

Jupiter vs. ISM chemistry

how same + different? why different?

CO vs CH<sub>4</sub>

C<sub>4</sub>H<sub>2</sub> vs C<sub>3</sub>H<sub>8</sub>

Dust heating + cooling example - protoplanetary disk

black puzzle:

$$f_{abs} = \pi a^2 F_* = \pi a^2 L_* / 4\pi d^2 = \pi a^2 \sigma T_*^4 R_*^2 / d^2$$

$$\langle f_{em} \rangle = 4\pi a^2 \sigma T^4$$

$$\text{equil: } T^4 = \frac{R_*^2}{4d^2} T_*^4 \Rightarrow T \propto d^{-0.5}$$

Dust + grain:

$$f_{abs} = \pi a^2 \langle \langle f_{abs} \rangle \rangle F_* = \pi a^2 \int \langle f_{abs}(v) \rangle F_v dv$$

$$\langle f_{em} \rangle = 4\pi a^2 \langle \langle f_{em} \rangle \rangle \tau T^4 = 4\pi a^2 \int \langle f_{em}(v) \rangle \pi b_v(\tau) dv$$
$$\langle \langle f_{em} \rangle \rangle = g_{em}(v)$$

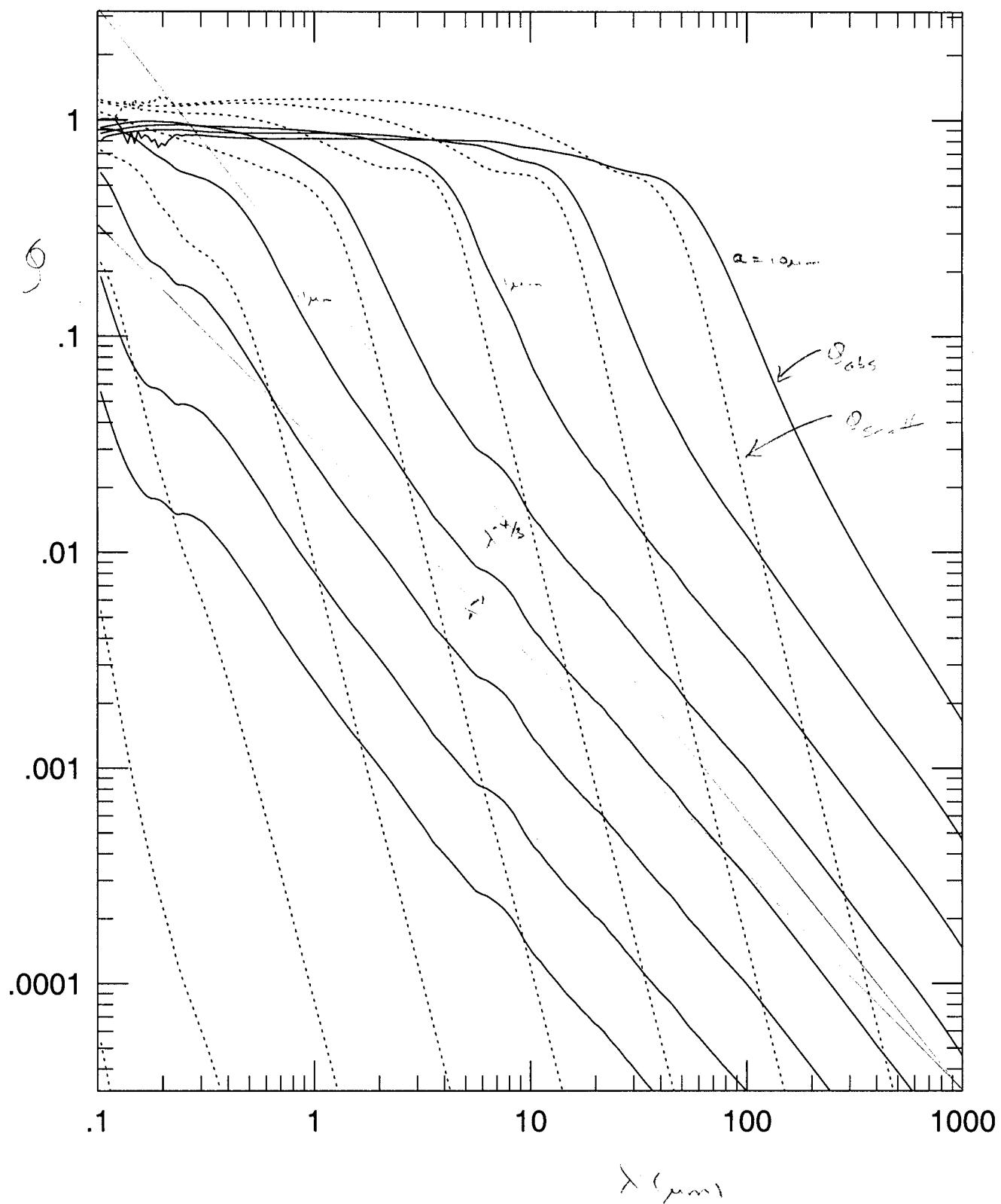
$$T^4 = \frac{R_*^2}{4d^2} T_*^4 \frac{\langle \langle f_{abs} \rangle \rangle}{\langle \langle f_{em} \rangle \rangle}$$

