

Carbon Monoxide Line Emission as a CMB Foreground:

Tomography of the Star Forming Universe

Rashid Sunyaev

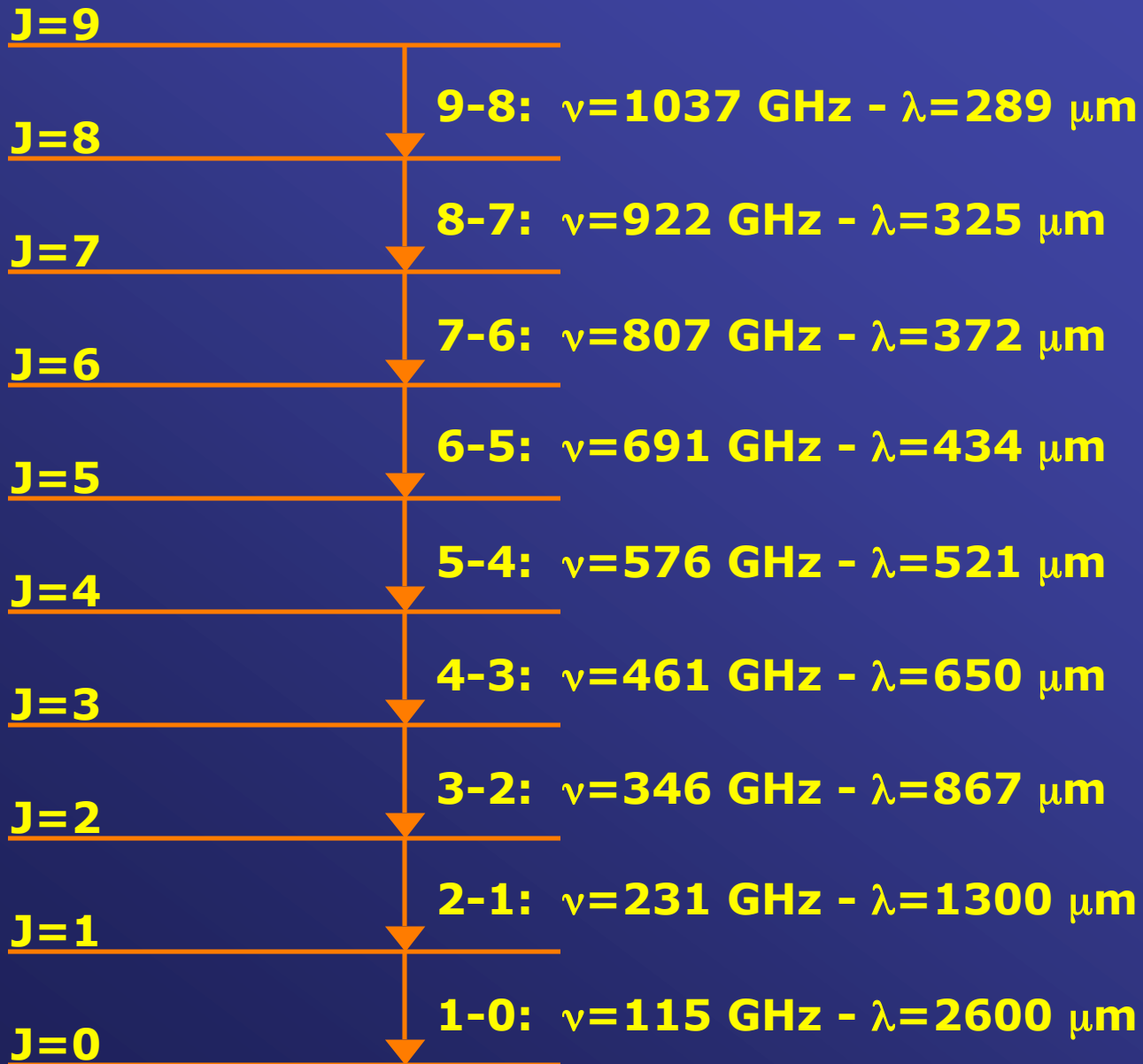
In collaboration with **Mattia Righi** and
Carlos Hernandez-Monteagudo



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Space Research Institute (IKI), Moscow**

October 29, 2009

CO rotational levels



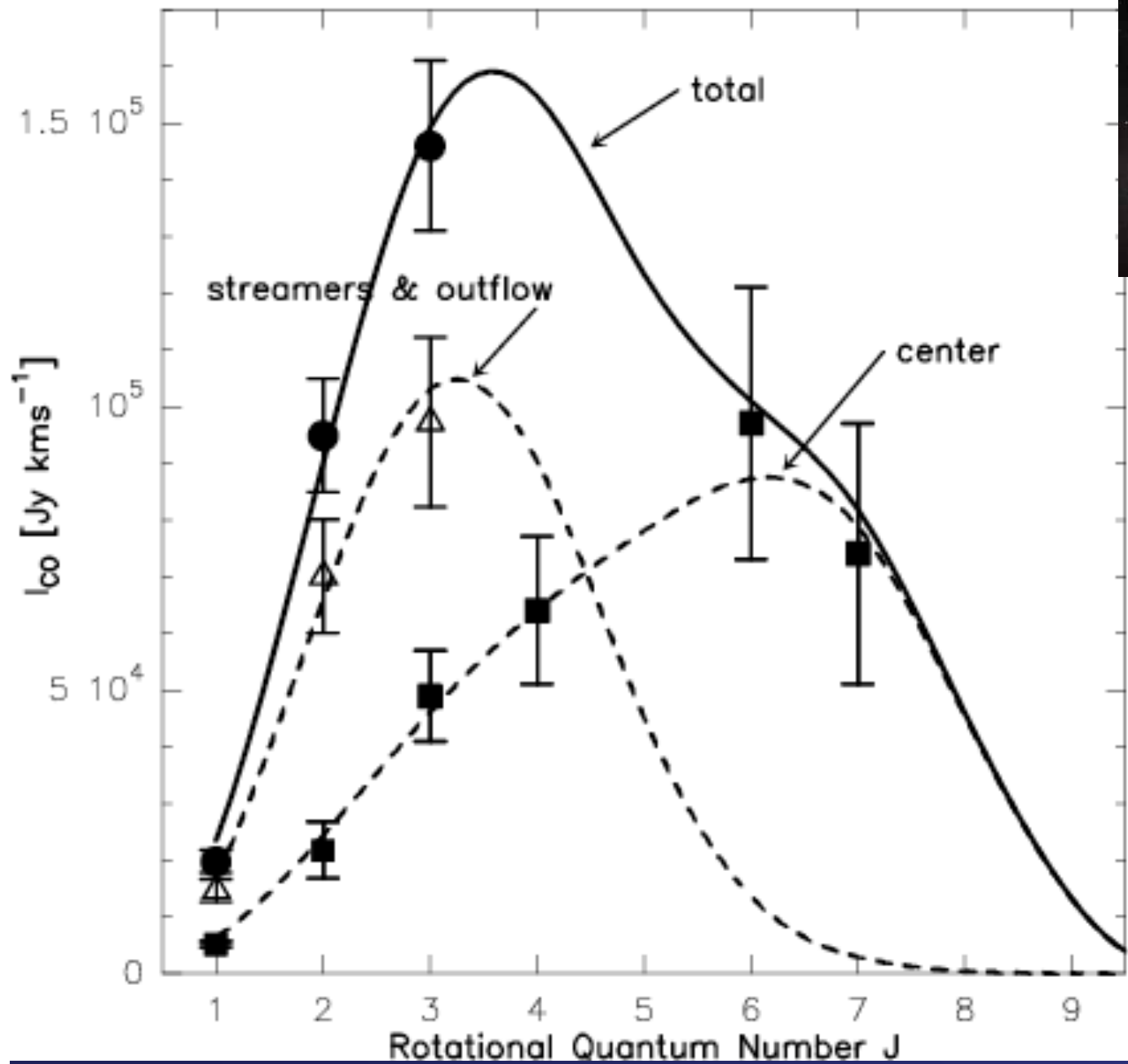
Population III stars produce very effectively C and O

Heger and Woosley,
2002

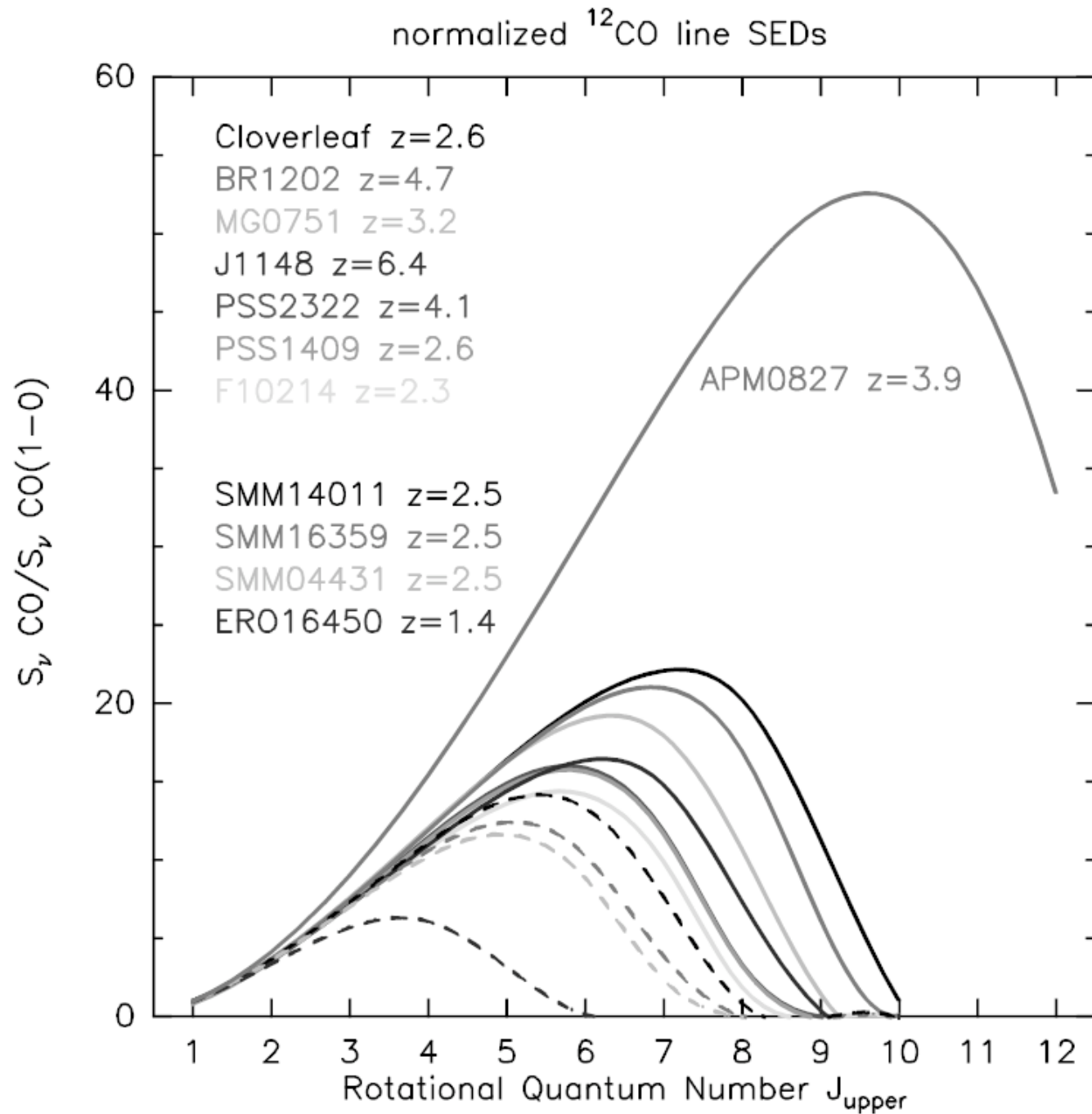
Are there traces of CO emission in early Universe on CMB maps?

- practically

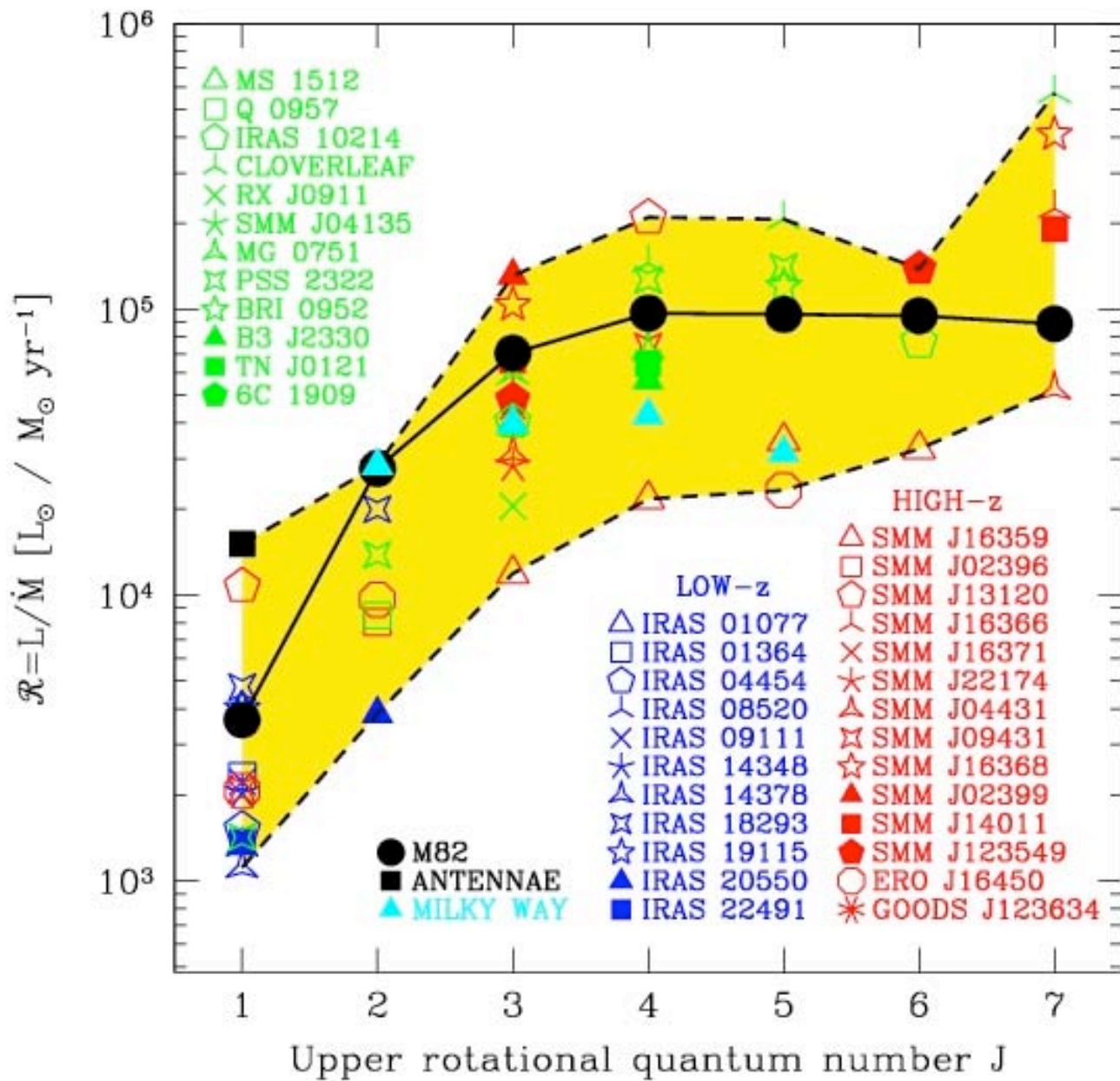
total ^{12}CO flux density in M82



Weiss and Scoville



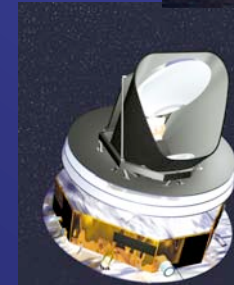
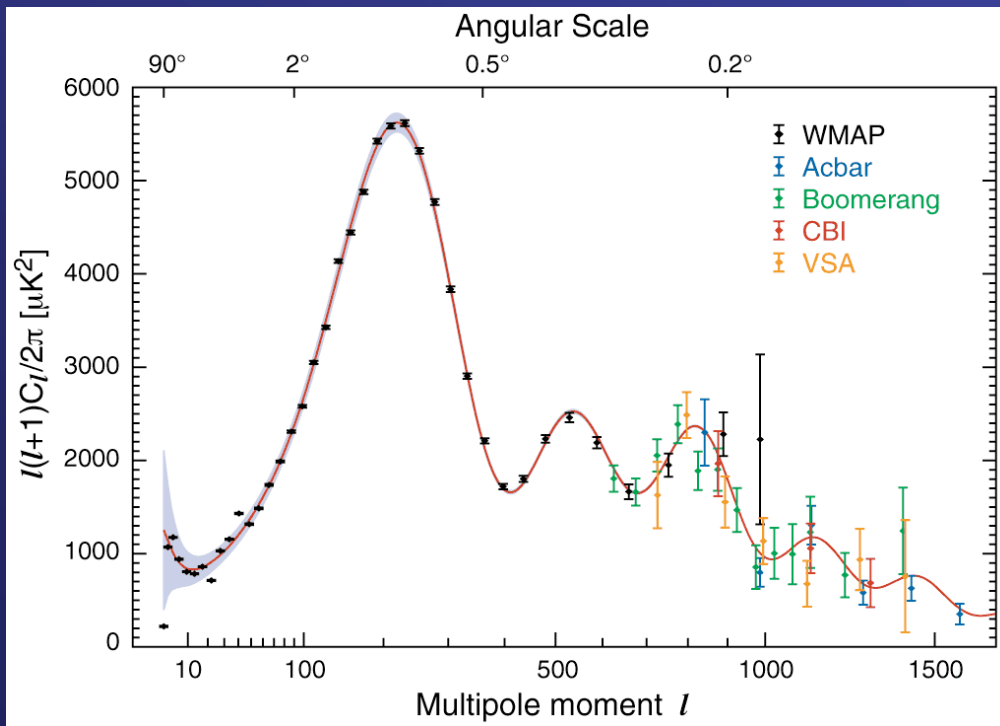
Distribution of \mathcal{R} ratios – CO lines



