

- 1) a) The free-fall timescale implies the density of a star, and also characterizes the timescale on which dynamical changes occur, such as pulsations and gravitational collapse.
- b) The maximum mass that a white dwarf may have - the Chandrasekhar Mass.
- c) Pressure is a consequence of energy (units are of energy density), and in GR both mass and energy are sources of gravity.
- d) The pressure exerted by electrons packed tightly together enough that the Pauli Exclusion Principle plays an important role in determining their bulk properties - degeneracy pressure.
- e) The Hertzsprung - Russell Diagram is a plot displaying the relation of temperature to luminosity for stars.

f) Photons are redshifted as they lose kinetic energy in climbing out of a gravitational potential well - gravitational redshift. (2)

g) i) White dwarf - $R \sim 10,000 \text{ km}$
ii) Neutron star - $R \sim 10 \text{ km}$

h) The Fermi momentum is the characteristic momentum of fermions packed together maximally in phase-space, as dictated by the Pauli Exclusion Principle and the Heisenberg Uncertainty Principle.

i) Spacetime is a 4-dimensional construct unifying the three spatial dimensions and time.

j) A black hole with mass equal to the Planck Mass has a position that is not defined to be within its own Schwarzschild radius, while GR demands that nothing interior to this radius can ever be found outside of it.

2 a) $\tau_{ff} = \sqrt{\frac{1}{G\rho}}$
 $= \sqrt{\frac{1}{G n m_H}}$
 $\sim 10^8 \text{ yr.}$

b) Minkowski said this!

c) The proper time for a photon never advances, in any frame - $d\tau = \frac{ds}{c} = 0$ always. Thus, the clock registers no passage of time.

3 a) Integration yields

$$\int_0^{P(r)} dP = \int_R^r -\frac{4\pi}{3} G \rho_c^2 r \exp\left(-\frac{r^2}{a^2}\right)$$

$$\Rightarrow P(r) = \frac{2\pi}{3} G a^2 \rho_c^2 \exp\left(-\frac{r^2}{a^2}\right) \Bigg|_R^r$$

$$= \frac{2\pi}{3} G a^2 \rho_c^2 \left[\exp\left(-\frac{r^2}{a^2}\right) - \exp\left(-\frac{R^2}{a^2}\right) \right]$$

For $P_c = P(0)$, we find

$$P_c = \frac{2\pi}{3} G a^2 \rho_c^2 \left[1 - \exp\left(\frac{-R^2}{a^2}\right) \right],$$

which for $R = 2 R_\odot$ gives $P_c \sim 10^{18}$ Pa.

b) As $r \rightarrow 0$, the mass enclosed $M(r) \rightarrow 0$, which implies that the gravitational force balanced by the pressure gradient goes to zero, as well. Thus, the pressure gradient must disappear, too.

[4] a) Just roughly, we have

$$E_{\text{pot}} \sim \frac{GM^2}{R}$$

$$\sim \frac{7 \times 10^{-11} \times 4 \times 10^{60}}{7 \times 10^6}$$

$$\sim 4 \times 10^{43} \text{ erg}$$

b) Similarly, $P_c \sim \frac{GM^2}{R^4} \sim 10^{23}$ Pa.

(5)

c) Once again very roughly, we balance the pressure from part b with the ideal gas pressure, $P = nk_B T$. We estimate that

$$n \sim \frac{M}{m_H R^3}, \text{ so that we obtain}$$

$$T = \frac{GMm_H}{k_B R}$$

$$\sim \frac{10^{-10} 10^{30} 10^{-27}}{10^{-23} 10^7} \sim 10^9 \text{ K}$$

[5] a) With respect to the coordinate time, t , measured at $r \rightarrow \infty$, a clock at radius r will read

$$\tau = \left(1 - \frac{2GM}{rc^2}\right) t.$$

Therefore, the ratio of the proper times at two radii, r_{out} and r_{in} , will be

$$\frac{\tau_{out}}{\tau_{in}} = \frac{\left(1 - \frac{2GM}{r_{out}c^2}\right)}{\left(1 - \frac{2GM}{r_{in}c^2}\right)} \quad (6)$$

In the case here $r_{in} = 5 R_{\odot}$, $r_{out} = 10'' m$,
and $\tau_{in} = 1 \text{ sec}$. Thus, here

$$\tau_{out} = 1.0001 \text{ sec}.$$

b) Rewriting the metric, we have

$$ds^2 = -c^2 d\tau^2, \text{ which yields}$$

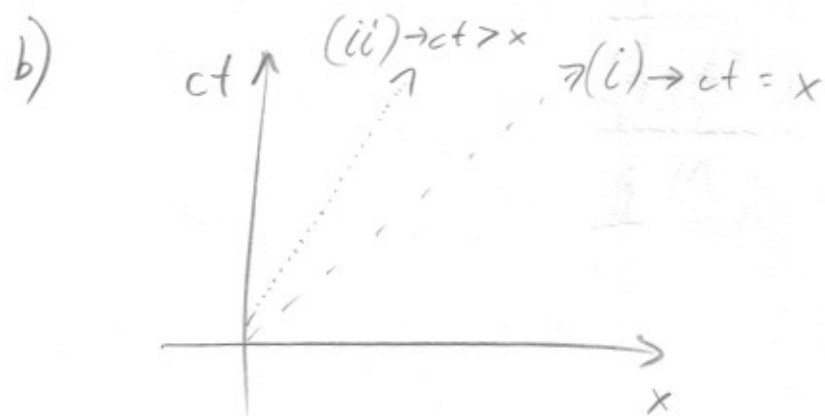
$$d\tau^2 = \left(1 - \frac{2GM}{rc^2}\right) dt^2 - \frac{1}{c^2} \left(1 + \frac{2GM}{rc^2}\right) (dx^2 + dy^2 + dz^2).$$

If $c \rightarrow \infty$, then we obtain

$$d\tau^2 = dt^2 \Rightarrow d\tau = dt, \text{ which is}$$

intended to mean that there is a universal
time measured to be the same by all observers,
independent of r , x , y , or z .

$$\boxed{6} \text{ a) } ds^2 = -c^2 dt^2 + dx^2 + dy^2 + dz^2 \quad (7)$$



c) An observer will see the photon traveling, say, in the x -direction with speed c , such that $dx = c dt$ (with $dy = dz = 0$). Then,

$$\begin{aligned} ds^2 &= -c^2 dt^2 + dx^2 \\ &= -c^2 dt^2 + c^2 dt^2 \\ &= 0 \end{aligned}$$