If  $\langle f_{abs} \rangle \propto v$ ,  $\langle \langle f_{abs} \rangle \rangle \propto T$

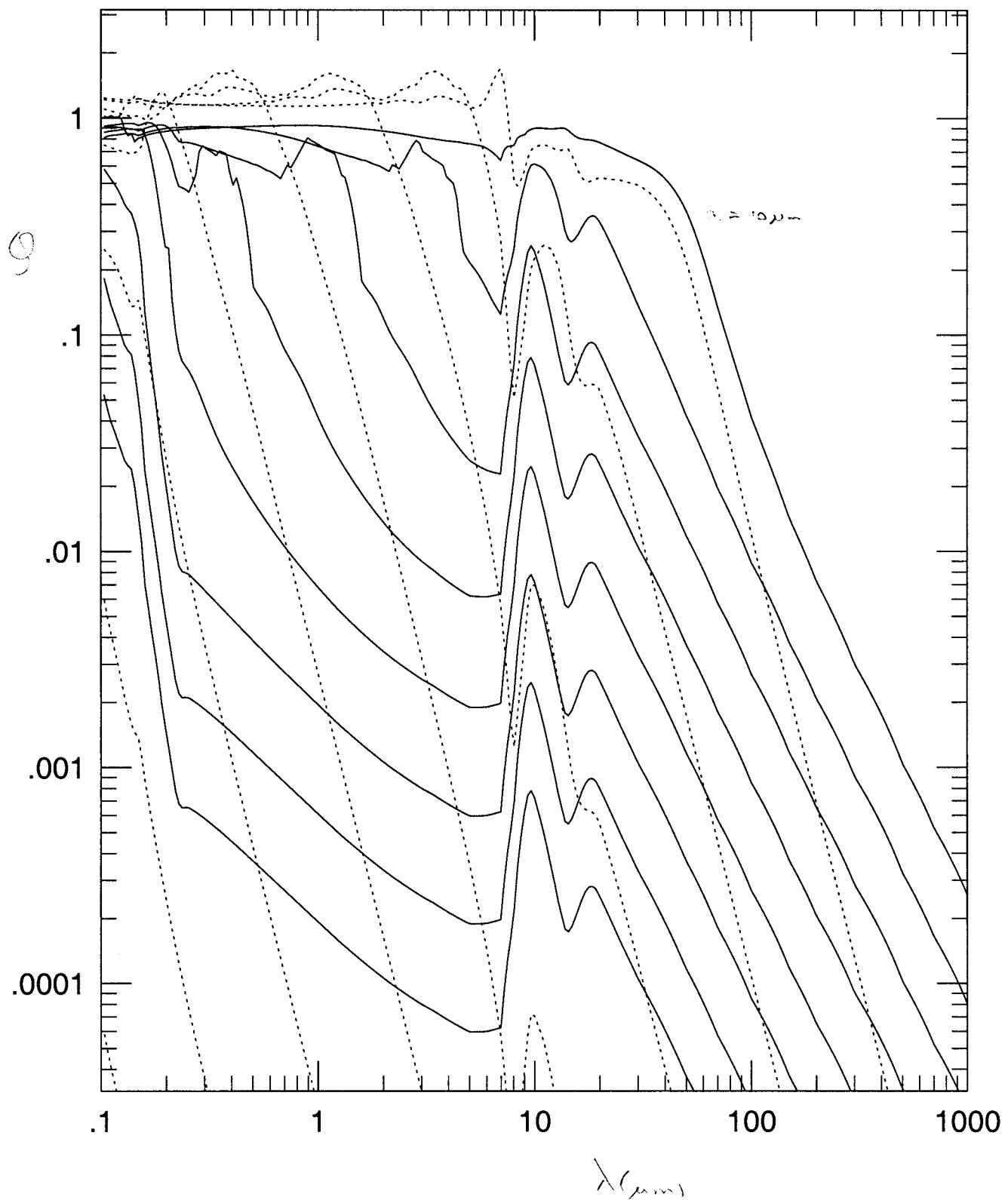
$$T^4 = \frac{R_*^2}{4d^2} T_*^4 \frac{T_*}{T} \Rightarrow T = T_* \left( \frac{R_*}{2d} \right)^{0.4}$$