Motivation

Forthcoming CMB experiments will reach very high angular resolution, probing the **high multipole** region ($l \sim 10^3 - 10^4$) of the angular power spectrum

Hinshaw et al. (2007)



SPT

PLANCK



ACT

High frequency channels (150-350 GHz) will be affected by thermal dust emission from star-forming galaxies
Clustering of such sources could be one of the most important **foregrounds** for these channels

CO lines are more important than **dust** as a foreground at low frequencies (15 – 90 GHz):

WMAP, Planck LFI, Quiet and planned modification of CBI

CO line foreground observations require better spectral resolution than WMAP, PLANCK LFI, SPT and ACT have

In principle this is possible

Cosmic Background Imager had 10 spectral channels at 30 GHz

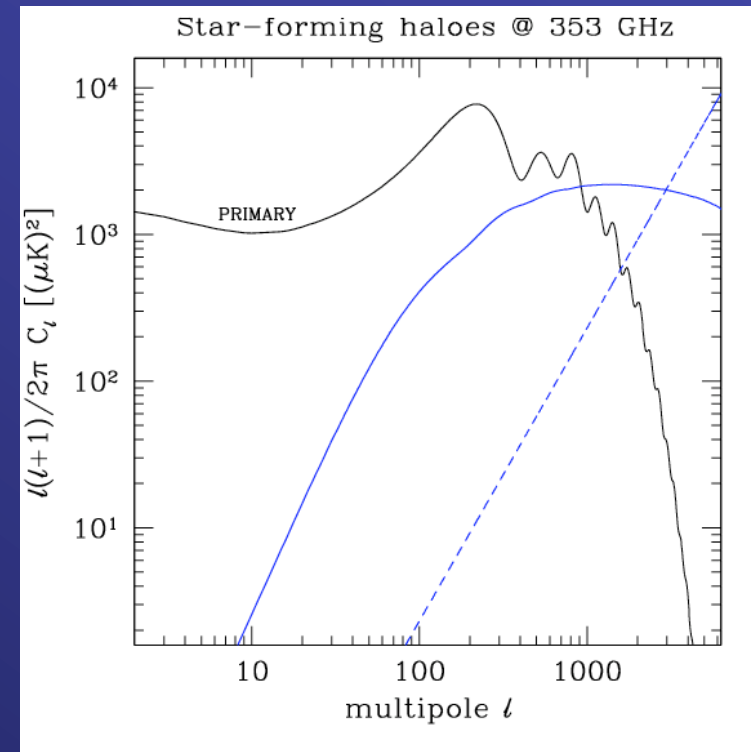
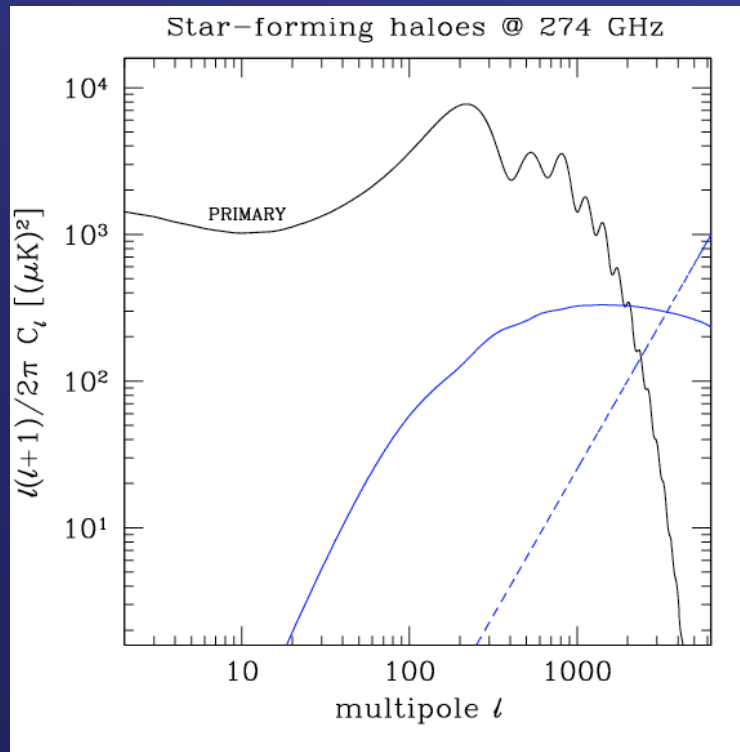
SZ-Array in Owens Valley, California has 145 spectral channels at 30 GHz



Results: power spectra

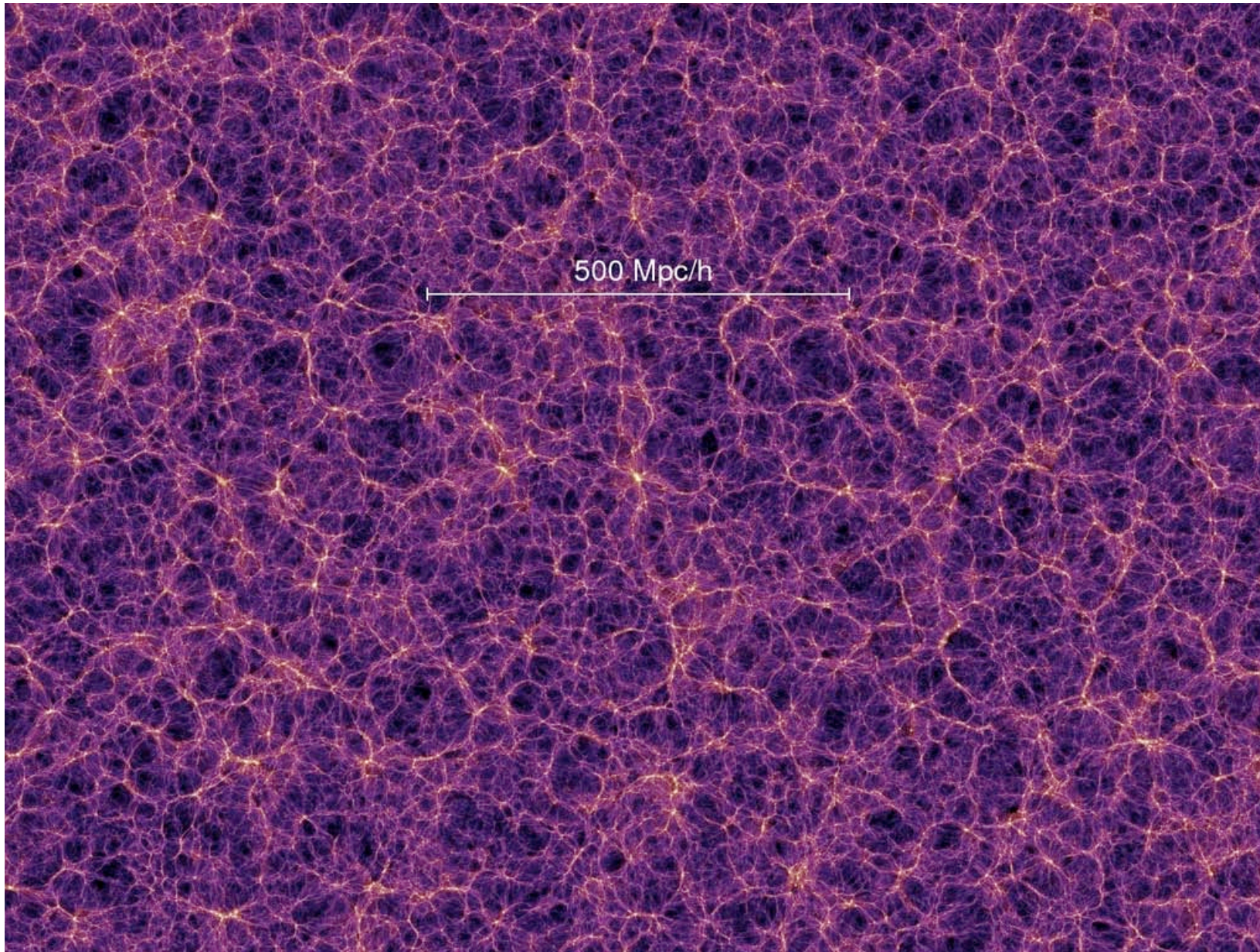
Two terms: **POISSON** + **CORRELATION**

$$\langle \Delta I_\nu(\vec{n}_1) \Delta I_\nu(\vec{n}_2) \rangle = \sum_l \frac{2l+1}{4\pi} (C_l^P + C_l^C) P_l(\vec{n}_1 \cdot \vec{n}_2)$$



Dust emission will be one of the most important small-scale foregrounds for the high-frequency channels of the future experiments

Good agreement with earlier estimates (Haiman & Knox 00)



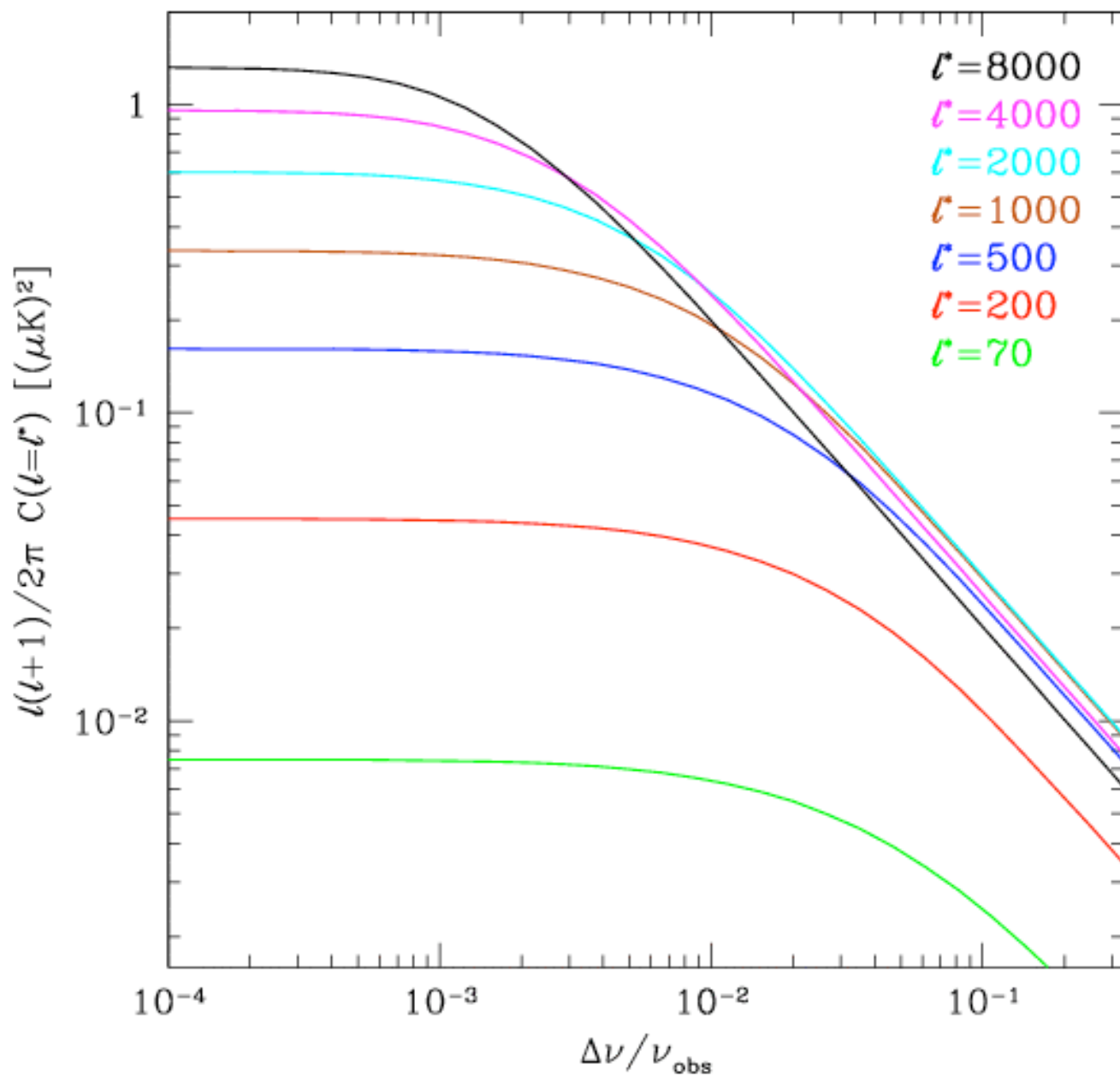
Best
Contrast!!!

spectral
resolution
 10^{-3}
corresponds
to similar
 dz/z

higher
Cl's

This slice through the density field is 15 Mpc/h thick (Volker Springel).

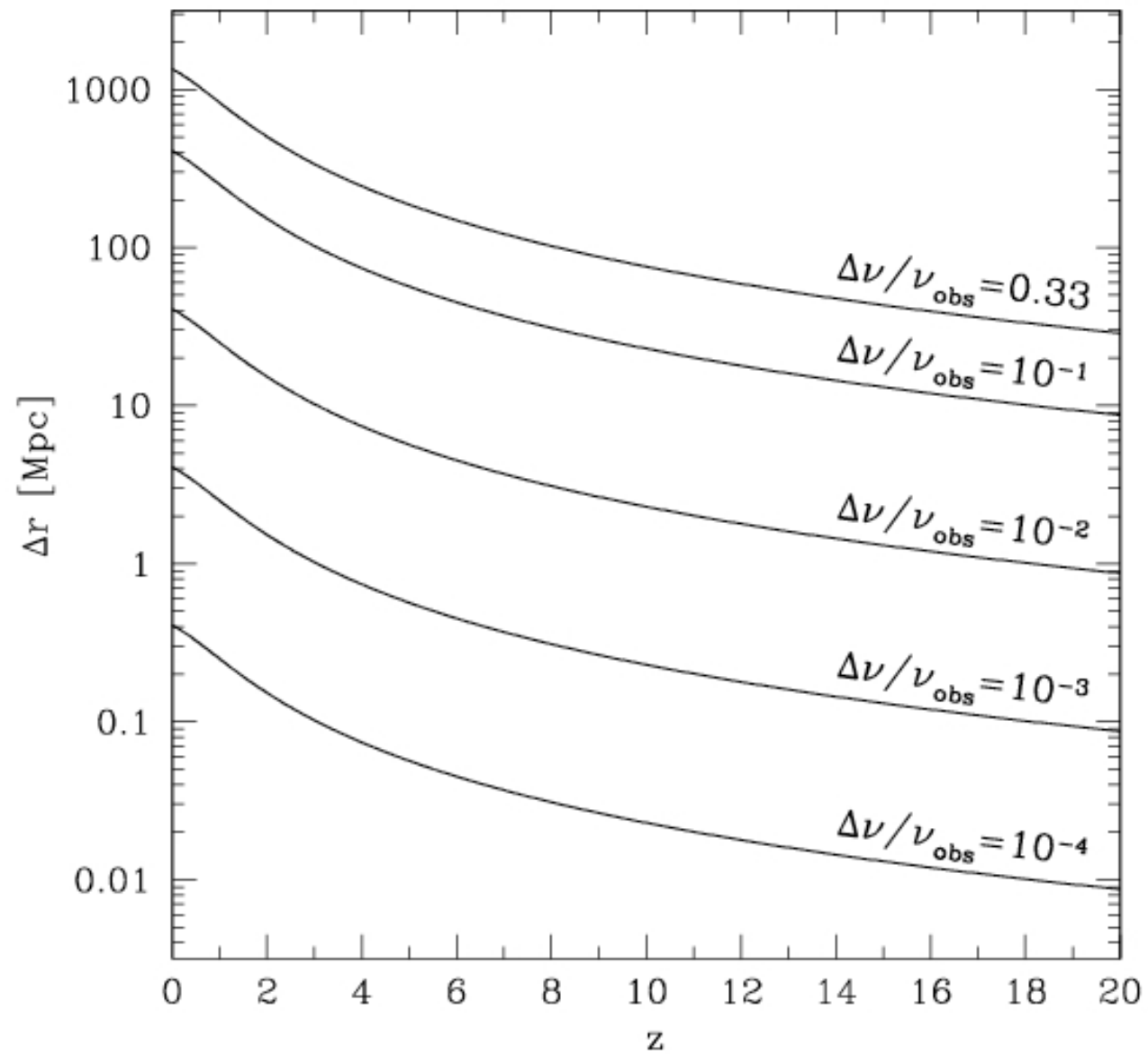
Spectral resolution – CO 1–0 line – 30 GHz



