in general: accelerated charges radiate

(3)

$$\Rightarrow \frac{dE}{dt} = \frac{1}{6\pi\epsilon_0 c^3} q^2 \underbrace{|\vec{a}|^2}_{\text{acceleration}}, \text{ Larmor's formula.}$$

$$(\epsilon_0 = 8.85 \times 10^{-12} \text{ C}^2 \text{ N}^{-1} \text{ m}^{-2})$$

Consider an electric dipole

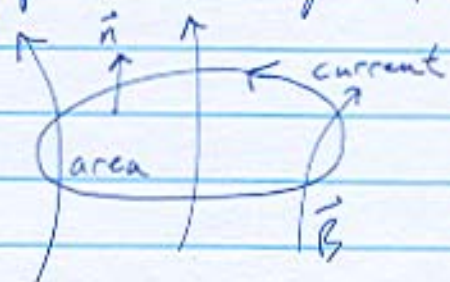


electric dipole moment $\vec{d} = q \vec{r}$

$$\boxed{\frac{dE}{dt} = \frac{1}{6\pi\epsilon_0 c^3} |\ddot{\vec{d}}|^2}$$

→ electric dipole radiation

→ magnetic analogue, "magnetic dipole moment"



$$\begin{aligned} \vec{\mu} &= \text{area} \times \text{current} \hat{n} \\ &= \pi R^2 \times i \hat{n} \end{aligned}$$

from Ampere's law:

$$\oint \vec{B} \cdot d\vec{l} = \mu_0 i \quad \left(\mu_0 = 4\pi \times 10^{-7} \text{ N A}^{-2} \right)$$

$$\Rightarrow 2\pi R B = \mu_0 i$$

$$\Rightarrow \mu \equiv |\vec{\mu}| = \frac{1}{\mu_0} R^3 B$$

→ magnetic dipole radiation:

$$\frac{dE}{dt} = \frac{1}{6\pi\epsilon_0 c^3} \frac{\mu_0}{4\pi} |\ddot{\mu}|^2$$

- Consider sinusoidal variation

$$\mu = A \sin \omega t$$

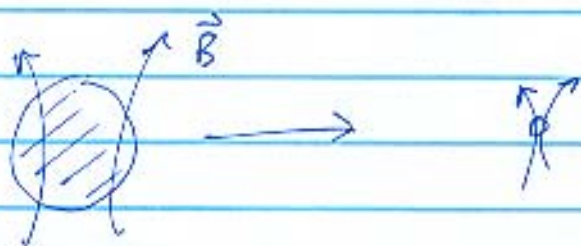
$$\ddot{\mu} = -\omega^2 \mu$$

$$\frac{dE}{dt} = \frac{2}{3c^3} \frac{\mu_0}{4\pi} \omega^4 \frac{R^6 B^2}{\mu_0^2}$$

For Crab, we find that $B = 10^9 \text{ T}$

Q: Where do such strong B-fields come from?

A: Collapse of star!



- magnetic fields are "frozen" into fluid
→ MHD (magnetohydrodynamics)

→ magnetic flux $= \phi = \text{area} \times B$
 $= \pi r^2 B = \text{constant}$

$$B_f = B_i \left(\frac{r_i}{r_f} \right)^2$$

$$r_i = 10^6 \text{ km}$$

$$r_f = 10 \text{ km}$$

$$B_i = 10^{-2} \text{ T}$$

$$\Rightarrow B_f = 10^8 \text{ T}$$

• Magnetars

- Boost of $\vec{B}$ -field by another factor of 1,000 by dynamo action. $\Rightarrow \vec{B} = 10^{11} \text{ T}$

Note: This is larger than "quantum-critical field," where QED effects kick in!

Consider: motion of electron in $\vec{B}$ -field



$$e v B = \frac{m_e v^2}{r_L}$$

$$\Rightarrow r_L = \frac{m_e v}{e B} \quad \text{"Larmor radius"}$$

If $\vec{B}$ -field is very strong, we can have

$$r_L = \lambda_{dB} = \frac{h}{p}$$

Consider NR/UR boundary: $v \sim c$, $p \sim m_e c$

$$\frac{m_e c}{e B_{QED}} = \frac{h}{m_e c}$$

$$\Rightarrow B_{QED} = \frac{m_e^2 c^2}{e \hbar}$$

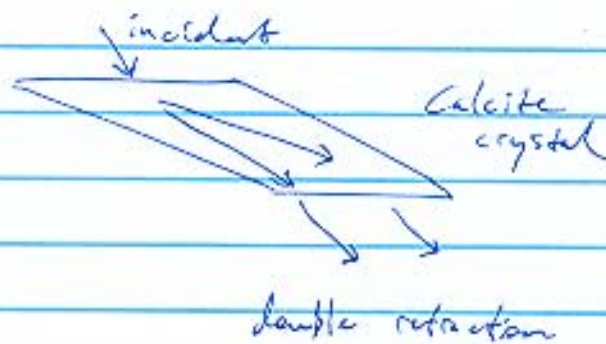
$$\approx 4 \times 10^9 \text{ T}$$

Thus: $B_{magnetar} > B_{QED}$.

→ Weird effects

(1) Vacuum birefringence

(2) Photon splitting



→ Distort shape of atoms
(e.g. Fe in NS crust)