$$\text{at 1 AU for } \Sigma_{\text{crit}}: \frac{R_*}{2d} = 2.3 \times 10^{-3} \quad T = 0.088 T_* \\ = 513 \text{ K}$$

Aerosol carbon



Amorphous Silicates



color difference.  
is; viz.

(5.88)

straightforward.  
alized extinction  
is

(5.89)

color excess as,

(5.90)

into the infrared,

(5.91)

by extrapolating  
on ratio  $A_\lambda/A_V$

i. this parameter

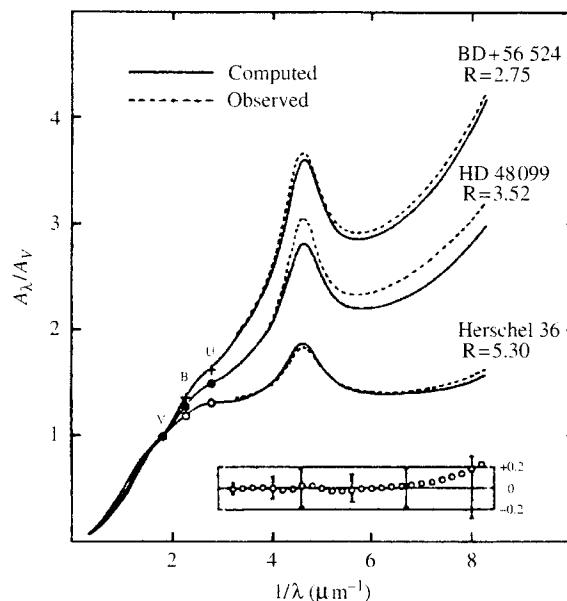


Figure 5.7 Three observed extinction curves are shown as a function of  $\lambda^{-1}$ . These curves show the range in wavelength behavior of the extinction laws in the interstellar medium. The solid lines show, for comparison, the computed parameterized extinction. The insert shows the deviations. Figure courtesy of J. S. Mathis; reprinted with permission from *Ann. Rev. Astron. Astrophys.*, **28**, p. 37, ©1990 by *Ann. Rev.* ([www.annualreviews.org](http://www.annualreviews.org)).

For a diffuse radia field :

$$P_{\text{abs}} = 4\pi a^2 \int Q_{\text{abs}}(v) \pi J_v dv$$

→ heated from all sides

Is this consistent w/  $\pi a^2 \int Q_{\text{abs}}(v) F_v dv$  ?

$$J_v = \frac{\pi R_*^2}{4\pi d^2} B_v(T_*)$$

$$F_v = \frac{L_v}{4\pi d^2} = \frac{4\pi R_*^2}{4\pi d^2} \pi B_v(T_*)$$

$$= 4\pi J_v. \quad \checkmark \quad (\text{as desired})$$

What if  $J_v = B_v(T)$  ?

$$4\pi a^2 \int Q_{\text{abs}}(v) \pi J_v dv = 4\pi a^2 \int Q_{\text{em}}(v) \pi B_v(T) dv$$

$$\Rightarrow T = T_* \quad \checkmark \quad (\text{as required})$$

How to set up program to calculate  $T$  inside an optically thick dust disk?