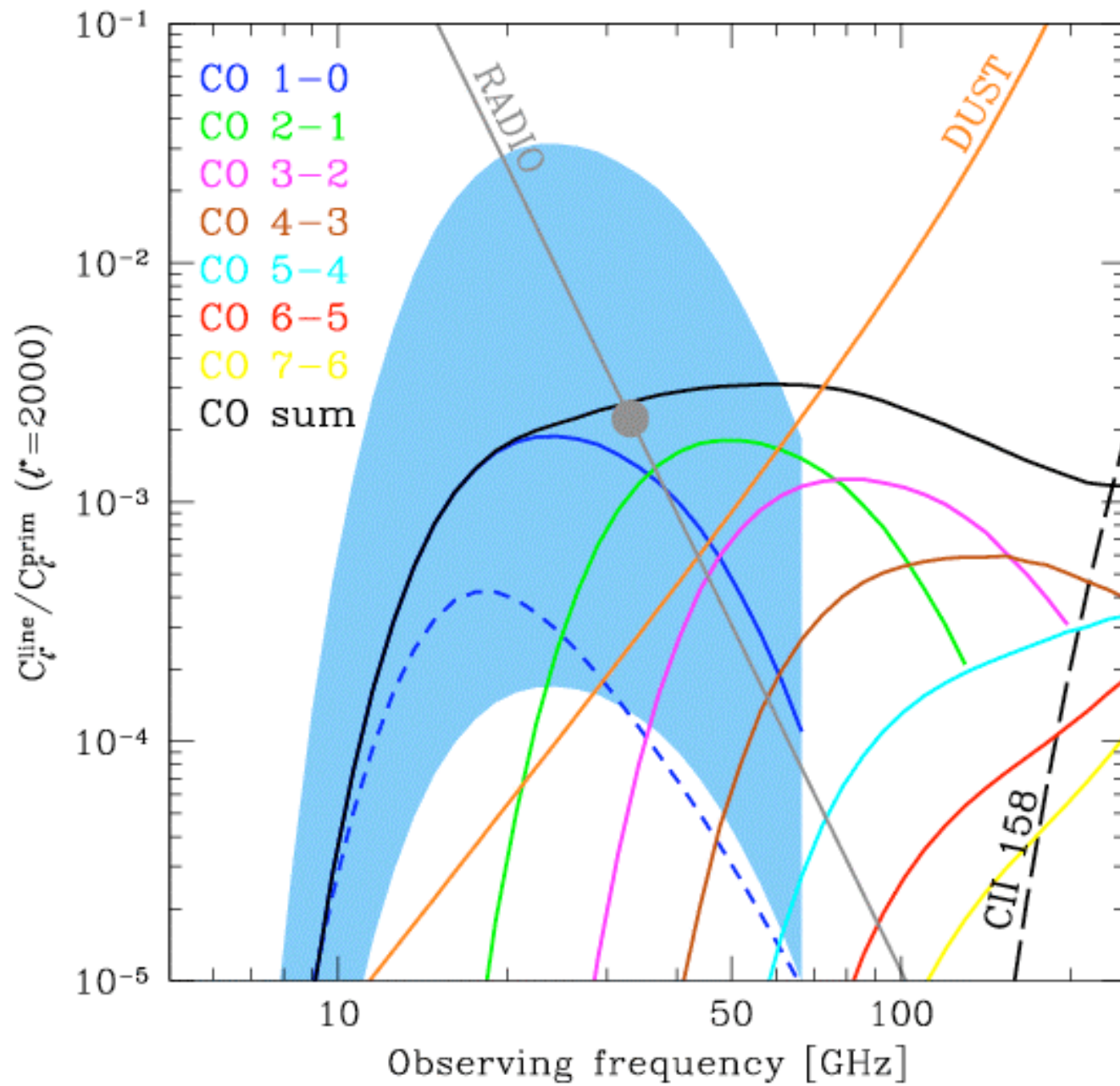
Increase of spectral resolution 300 times provides 100 – 300 times increase in sensitivity of CMB measurement

But only for high $l > 1000$

Comoving width of the slice



Frequency dependence - $\Delta\nu/\nu_{\text{obs}} = 10^{-3}$



Outline

Merging model for star-formation inside haloes

Observational data for CO line emission from
merging galaxies

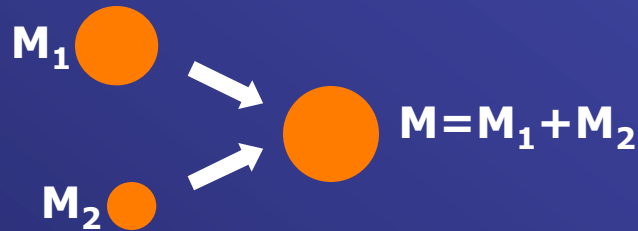
Observational tests and model calibration

Results: power spectra
of angular fluctuations

Star formation from mergers

BASIC IDEA (Barkana & Loeb, 2000)

Derive star formation as the rate of accretion of baryonic mass into new haloes



Stellar mass produced in the merger is a function of the merging masses, limited by cooling time

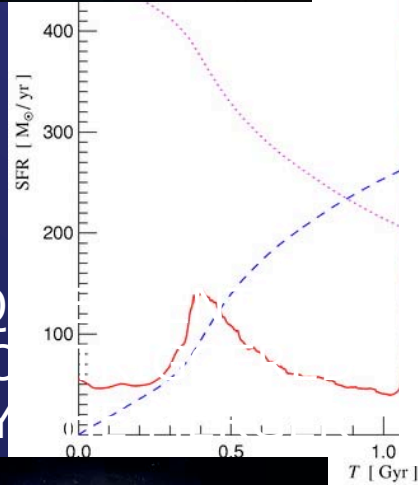
$$M_*^1 = \frac{\Omega_b}{\Omega_m} \eta M_1 \frac{M_2}{M/2} \quad M_*^2 = \frac{\Omega_b}{\Omega_m} \eta M_2 \frac{M_1}{M/2}$$

For a given cosmological model, rate of merging is derived in the context of the **extended-PS theory** (Lacey & Cole, 1993)

Future clusters of galaxies and superclusters contained more merging objects

Star formation timescales

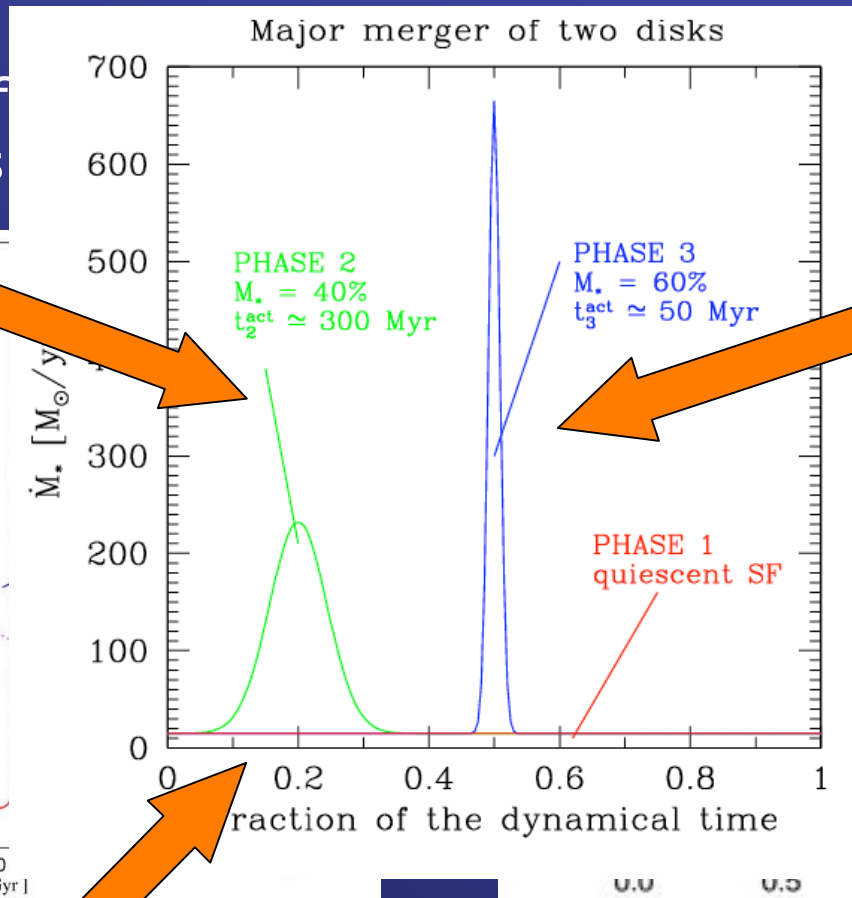
FIRST PASSAGE
MODERATE BURST



QUENCHED BY

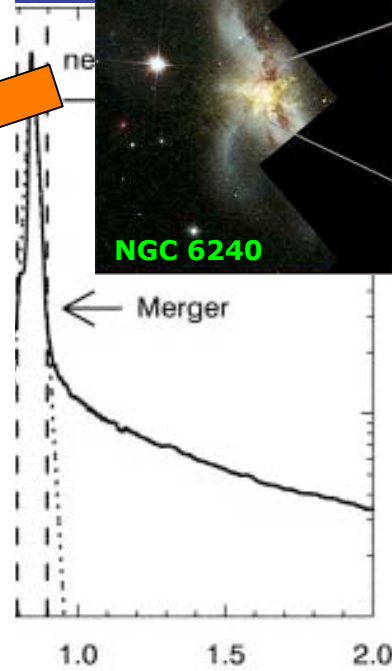


Springel & Hernquist (2005)

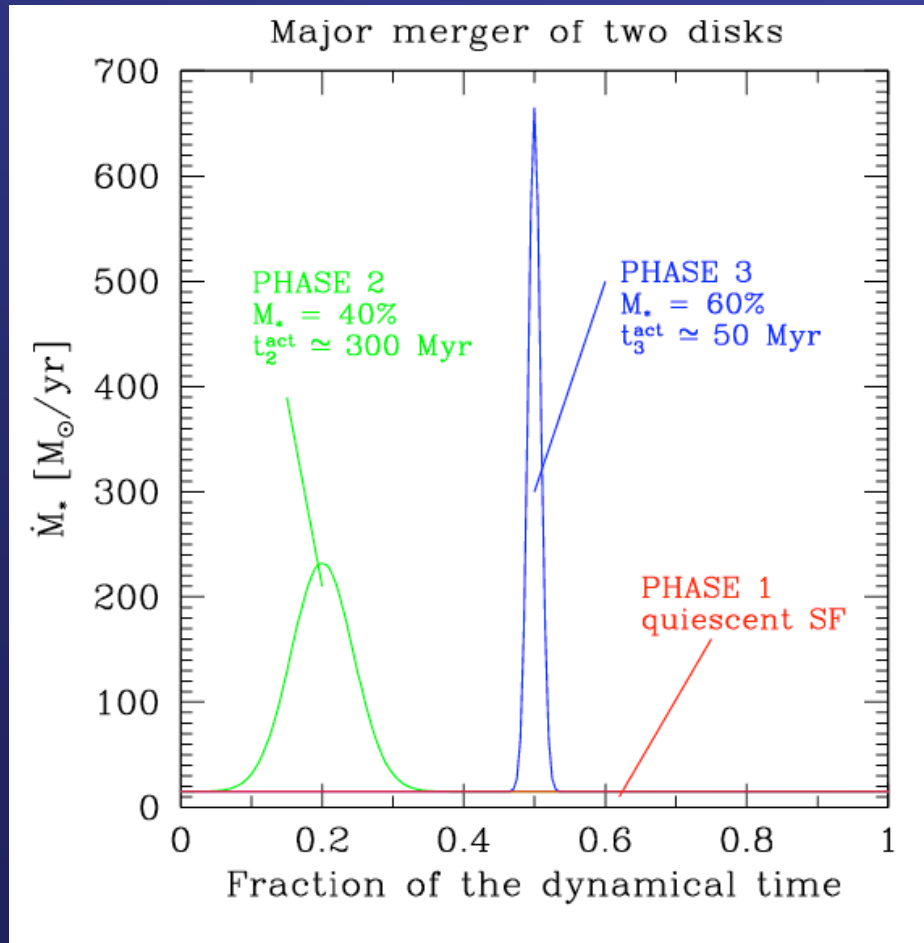


Robertson et al. (2006)

FINAL MERGER
STRONG BURST



Star formation timescales



TWO SEPARATE POPULATIONS

40% of M_* is produced during first passage (due to tidal deformation, shocks and gas compression)

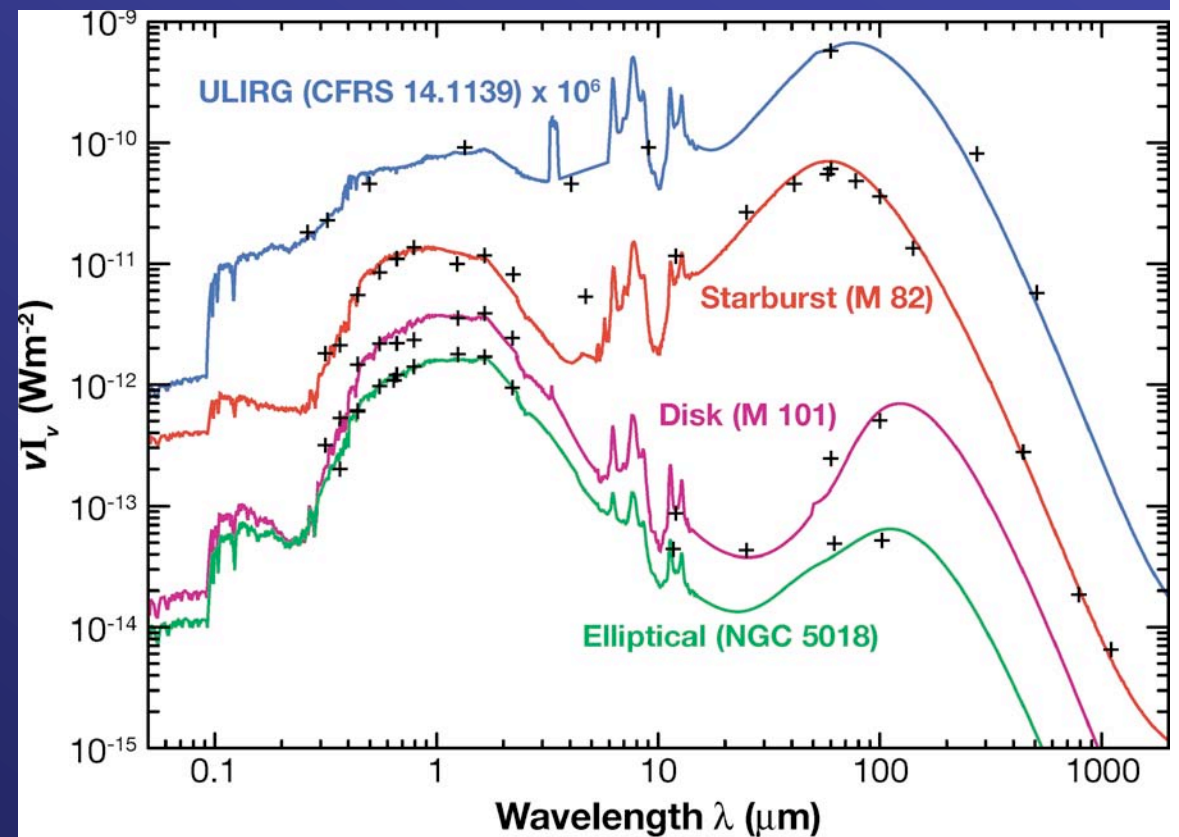
$$t_{\text{char}} \sim 3 \times 10^8 \text{ years}$$

60% of M_* is produced in the final coalescence (strong burst when the two galaxies finally merge together)

$$t_{\text{char}} \sim 5 \times 10^7 \text{ years}$$

Spectral Resolution 10^{-3}

- $DUST/CO(2-1) = 0.005$
- $DUST/CO(3-2) = 0.009$
- $DUST/CO(4-3) = 0.020$
- $DUST/CO(5-4) = 0.054$
- $DUST/CO(6-5) = 0.122$
- $DUST/CO(7-6) = 0.236$



Model summary

MERGING MODEL

Prescription for star formation in major mergers

$$L_{\text{line}} = \mathcal{R} \text{SFR}$$

RESULTS

Angular power spectra
Spectral resolution dependence
Frequency dependence
Foregrounds comparison

Line luminosity: assumed linear scaling with star formation rate

+

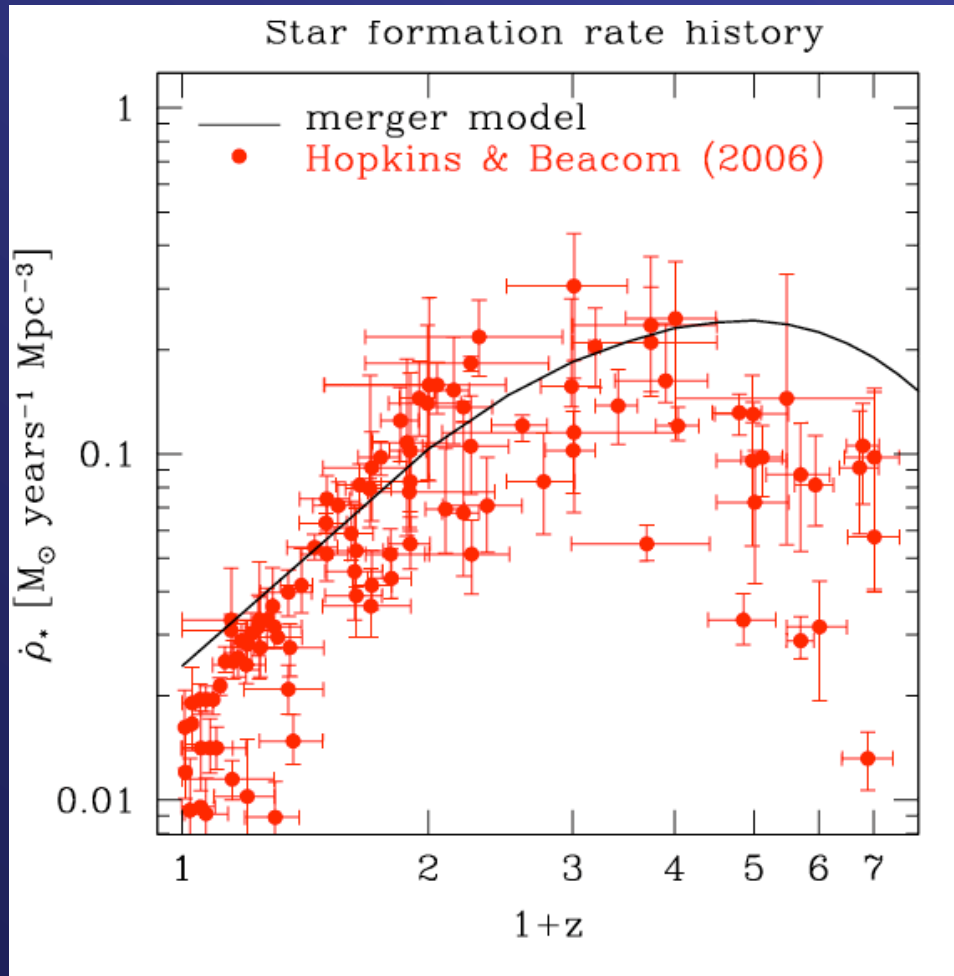
OBSERVATIONAL DATA

to calibrate the proportionality (\mathcal{R} ratio):
M82, Antennae, Milky Way, low-z IRAS galaxies, high-z sub-mm galaxies, QSO, LBGs, radio galaxies

=

\mathcal{R} ratio for CO transition (J+1=1-7) and atomic species (CII, NII, OI...)

