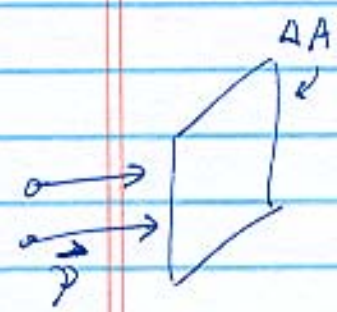


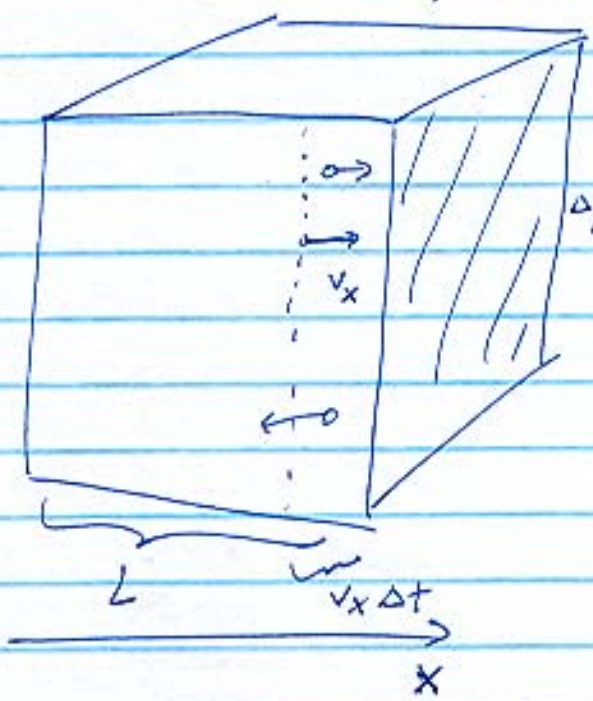
# Physics of compact objects (cont)

Pressure,  $P$ :

$$P = \frac{\Delta p}{\Delta A \Delta t} \quad (\text{"momentum flux"})$$



Consider a box, with side length  $L$ , filled with particles:



$$V = L^3$$

$$N = n L^3$$

total # of particles      number density of particles

$$\Delta V = \Delta A v_x \Delta t$$

$$\Delta N = \frac{1}{2} n \Delta V$$

Momentum in  $x$ -direction:

$$p_x = m v_x$$

where  $m$  is the mass of a particle.

$$= \frac{1}{2} n \Delta A v_x \Delta t$$

only  $\frac{1}{2}$  of the particles in the volume  $\Delta V$  travel in the  $+x$  direction

$$\Delta p_x = 2 \Delta N p_x = n \Delta A p_x v_x \Delta t$$

↑ multiply by 2 because the particles bounce off the 'wall of the box' and their velocity goes toward the negative  $x$ -direction, i.e.  
 $\Delta p_x = 2 p_x$  for each particle

Pressure,  $P$ , is thus:

$$P = \frac{\Delta p_x}{\Delta A \Delta t} = n \langle p_x v_x \rangle$$

Now, because the pressure is isotropic

$$\begin{aligned} \langle \vec{p} \cdot \vec{v} \rangle &= \langle p_x v_x \rangle + \langle p_y v_y \rangle + \langle p_z v_z \rangle \\ &= 3 \langle p_x v_x \rangle \end{aligned}$$

$$\Rightarrow P = \frac{n}{3} \langle \vec{p} \cdot \vec{v} \rangle$$

Pressure in an ideal gas:

- non-relativistic (NR) particles, i.e.

$$E_{\text{kin}} \ll m_0 c^2$$



$$\left. \begin{aligned} p &= m_0 v \\ \epsilon_{kin} &= \frac{1}{2} m_0 v^2 \end{aligned} \right\} \Rightarrow p v = 2 \epsilon_{kin}$$

$$\Rightarrow P = \frac{1}{3} \frac{N}{V} \langle \vec{p} \cdot \vec{v} \rangle$$

$$= \frac{2}{3} \frac{N}{V} \epsilon_{kin} = \frac{2}{3} U_{kin}, \quad \text{where}$$

$U_{kin} \equiv$  kinetic energy density

Now, for ultra-relativistic particles (UR):

$$\epsilon_{kin} \gg m_0 c^2 \quad v = c$$

$$\Rightarrow \epsilon_{kin} \approx \underbrace{\left( m_0^2 c^4 + p^2 c^2 \right)^{1/2}}_{\epsilon_{tot}} \approx pc$$

$$\Rightarrow p v = \epsilon_{kin}, \quad \text{and}$$

$$P = \frac{1}{3} \frac{N}{V} \langle \vec{p} \cdot \vec{v} \rangle = \frac{1}{3} \frac{N}{V} \epsilon_{kin} = \frac{1}{3} U_{kin}$$

Total Energy and Stability:

Consider NR gas:

For stability:  $\langle P \rangle = -\frac{1}{3} \frac{E_{\text{pot}}}{V}$

We have that  $\langle P \rangle = \frac{2}{3} \langle U_{\text{kin}} \rangle$   
 $= \frac{2}{3} \frac{E_{\text{kin}}}{V}$ ,

where  $E_{\text{kin}} \hat{=}$  total kinetic energy.

In hydrostatic equilibrium:

$$\frac{2}{3} \frac{E_{\text{kin}}}{V} = -\frac{1}{3} \frac{E_{\text{pot}}}{V}$$

$$\Rightarrow \boxed{2 E_{\text{kin}} = -E_{\text{pot}}}$$

often called the virial theorem. (VT)

Total energy,  $E_{\text{tot}} = E_{\text{kin}} + E_{\text{pot}}$

If system is gravitationally bound, then

$E_{\text{tot}} < 0$ , since one must have no velocity outwards,  $v$ , as  $R \rightarrow \infty$ :

$$E_{\text{tot}} = \frac{1}{2} m v^2 - \frac{G m^2}{R}$$

Now,  $\boxed{E_{\text{tot}} = -2 E_{\text{kin}} + E_{\text{kin}} = -E_{\text{kin}}}$ ,



from the virial theorem.

$\Rightarrow$  NR gas can be stable.

Consider UR gas:

For stability:  $\langle P \rangle = -\frac{1}{3} \frac{E_{pot}}{V}$

We have that  $\langle P \rangle = \frac{1}{3} \frac{E_{kin}}{V}$

$\Rightarrow$   $E_{tot} = E_{pot} + E_{kin} = 0$