Need to know  $J_v$

I followed many rays through disk, integrating egn. of rad. xfer:

$$\frac{dI_v}{ds} = - \left[ \sum_i \pi a_i^2 Q(a_i, v) n_i \right] I_v + \sum_i \pi a_i^2 g(a_i, v) n_i B_v(T_i)$$

or for finite  $\Delta s$ :

$$\Delta I_v = e^{-\alpha(v) \Delta s} I_v + (1 - e^{-\alpha(v) \Delta s}) S_v$$

$$S_v = \frac{\sum_i \pi a_i^2 g(a_i, v) n_i B_v(T_i)}{\sum_i \pi a_i^2 g(a_i, v) n_i}$$

By averaging rays through each grid cell in disk  
I calculated  $J_\nu$

Then I recalculated  $T$  and repeated until  
it converged

How to include scattering?

Negligible for emitted radiation in IR

Important for star light:

$$\frac{dI_\nu}{ds} = - \sum_i \pi a_i^2 (\alpha_{abs} + \alpha_{scat}) n_i I_\nu$$

$$+ \sum_i \pi a_i^2 \beta_{abs} n_i B_\nu(T)$$

$$+ \sum_i \pi a_i^2 \alpha_{scat} n_i J_\nu$$

$$\text{or } S_\nu = \frac{\alpha_{abs} B_\nu + \alpha_{scat} J_\nu}{\alpha_{abs} + \alpha_{scat}}$$

What next?

To calculate the disk structure, I need  $T_{gas}$   
so I can solve the eqn. of hydrostatic equil.

$$\frac{dP}{dz} = - \gamma \rho$$

$$\text{where } P = n k T, \rho = \langle m \rangle n$$

$$\text{At } L \ll n, T_{gas} = T_{disk}$$

At low  $n$ , w/ other heating sources, need to  
calculate  $T_{gas}$  by balancing heating + cooling.

## Small grains + PAHs

Grain size distribution:

$$MRN: \frac{dn}{da} \propto a^{-3.5} \rightarrow a_{\max} \sim 0.1 \mu\text{m}$$

$$\pi a^2 \frac{dn}{d\ln a} \propto a^{-0.5} : T_{\nu} \text{ dominated by small grains}$$

$$\frac{4\pi}{3} a^3 \frac{dn}{d\ln a} \propto a^{+0.5} : T_{\text{FIR}} \text{ dominated by } a_{\max}$$

grain T falls w/  $a$  b/c  $E_{\text{FIR}}/T_{\nu}$  increases

Sellgren: reflection nebulae w/ NIR albedo  $> 1$

explanation: emission w/  $T > T_{\text{equil.}}$

$\Rightarrow$  v. small grains

Similar effect: 3.3  $\mu\text{m}$  emission from dust illuminated  
by OV - PAHs