Cosmic star-formation rate

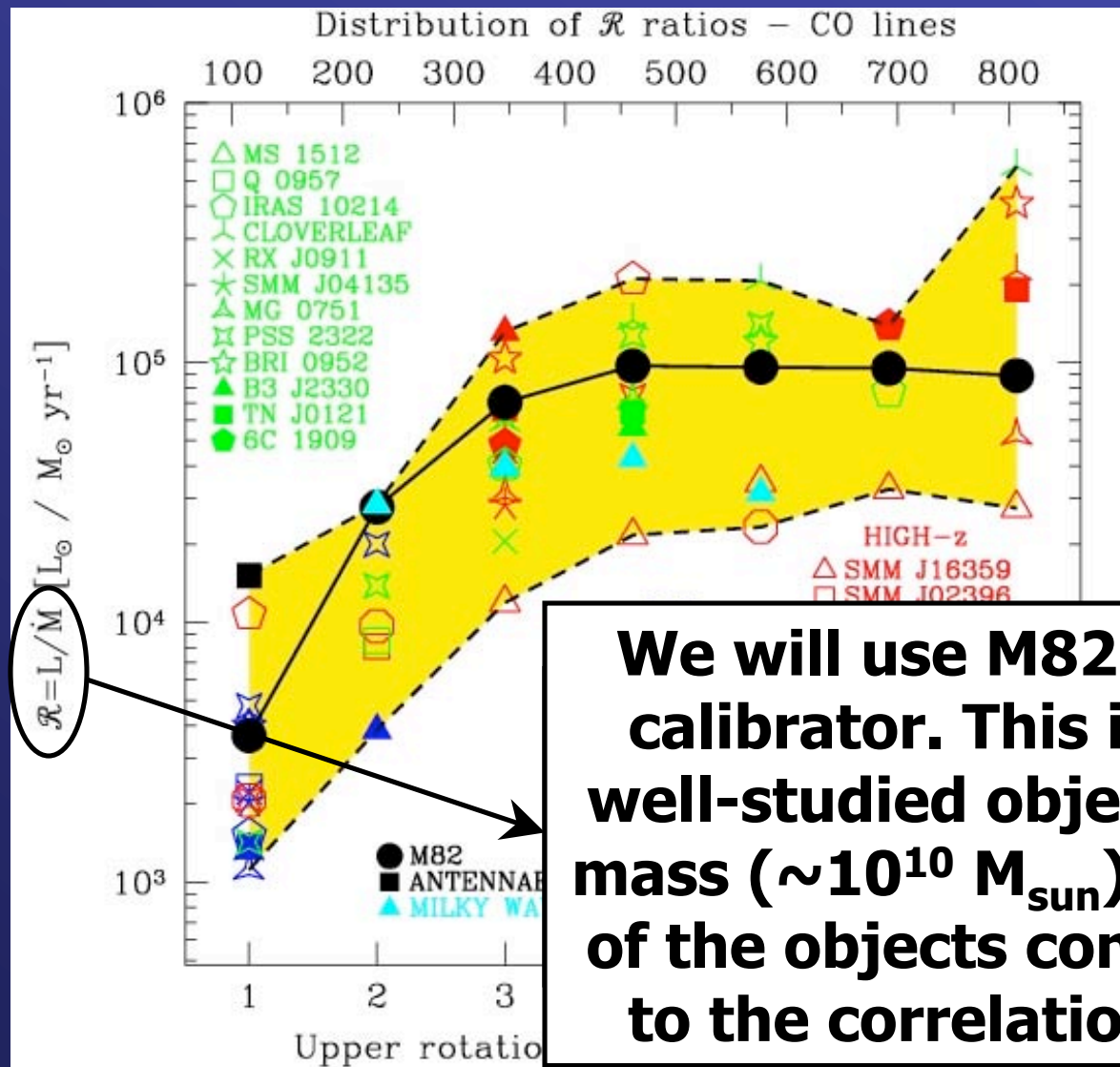


The star formation model can be compared with observations by computing cosmic SFR (Madau plot)

$$\dot{\rho}_*(z) = \int dM_1 \int dM \frac{dN_{\text{merg}}}{dM dt} \frac{dn}{dM_1} M_*(M_1, M_2, z)$$

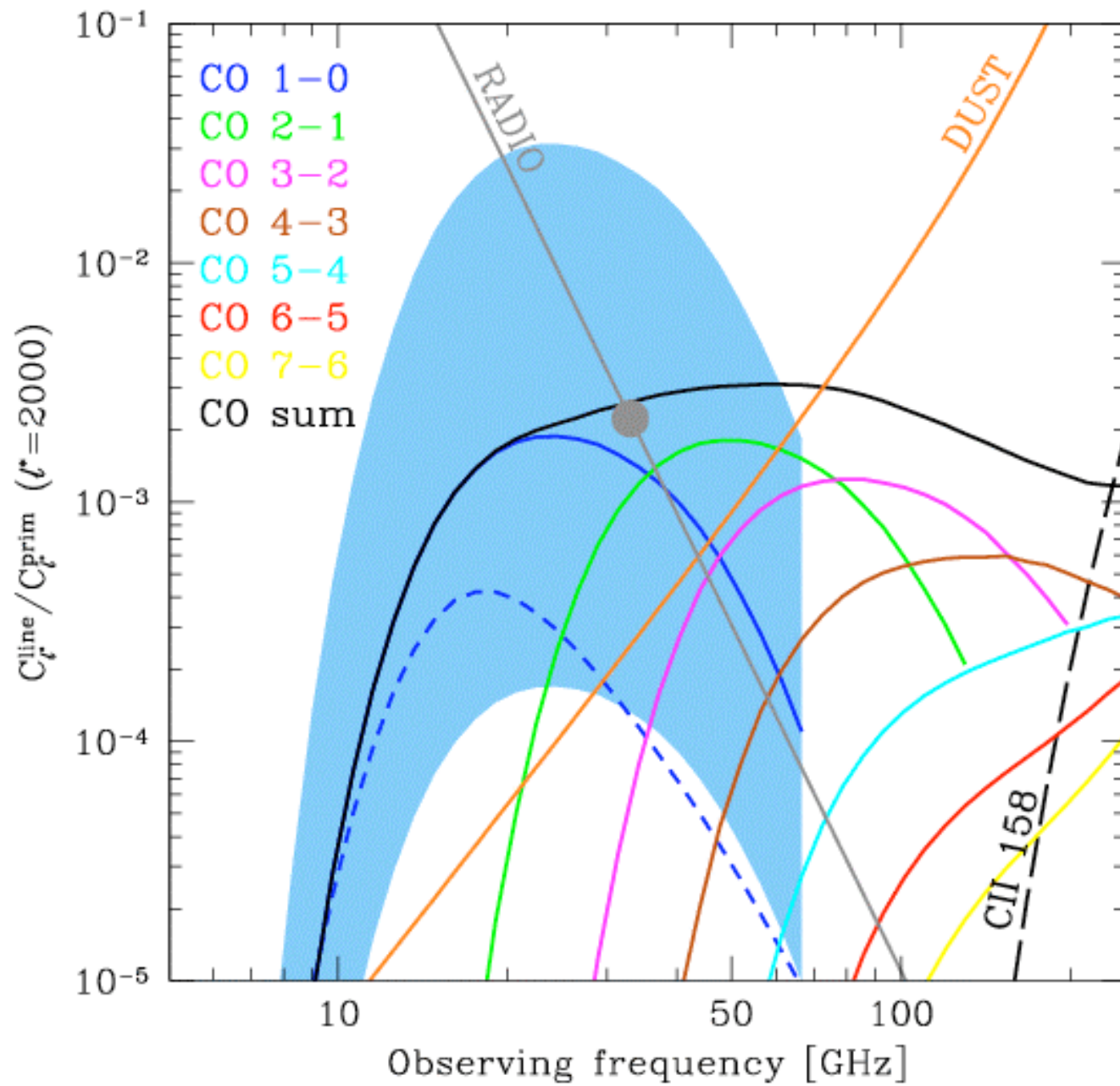
Rather good agreement is obtained with a star formation efficiency $\eta=5\%$

OBSERVATIONAL SAMPLE

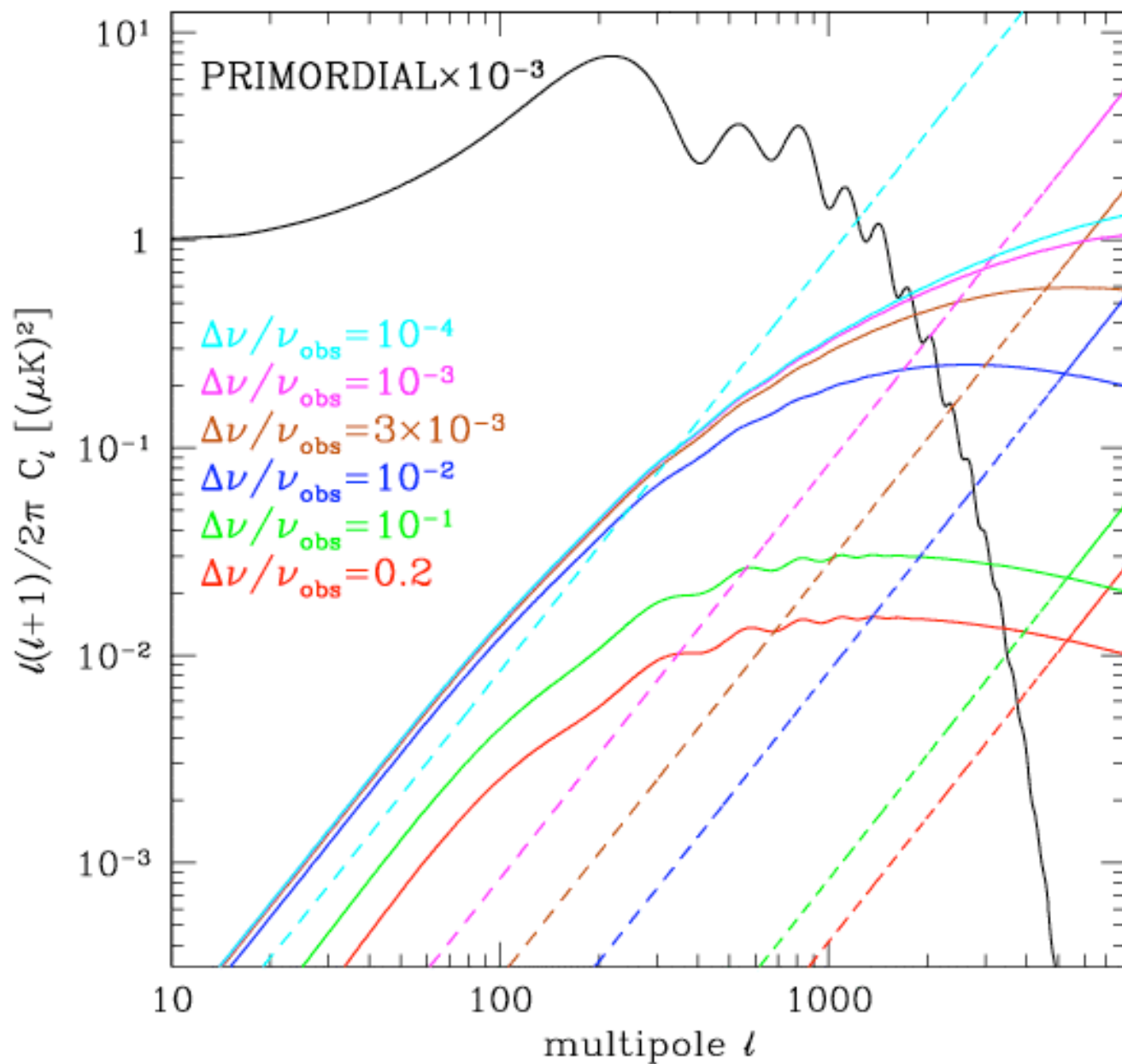


We will use M82 as main calibrator. This is a very well-studied object and its mass ($\sim 10^{10} M_{\text{sun}}$) is typical of the objects contributing to the correlation signal

Frequency dependence - $\Delta\nu/\nu_{\text{obs}} = 10^{-3}$



Spectral resolution – CO 1–0 line – 30 GHz



Correlated contribution
(solid curves)

Poisson contribution
(dashed lines)
behaves differently

too high spectral
resolution is useless
for correlated
contribution

CORRELATION TERM

Stars are formed in highly clustered overdense regions.
The correlation (or clustering) term accounts for this

It depends on the halo number density **squared**.
On the other hand, the Poisson term is linear with the
density.



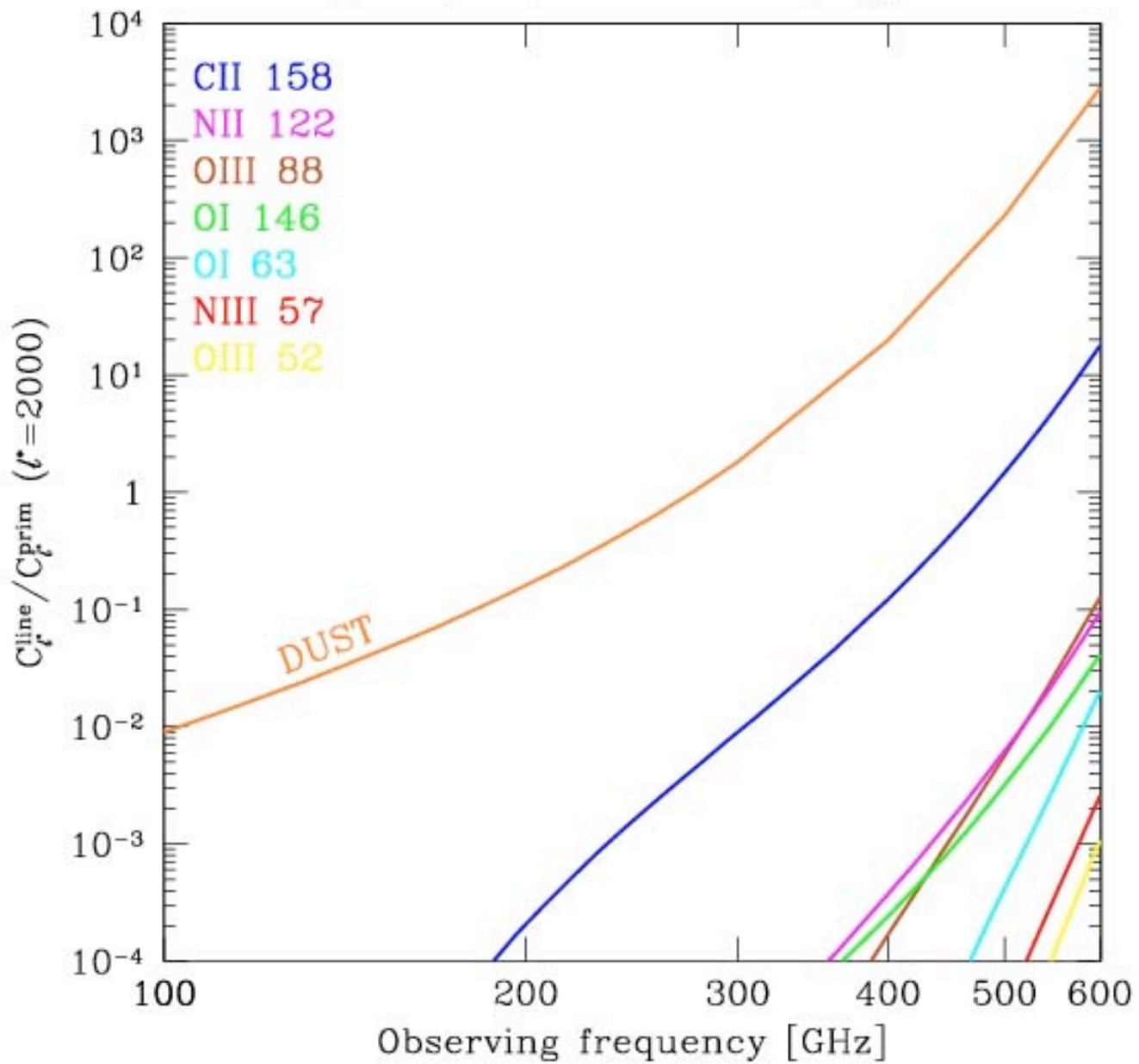
For large number densities (not too rare objects) the
clustering term will dominate over the Poisson

