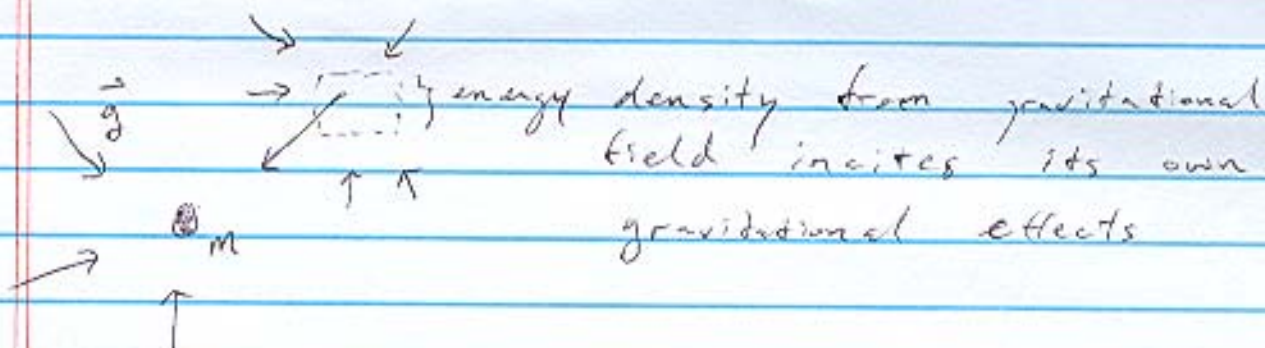


# General Relativity (final)

3/20/07


- Very complicated!  $\rightarrow$  Non-linear!



## Hydrostatic equilibrium in GR

- generalize Newtonian expression  $\frac{dP}{dr} = -\frac{Gm}{r^2} \rho$

(1)  $\rho \hat{=}$  "measure of inertia"  $\rightarrow$  inertial mass density

- in relativity  } total energy density,  $\rho c^2 + P$

$\rightarrow$  Replace  $\rho \rightarrow \rho + \frac{P}{c^2}$

(2)  $m \hat{=}$  "source of gravity"

$\rightarrow$  Replace  $m \rightarrow m_{\text{eff}} = \frac{4\pi}{3} r^3 \rho_{\text{eff}} = \frac{4\pi}{3} r^3 \rho + 4\pi r^3 \frac{P}{c^2}$

$$\Rightarrow m_{\text{eff}} = m + \frac{4\pi r^3 \rho}{c^2}$$

$$= m \left( 1 + \frac{4\pi r^3 \rho}{m c^2} \right)$$

(3) Space is curved in a gravitational field

$$\rightarrow \text{replace } r^2 \rightarrow r^2 \left( 1 - \frac{2GM}{c^2 r} \right)$$

Now,

$$\frac{dP}{dr} = -\rho \left( 1 + \frac{P}{\rho c^2} \right) \frac{GM \left( 1 + \frac{4\pi r^3 \rho}{m c^2} \right)}{r^2 \left( 1 - \frac{2GM}{r c^2} \right)}$$

$$= -\rho \frac{GM}{r^2} \left( 1 + \frac{P}{\rho c^2} \right) \left( 1 + \frac{4\pi r^3 \rho}{m c^2} \right) \left( 1 - \frac{2GM}{r c^2} \right)^{-1}$$

"Oppenheimer-Volkoff (OV) equation"

Notice: • You can recover the Newtonian form of this equation if  $c \rightarrow \infty$   
 $\rightarrow$  Newton's theory is "action at a distance"

• You can recover Newton's form also for a weak gravitational field:

$$\frac{GM}{c^2 r} \ll 1 \quad \text{and} \quad P \ll \rho c^2$$

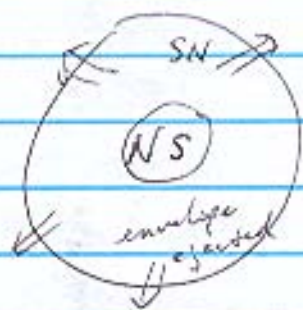
⇒  $OV$  equation describes structure of neutron stars!