PAHs: carbon bonds, benzene  $\text{sp}^2, \text{sp}^2 + \text{p}$

Excitation:  $\pi \rightarrow \pi^*$  electronic excitation

followed by radiative or non-radiative xition

to excited vib. state in another electronic state

describe vibrational excitation by  $\frac{1}{2}kT = E/\text{mode}$

not quite right: Canonical vs. microcanonical ensembles

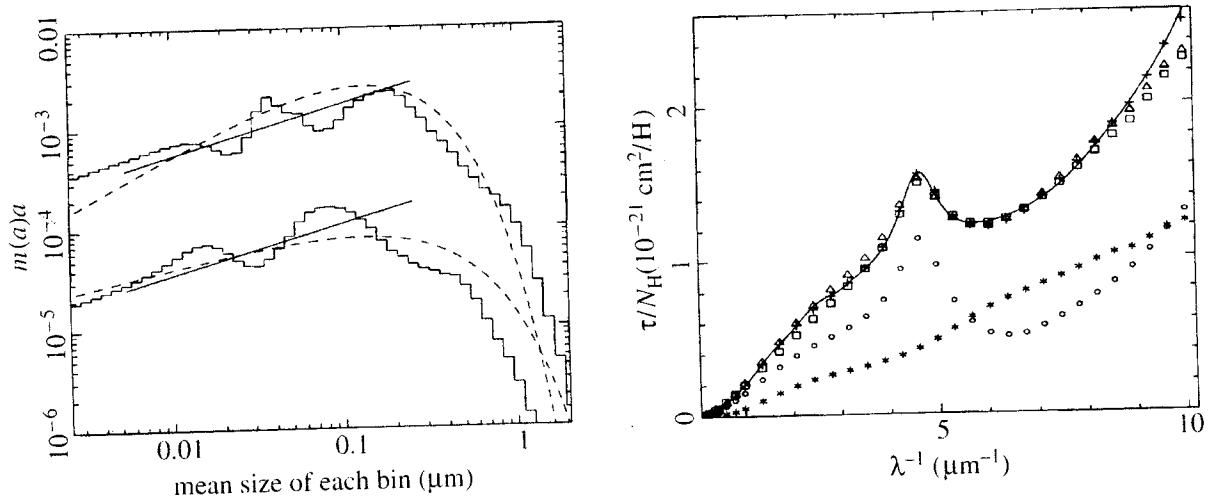


Figure 5.15 The interstellar grain size distribution – plotted as mass fractions – derived from different types of fits to the observed extinction curve (left panel). The top histogram and dashed line show silicates; the bottom histogram and dashed line shows graphite, displaced downwards by factor 10. The solid lines show for comparison the simple MRN power law size distributions. The calculated extinction curve is compared to the observations in the right panel. The contributions due to silicates (\*) and graphite (○) are shown separately. Figure reprinted with permission from S.-H. Kim, P. G. Martin and P. D. Henry, 1994, *Ap. J.*, **422**, p. 164.

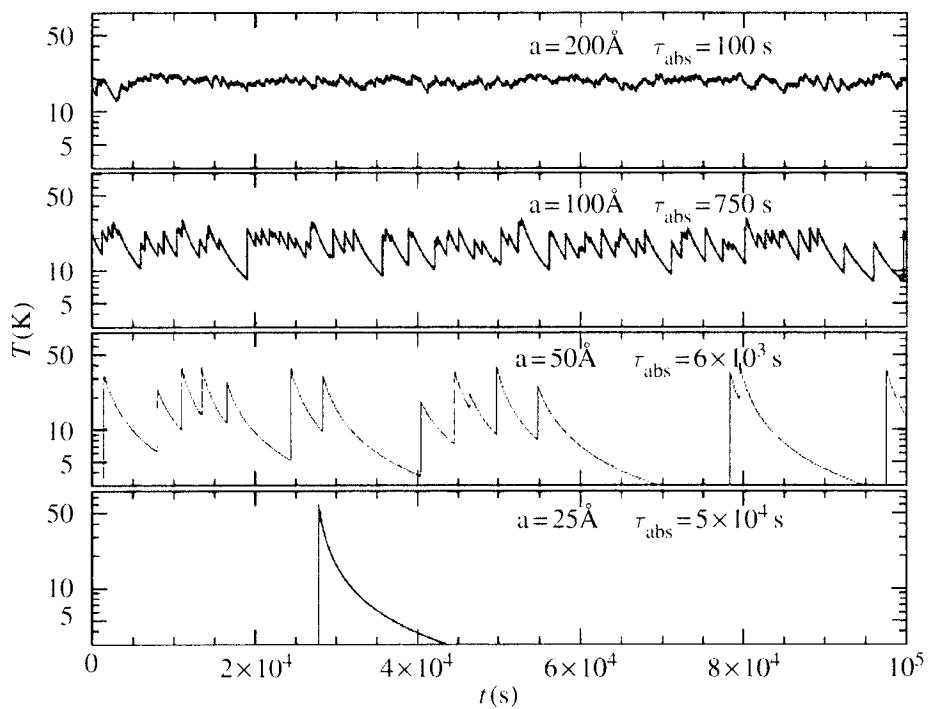


Figure 6.5 The time-dependent behavior of the temperature for various sizes of the species. The time axis corresponds to approximately one day. These calculations pertain to the diffuse ISM. Figure courtesy of B. T. Draine; reprinted, with permission, from the *Ann. Rev. Astron. Astrophys.*, **41**, p. 241 ©2003 by *Ann. Rev.* ([www.annualreviews.org](http://www.annualreviews.org)).