Conclusions

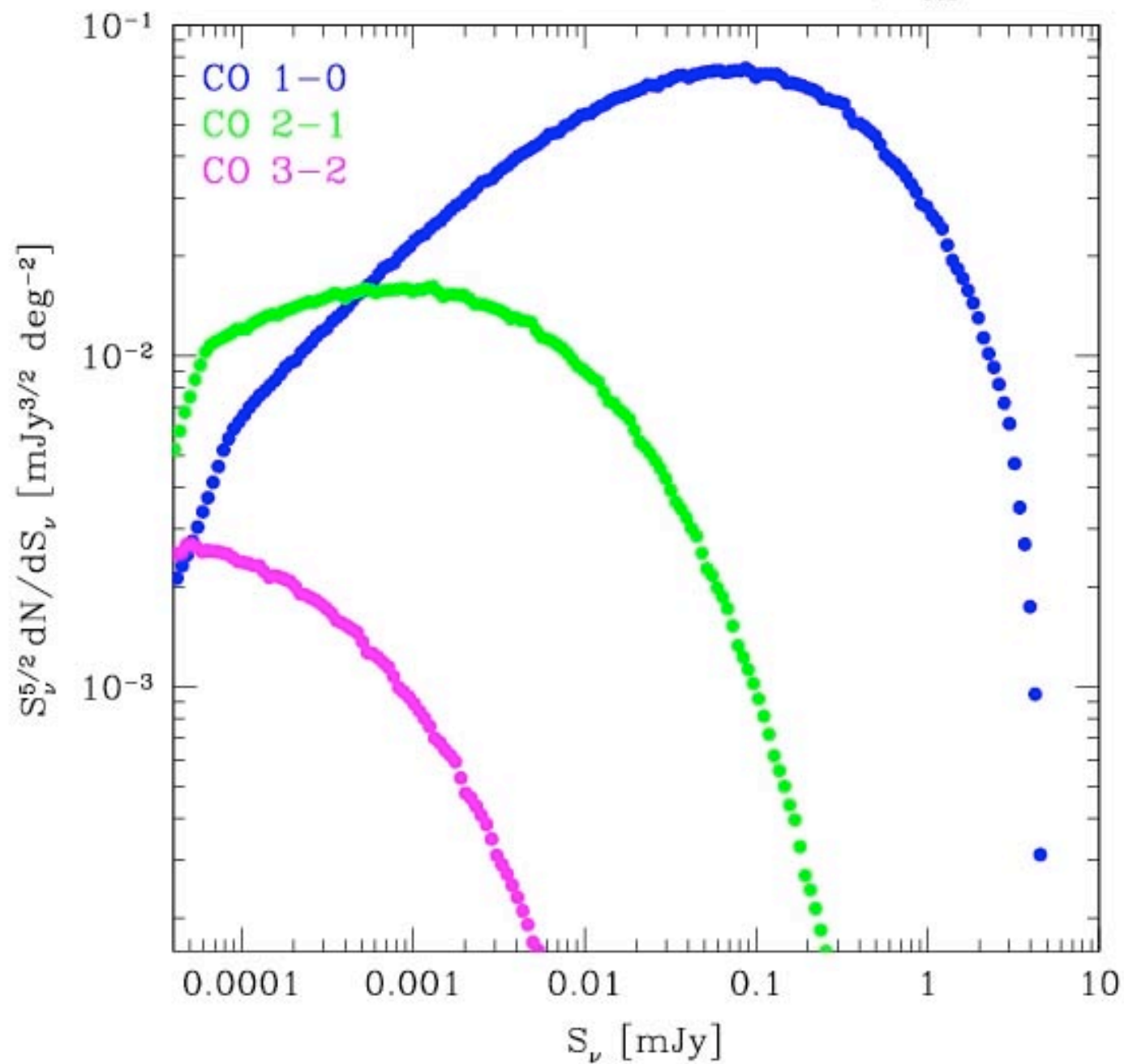
1. It is possible to separate CO contribution to Cl's from dust
2. The way to measure the rate of C and O production in early Universe
3. To measure σ_8
4. To follow star formation rate in the Universe

Righi, M.; Hernández-Monteagudo, C., Sunyaev, R. A., 2008
Astronomy and Astrophysics, 489, pp.489-504

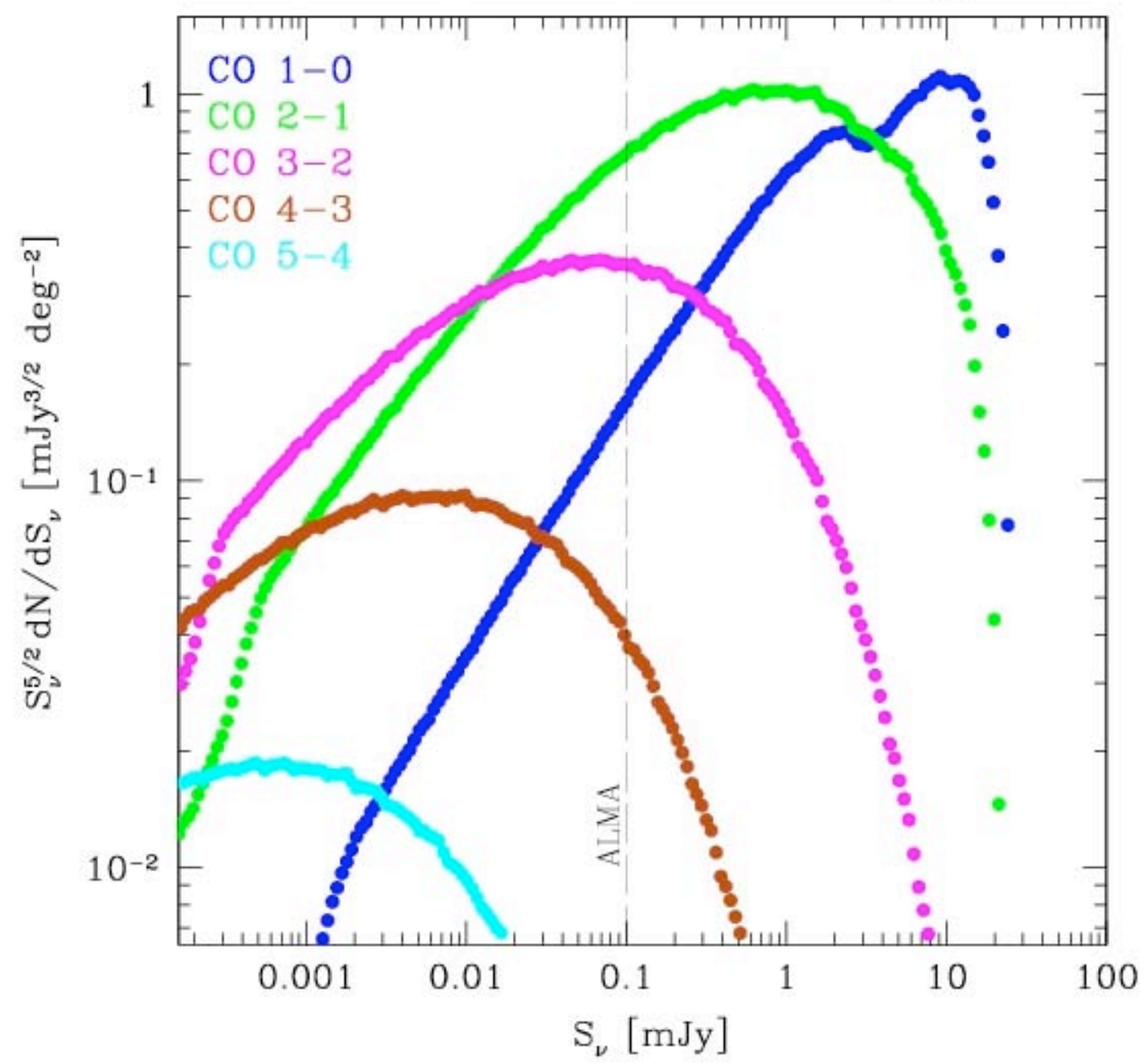
Frequency dependence - $\Delta\nu/\nu_{\text{obs}} = 10^{-3}$



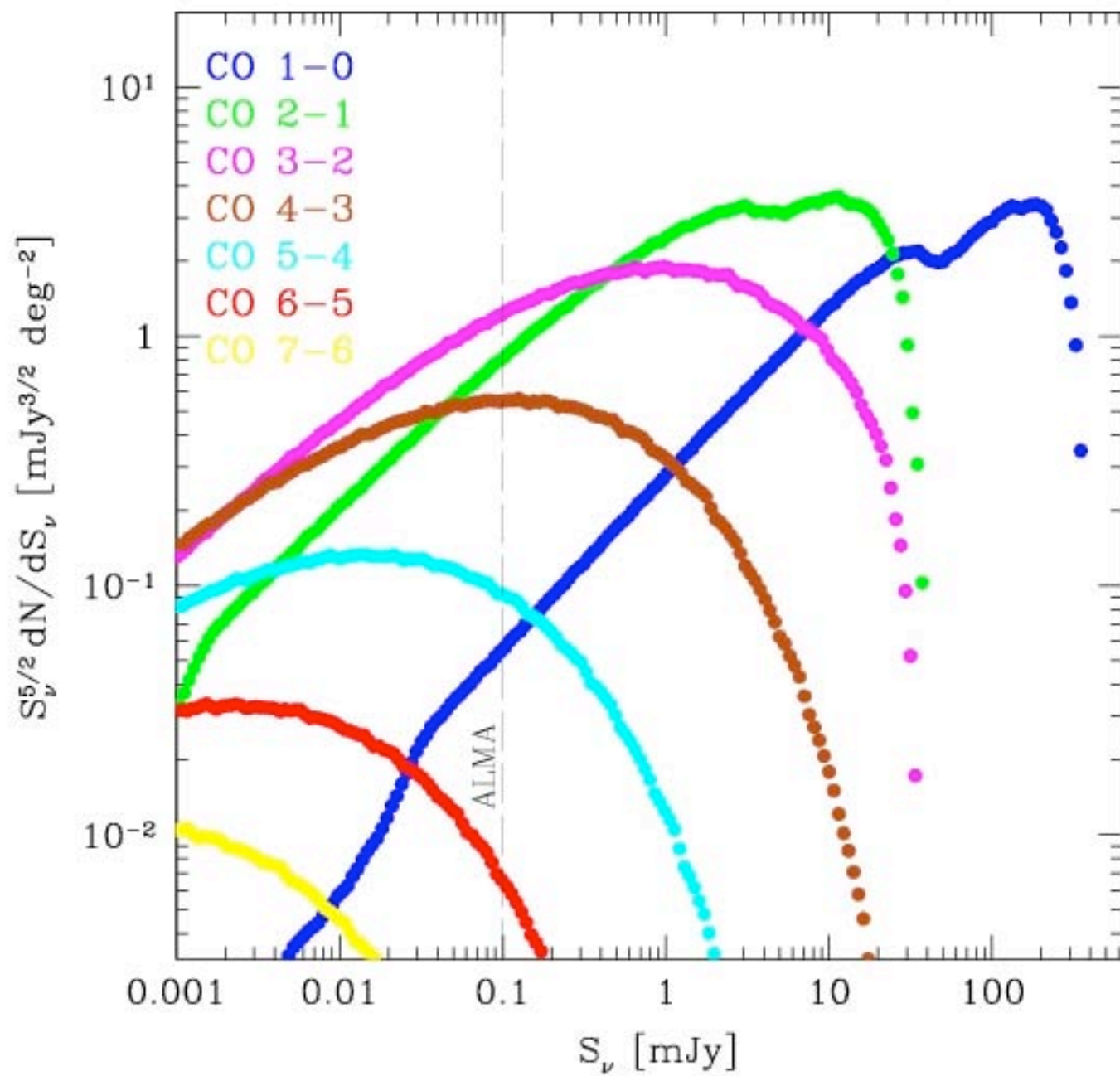
Differential counts - 30 GHz - $\Delta\nu/\nu_{\text{obs}}=10^{-3}$



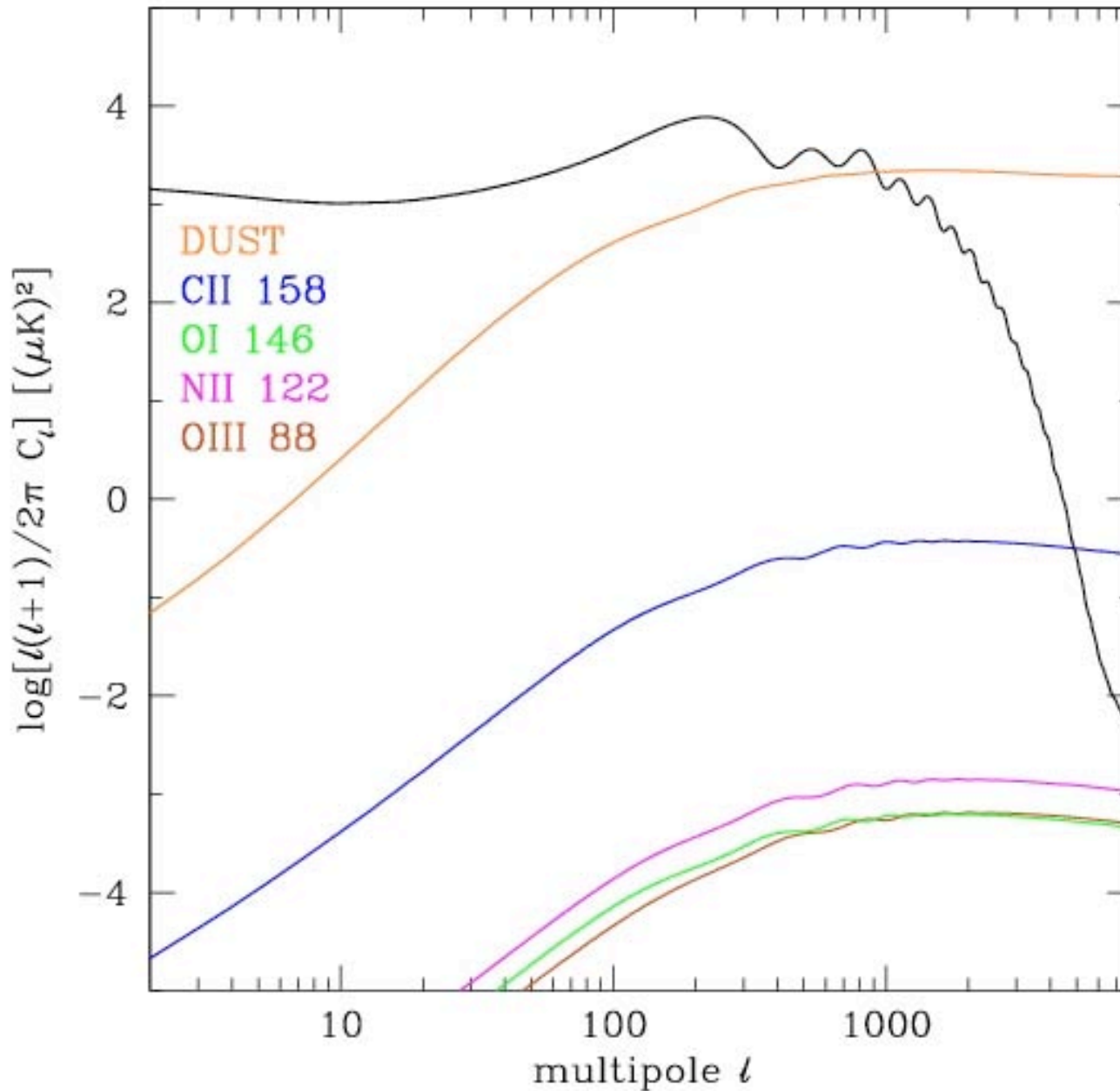
Differential counts - 70 GHz - $\Delta\nu/\nu_{\text{obs}}=10^{-3}$