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## Neutron Stars

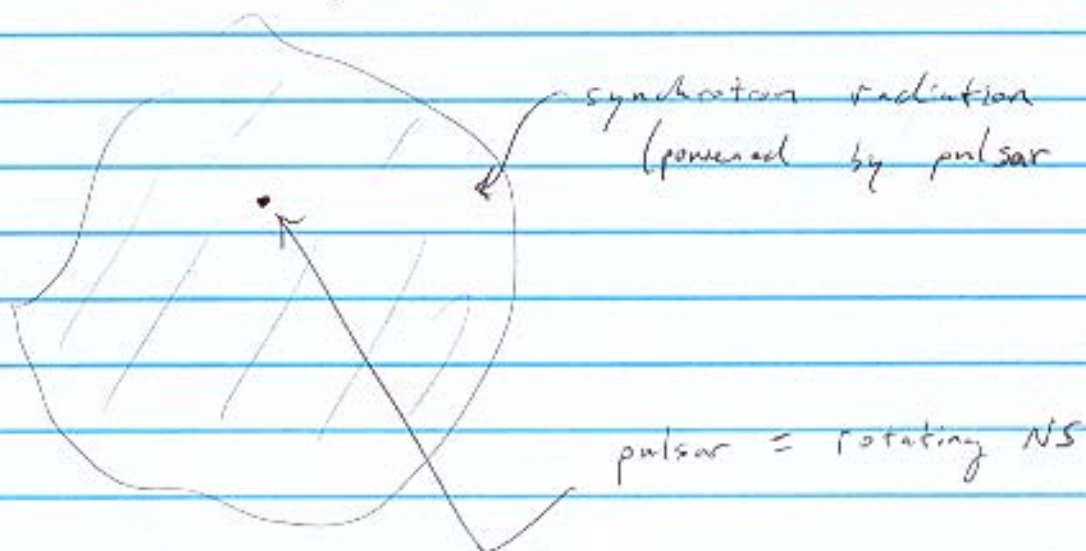
- Discovery: first postulated by theory (Bartlett & Zwicky) → 1934 -
- Idea (now known to be correct): (London, 1938)

"Neutron stars (NS) are born when massive stars die in a supernova (SN) explosion  
→ Initial mass of star  $\geq 8 M_{\odot}$



- By serendipity, radio astronomers observe pulsars (Bell & Hewish 1967)

e.g. in Crab nebula



→ Can not be WDs!

→ Rotation would have torn a WD apart

→ Oscillations too fast for WD pulsational period:  
 $\sim \tau_{\text{fs WD}} \sim 1 \text{ sec}$

### Basic Properties

→ pulsar observations

-  $M_{\text{NS}} \sim 2 M_{\odot}$  (from binary pulsars)

- Estimate radius from observed pulsation period:

$$P_{\text{NS}} = 1/30 \text{ sec}$$



centrifugal acceleration  $<$  gravity

rotation speed at surface:  $V_{\text{rot}} \approx \frac{2\pi R_{\text{NS}}}{P_{\text{NS}}}$

$$\frac{V_{\text{rot}}^2}{R_{\text{NS}}} < \frac{GM_{\text{NS}}}{R_{\text{NS}}^2} \Rightarrow R_{\text{NS}} < \left( \frac{GM_{\text{NS}} P_{\text{NS}}}{4\pi^2} \right)^{1/3} \approx 100 \text{ km}$$

Now known that  $R_{\text{NS}} \sim 10 \text{ km}$

$$\text{Average density } \langle \rho \rangle_{NS} = \frac{M_{NS}}{\left(\frac{4\pi R_{NS}^3}{3}\right)} \sim 10^{17} \text{ kg m}^{-3}$$

For comparison,  $\langle \rho \rangle_{\odot} \approx 1,000 \text{ kg m}^{-3}$

$$\langle \rho \rangle_{WD} = 10^9 \text{ kg m}^{-3}$$

An extreme density!

$$\rightarrow \text{Nuclear density} \sim \frac{M_H}{\frac{4\pi}{3} r_0^3}$$

$$\rightarrow \text{Nuclear radius } R = r_0 A^{1/3}$$

low  
nucleon number (n+p)

$$r_0 = 1 \text{ Fermi} = 10^{-15} \text{ m}$$

$\rightarrow$  NS are like giant atomic nuclei, but:

- NS is held together by gravity
- Nucleus is held together by strong force