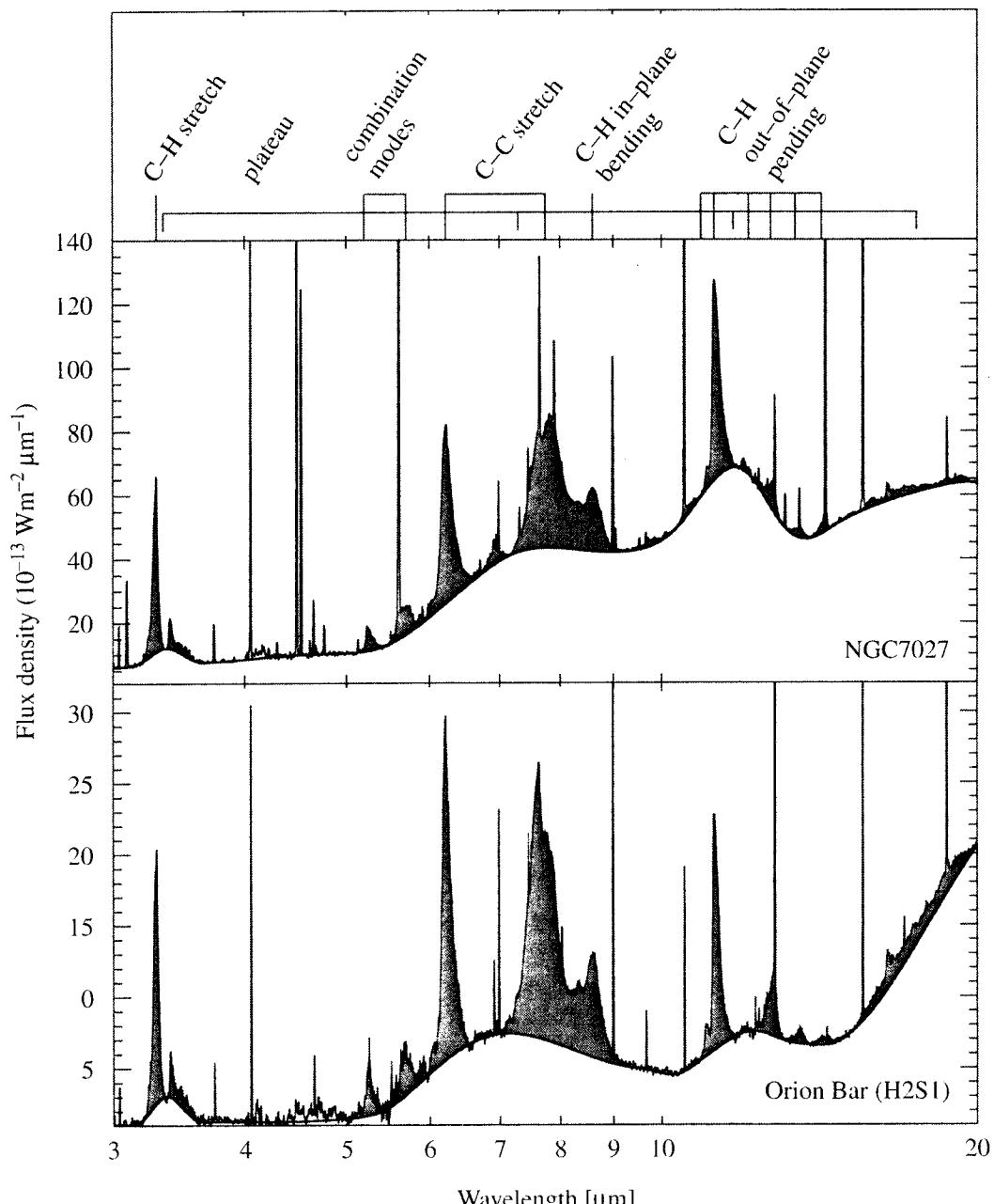


Figure 6.16 The 3–15  $\mu\text{m}$  spectra of the PDRs associated with the Orion Bar and NGC 7027, illustrating the ubiquitous nature as well as the richness of the IR emission features. The IR emission features are shaded. The narrow lines are HI recombination lines,  $\text{H}_2$  pure rotational and rotational–vibrational lines, and atomic fine-structure lines. The top panel indicates the PAH vibrational modes associated with each feature. Note the presence of broad plateaus underneath the narrow emission features. Figure adapted from E. Peeters, *et al.*, 2002, *A. & A.*, **390**, p. 1089.

These modes “measure” nearest-neighbor interactions to a large extent. Hence, the infrared signatures of carbonaceous solids with an abundant aromatic component – such as coal, amorphous carbon soots, and charcoals – also bear a general