Differential counts - 100 GHz - $\Delta\nu/\nu_{\text{obs}}=10^{-3}$

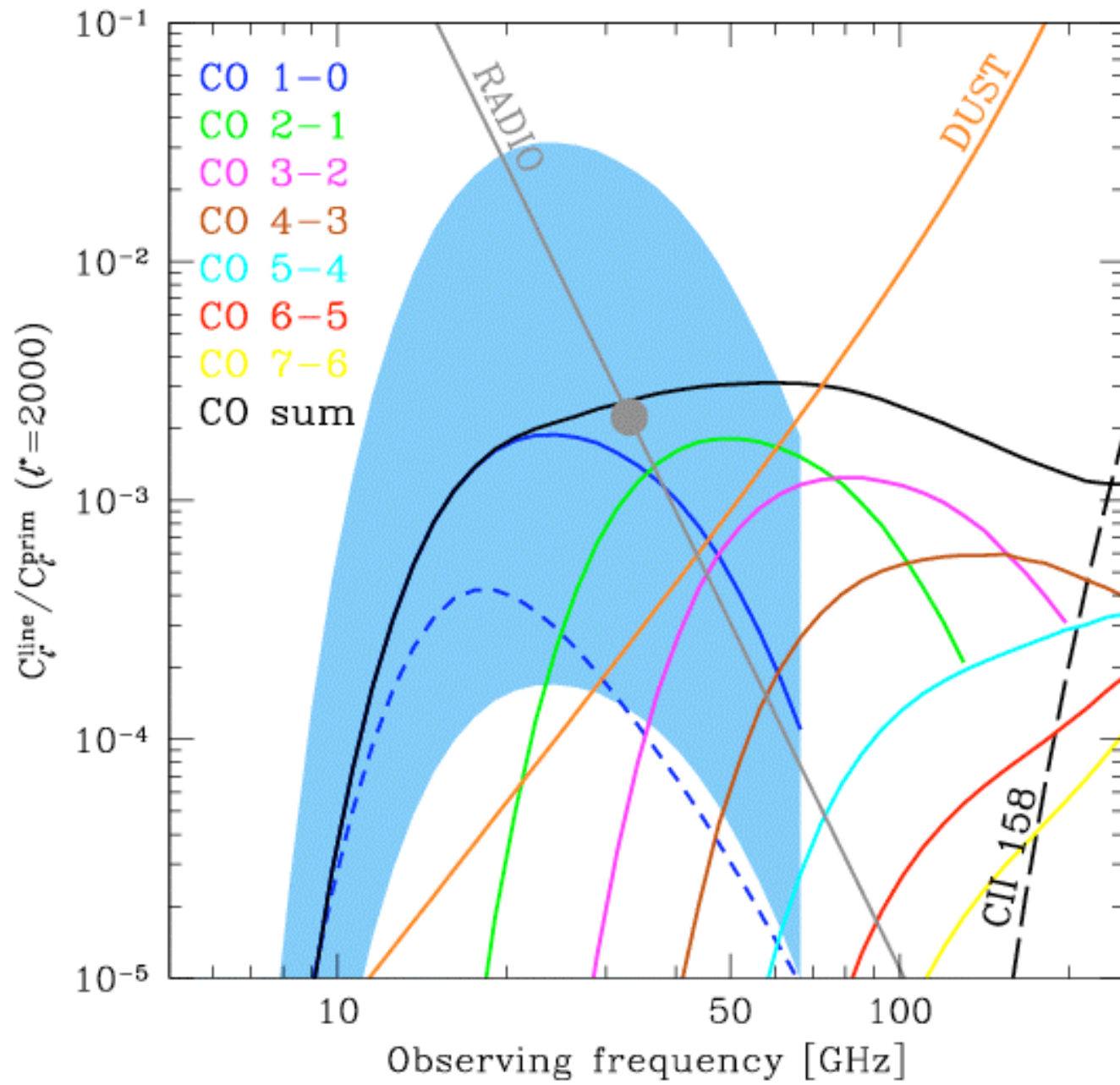


Atomic lines - $\Delta\nu/\nu_{\text{obs}}=0.33$ - 353 GHz

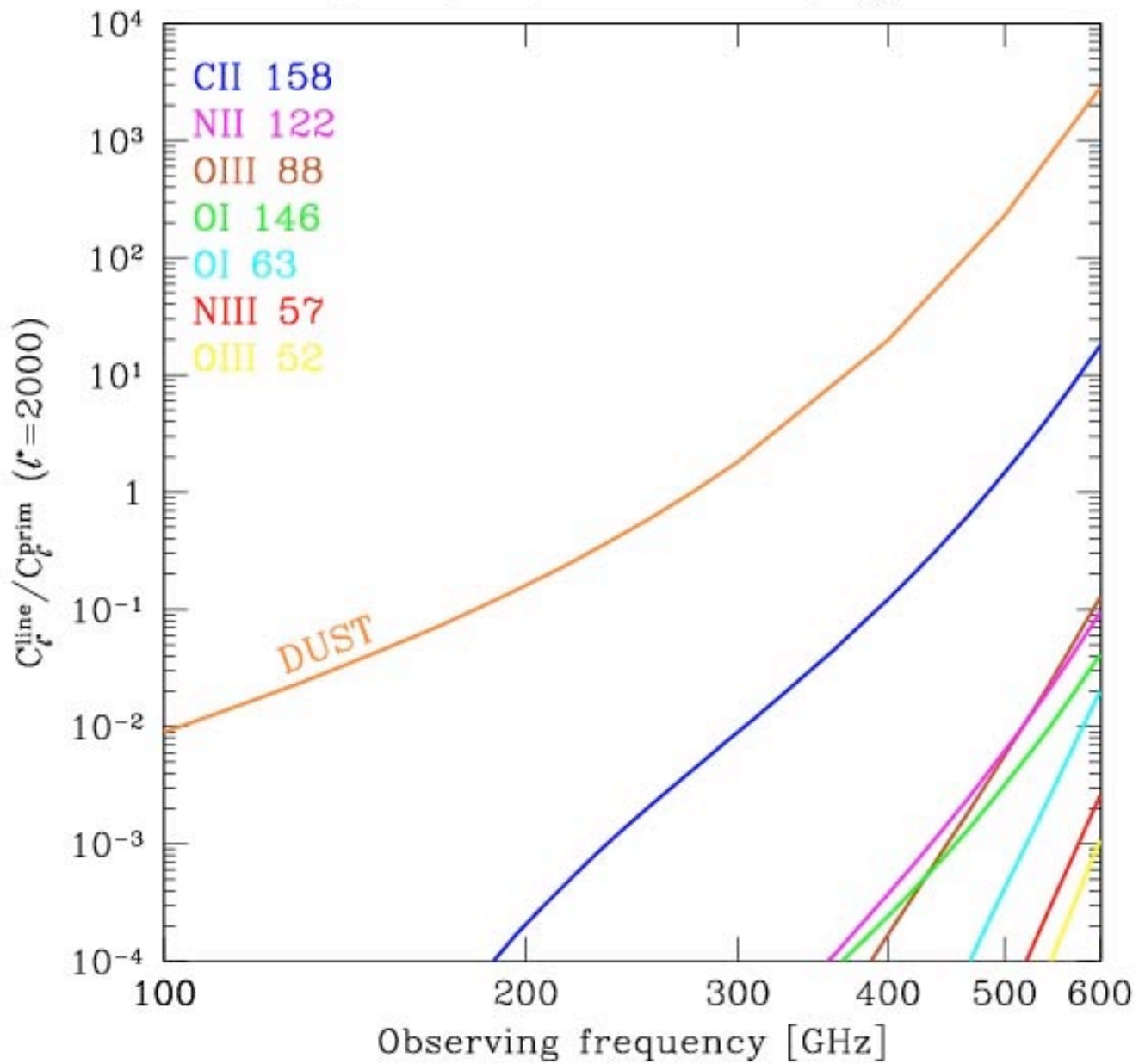


There is a possibility to increase Cl's for the lines 100 times
Increasing spectral resolution

Frequency dependence - $\Delta\nu/\nu_{\text{obs}} = 10^{-3}$



Frequency dependence - $\Delta\nu/\nu_{\text{obs}} = 10^{-3}$



“The EPS theory represents the only fully analytic model of the hierarchical growth of structure. While its derivation requires making several gross approximations and assumptions (Bond et al. 1991; Lacey & Cole 1993), it is remarkable that it captures well the qualitative dependences of progenitor mass distributions on redshift and final halo mass and of final halo mass distributions on initial progenitor mass and redshift. However, its accuracy is not sufficient for the present era of precision cosmology. For example, at high redshift, $z = 4$, it can underestimate the typical progenitor mass by factors of 3 or 4, or equivalently the abundance of the most massive progenitors by factors of a few.”

Cole et al. 2008

POWER SPECTRUM OF THE FLUCTUTATIONS

Two terms: **POISSON + CORRELATION**

$$\langle \Delta I_\nu(\vec{n}_1) \Delta I_\nu(\vec{n}_2) \rangle = \sum_l \frac{2l+1}{4\pi} (C_l^P + C_l^C) P_l(\vec{n}_1 \cdot \vec{n}_2)$$

Generated by the
brightest and rare
sources

Generated by the most
abundant and less bright
sources: probes the
clustering

CORRELATION TERM

$$C_l^c = \frac{2}{\pi} \int k^2 dk P(k) |\Delta_l(k)|^2 \rightarrow \text{Matter power spectrum}$$

Transfer function: contains l dependence and line profile φ

$$\Delta_l(k) = \int dr j_l(kr) \varphi(r) [S(r) \delta_k]$$

Source function: contains information about the population

$$S(r) = \int dL_{v(1+z)} \frac{dn}{dL_{v(1+z)}} \frac{L_{v(1+z)}}{4\pi} cH^{-1}(z) b(M, z)$$

Bias factor => clustering

POISSON TERM

$$C_l^P = \left[r^2 \frac{\Delta \nu}{\nu_{obs}} \right]^{-1} \times \int dL_{\nu(1+z)} cH^{-1}(z) \left(\frac{L_{\nu(1+z)}}{4\pi} \right)^2 \frac{d\bar{n}}{dL_{\nu(1+z)}}$$

Linear dependence on
the spectral resolution

Proportional to the
number of sources

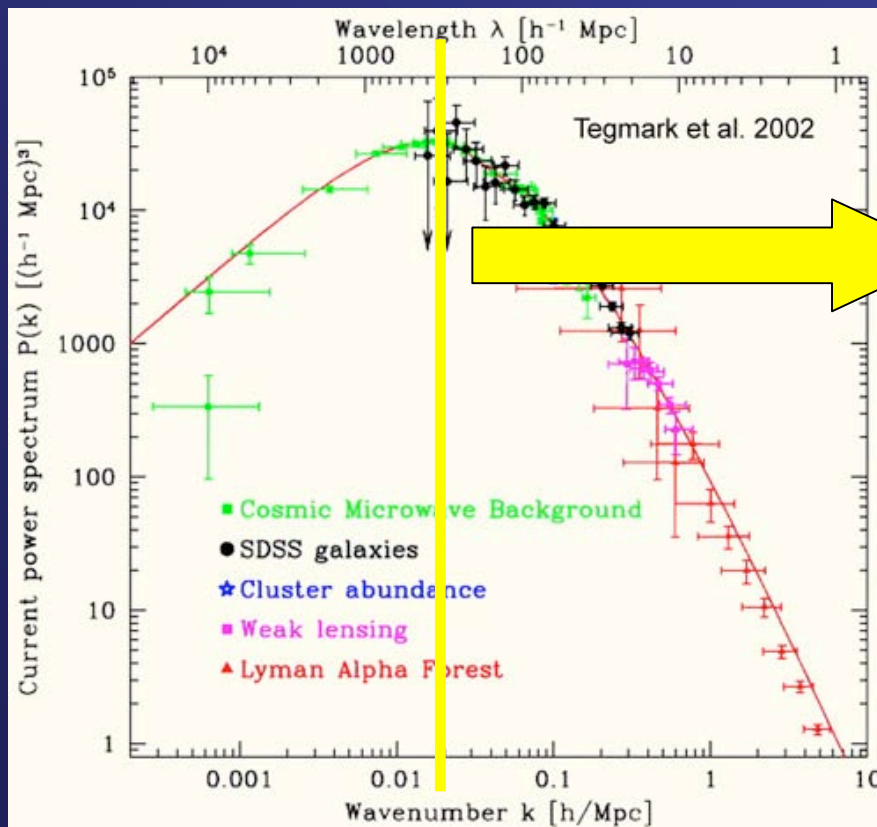
$$\frac{C_l^P}{C_l^C} \propto \frac{nL^2}{(nL)^2} = \frac{1}{n}$$

Poisson power is dominated by rarest (brightest) sources!

The fluctuations below the scale probed by the given $\Delta v/v_{\text{obs}}$ (along the line of sight) are “smeared out”.

$$C_l^C \propto \int_0^{k_z^{\text{MAX}}} \frac{dk_z}{2\pi} P(k_{\perp}, k_z) |W(k_z)|^2$$

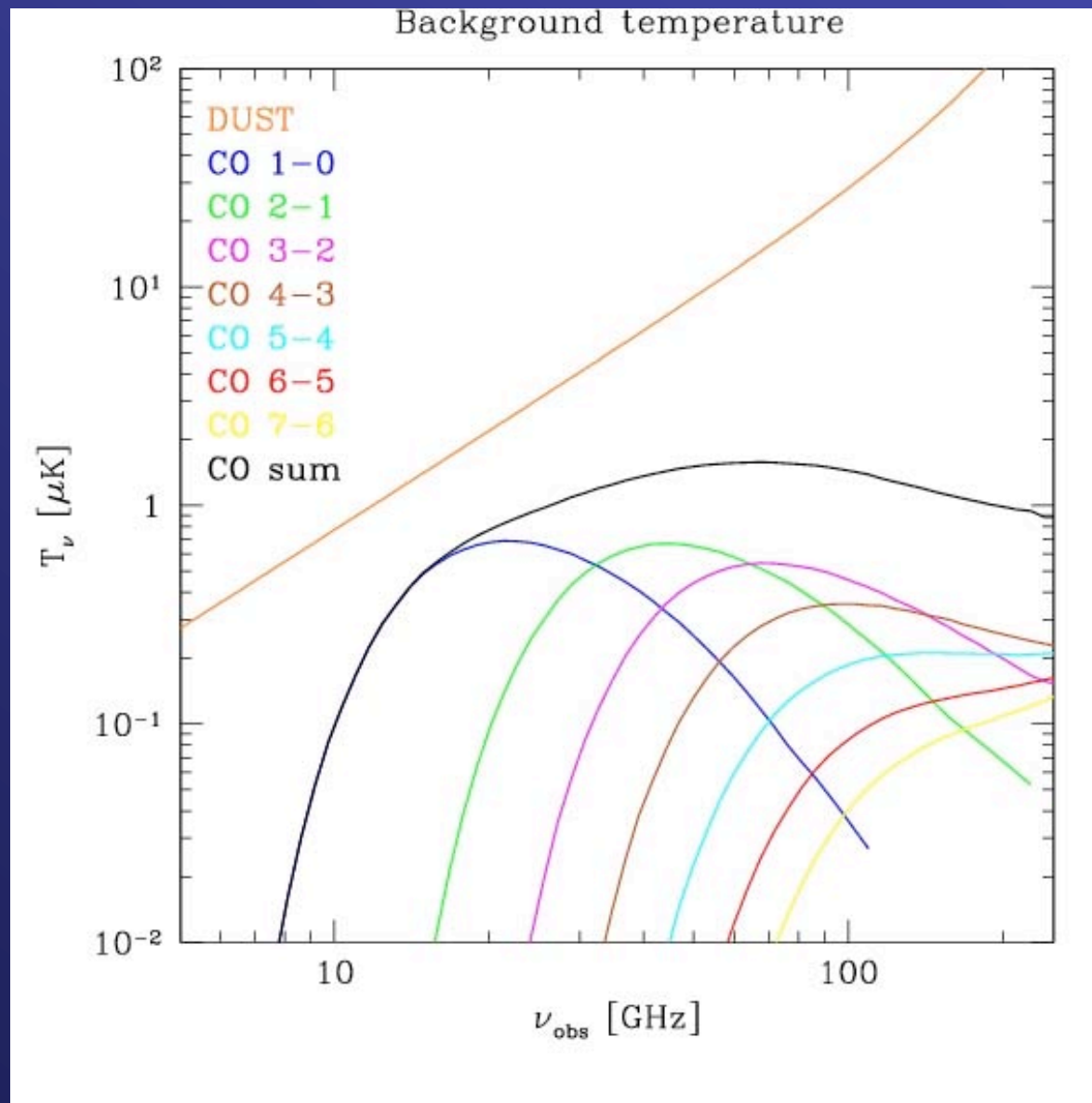
The effect of the window function W is to change the upper limit of integration to a **maximum k** (minimum scale).



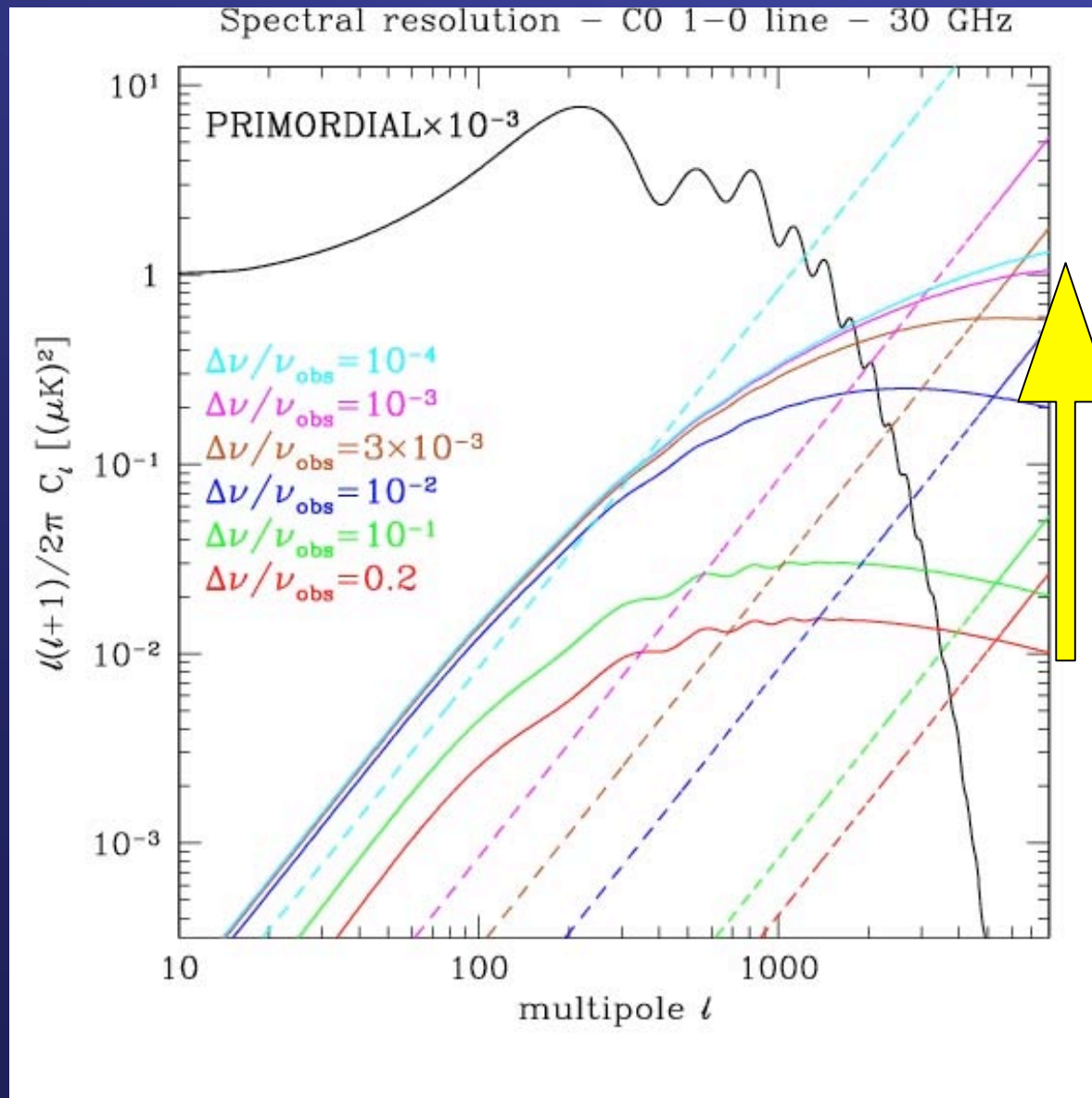
When k_{MAX} becomes larger than the than $\sim 0.01 \text{ h/Mpc}$ (peak of power spectrum) there is no gain in further improving the spectral resolution.

Plateau is reached!

BACKGROUND INTENSITY



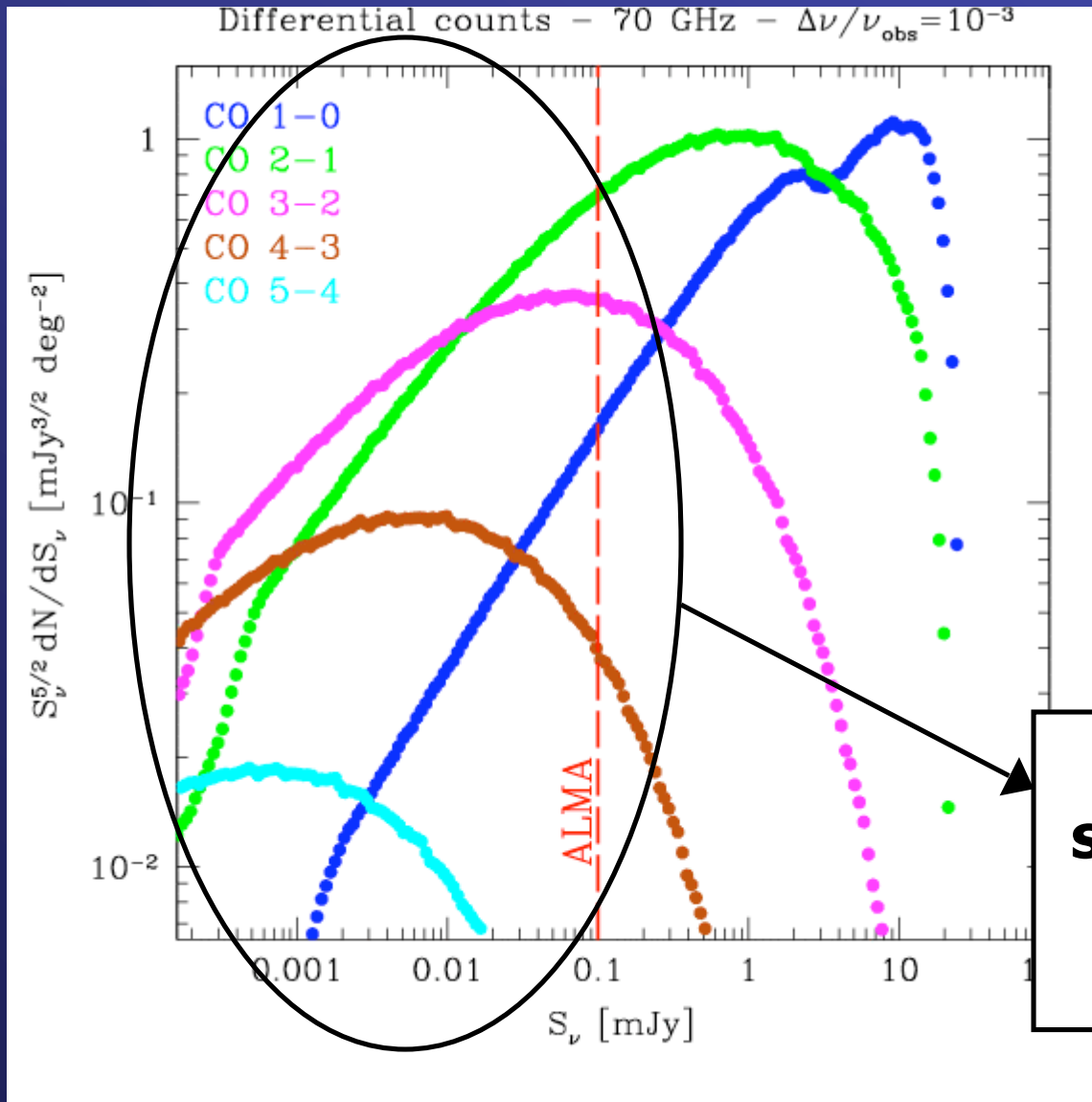
SPECTRAL RESOLUTION



The amplitude of the correlation term grows of ~ 2 orders of magnitude if the spectral resolution is improved down to $\Delta\nu/\nu_{\text{obs}} = 10^{-3}$

Primordial fluctuations and other continuum foreground DO NOT DEPEND ON $\Delta\nu/\nu_{\text{obs}}$

SOURCE NUMBER COUNTS



$\nu_{\text{obs}} = 70 \text{ GHz}$

1-0 - $z_{\text{em}} \sim 0.6$

2-1 - $z_{\text{em}} \sim 2.3$

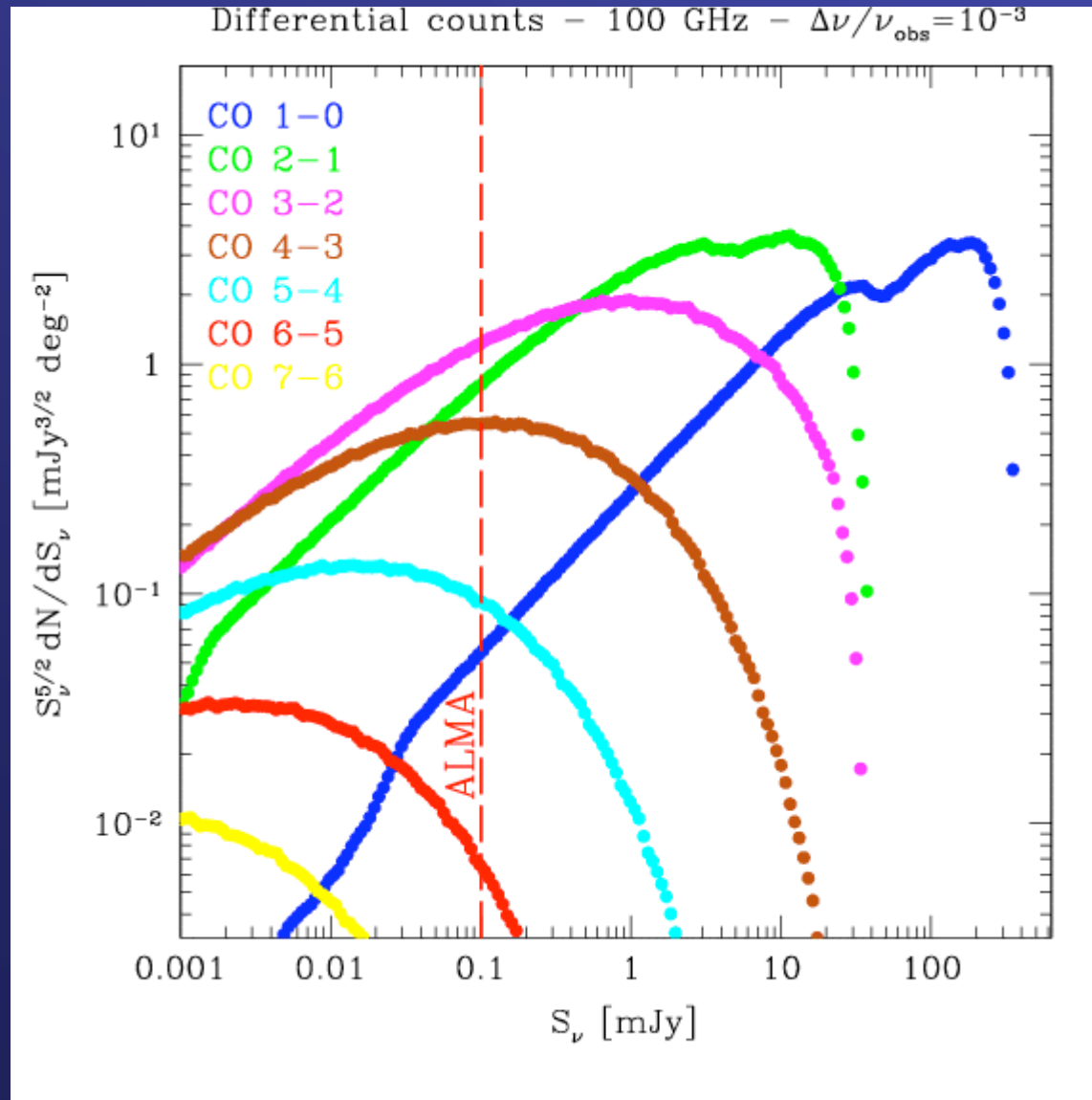
3-2 - $z_{\text{em}} \sim 3.9$

4-3 - $z_{\text{em}} \sim 5.6$

5-4 - $z_{\text{em}} \sim 7.2$

Bulk of the sources is below the sensitivity limits of ALMA

SOURCE NUMBER COUNTS



$\nu_{\text{obs}} = 100 \text{ GHz}$

1-0 - $z_{\text{em}} \sim 0.2$

2-1 - $z_{\text{em}} \sim 1.3$

3-2 - $z_{\text{em}} \sim 2.5$

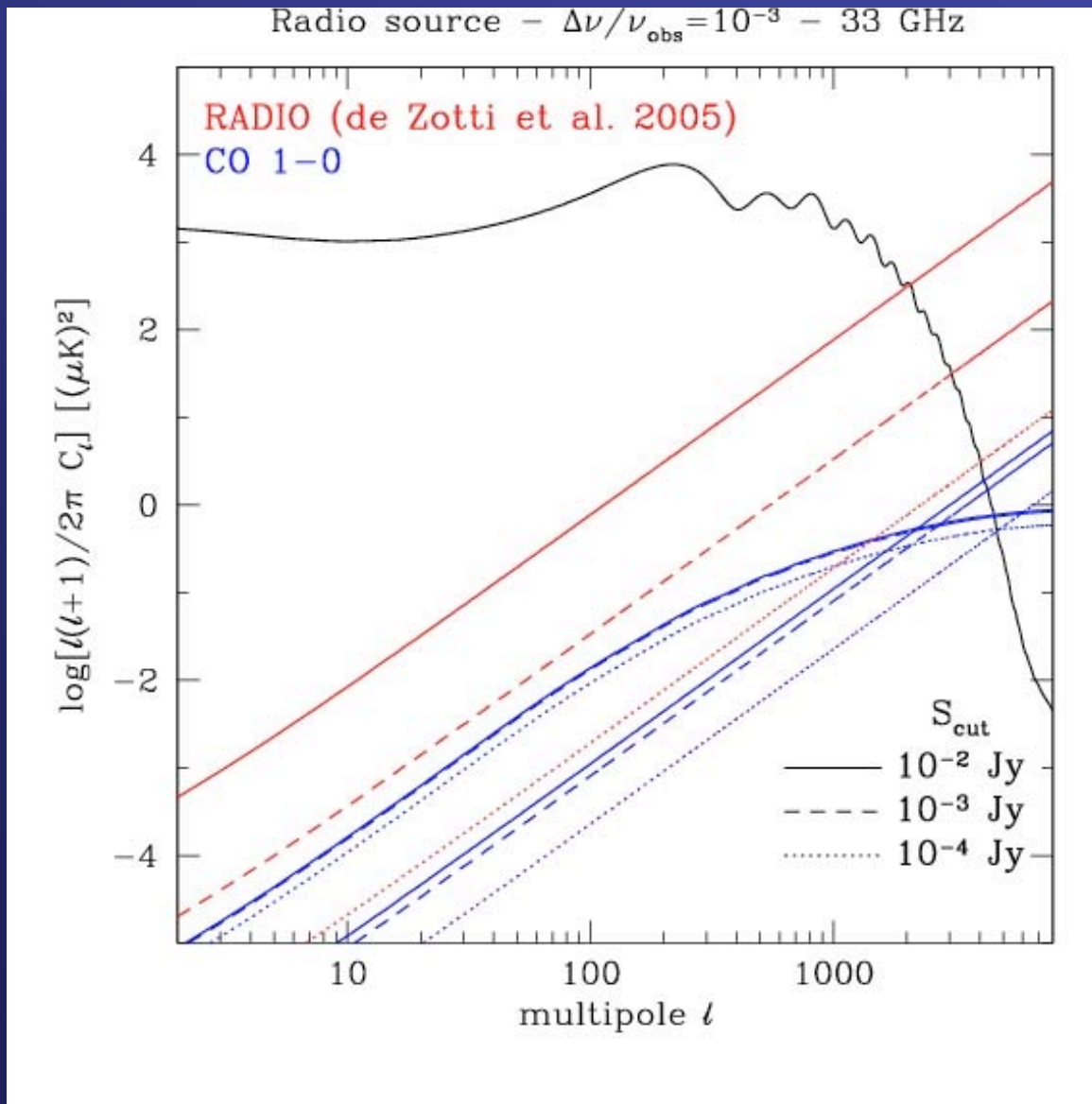
4-3 - $z_{\text{em}} \sim 3.6$

5-4 - $z_{\text{em}} \sim 4.8$

6-5 - $z_{\text{em}} \sim 5.9$

7-6 - $z_{\text{em}} \sim 7.1$

RADIO CONTAMINATION

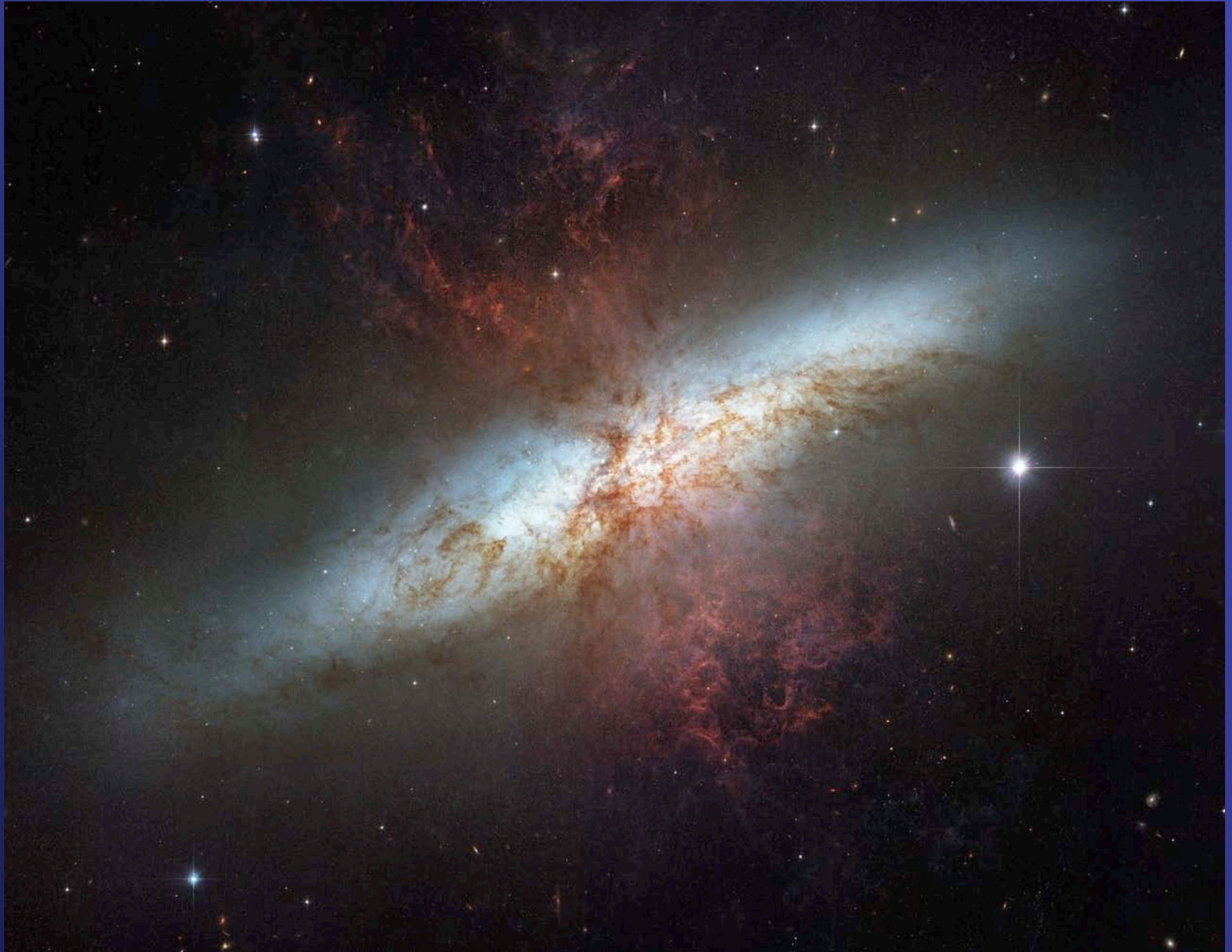


Possibility to decrease the **radio** contamination by **removing** bright sources from the maps.

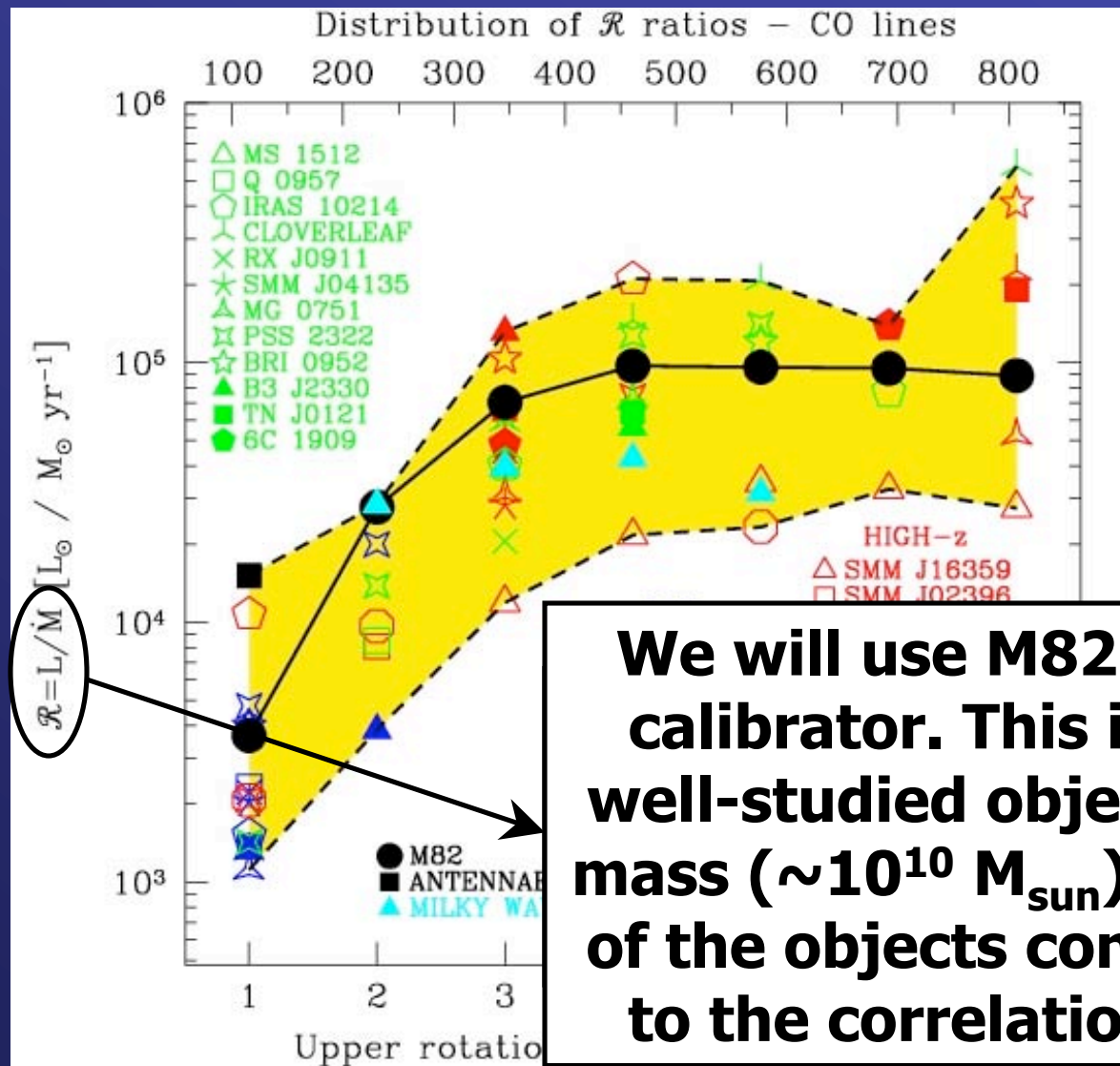
TWO ISSUES

Confusion noise
(20-200 sources/deg² at 10^{-3} - 10^{-4} Jy, should not be a problem)

Clustering properties, not very well known for this population

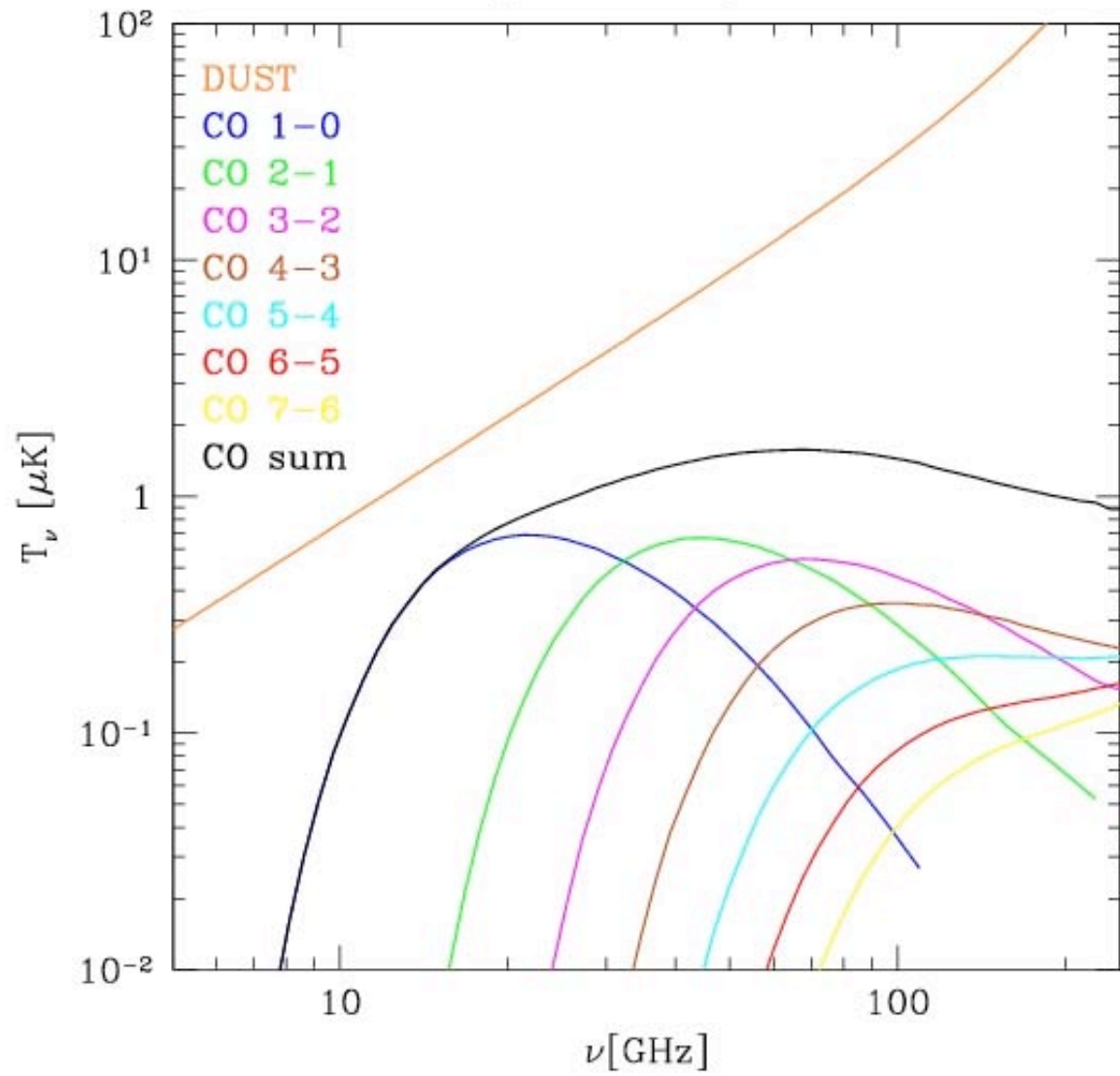


OBSERVATIONAL SAMPLE

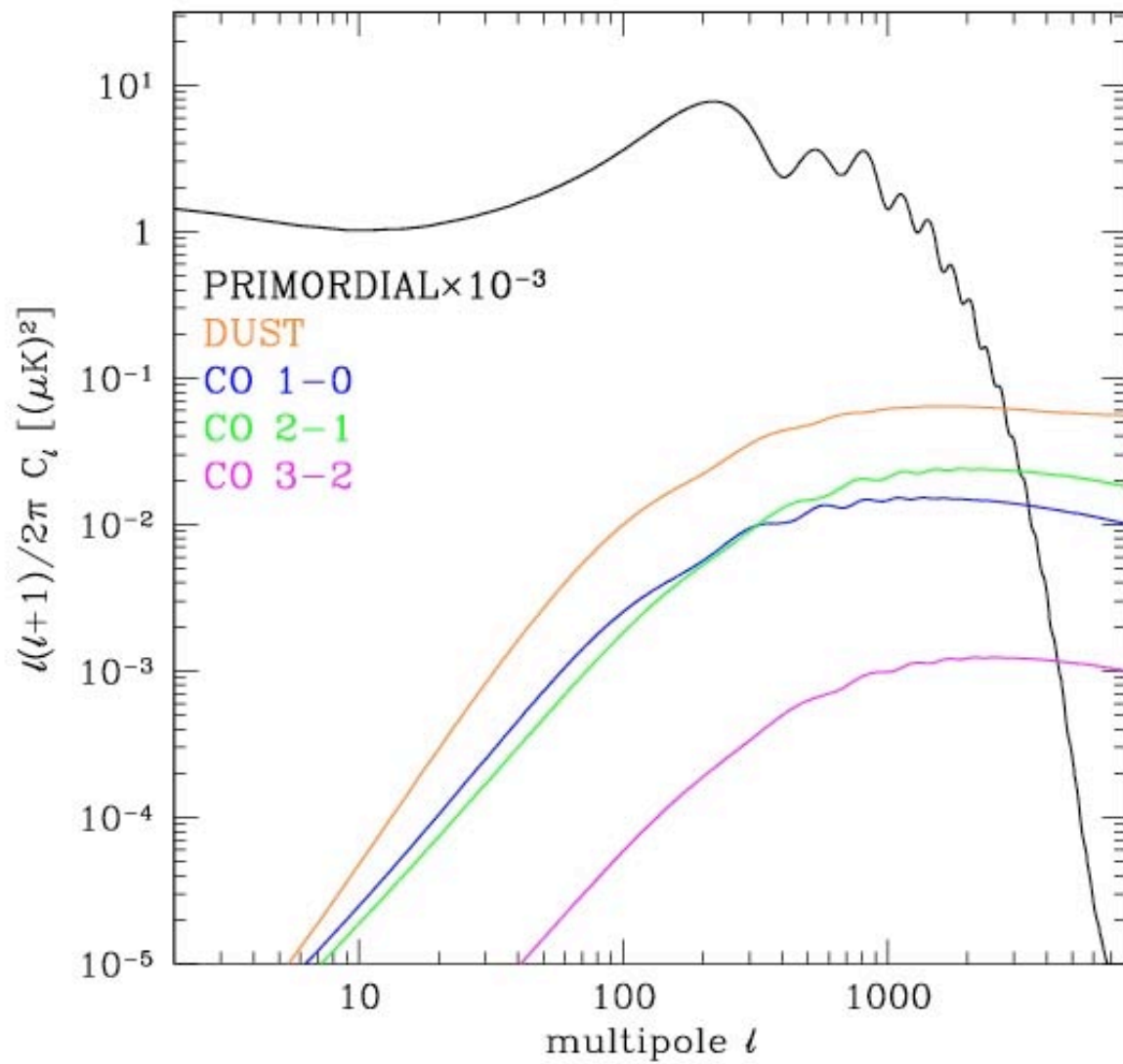


We will use M82 as main calibrator. This is a very well-studied object and its mass ($\sim 10^{10} M_{\text{sun}}$) is typical of the objects contributing to the correlation signal

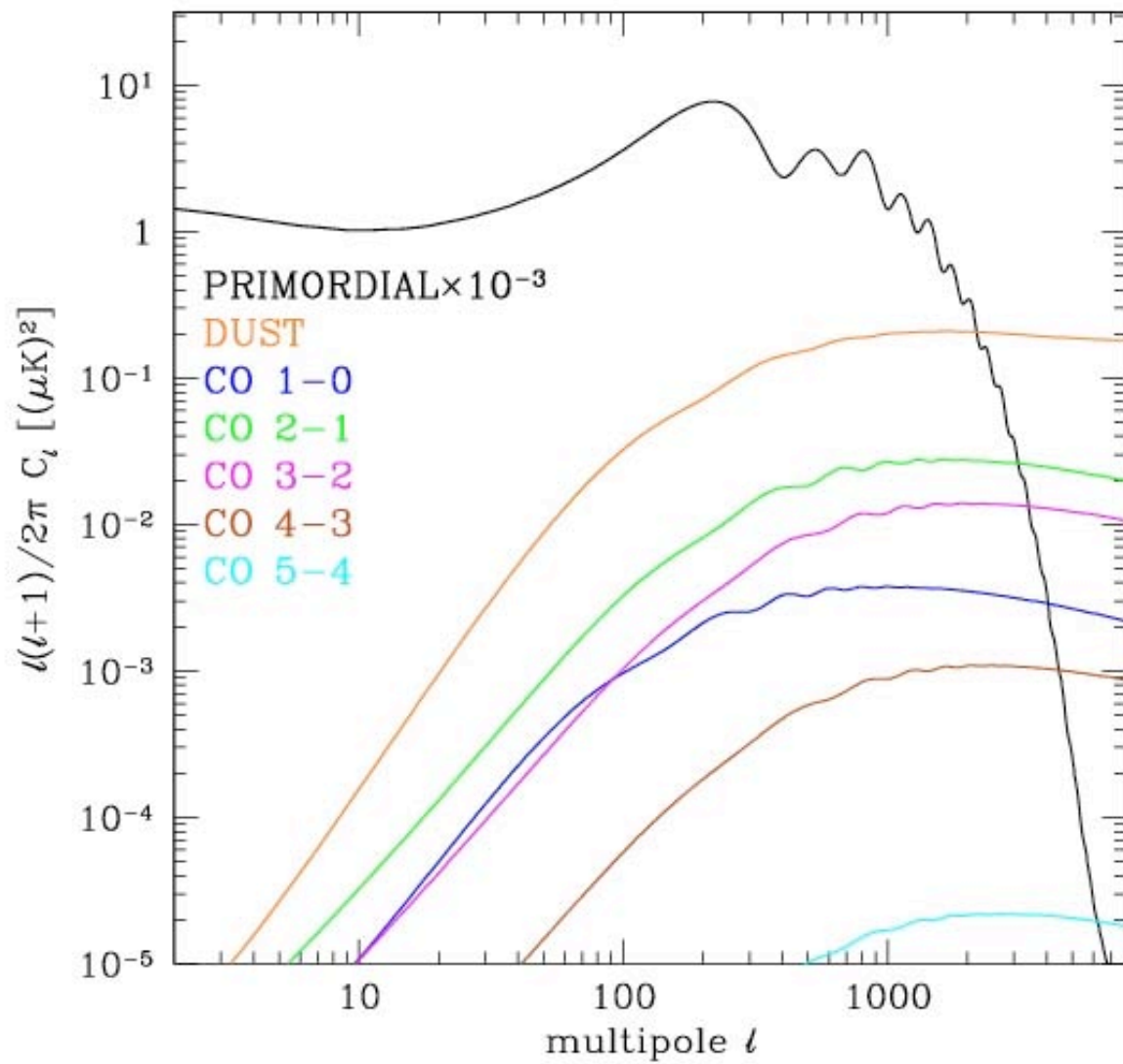
Background temperature



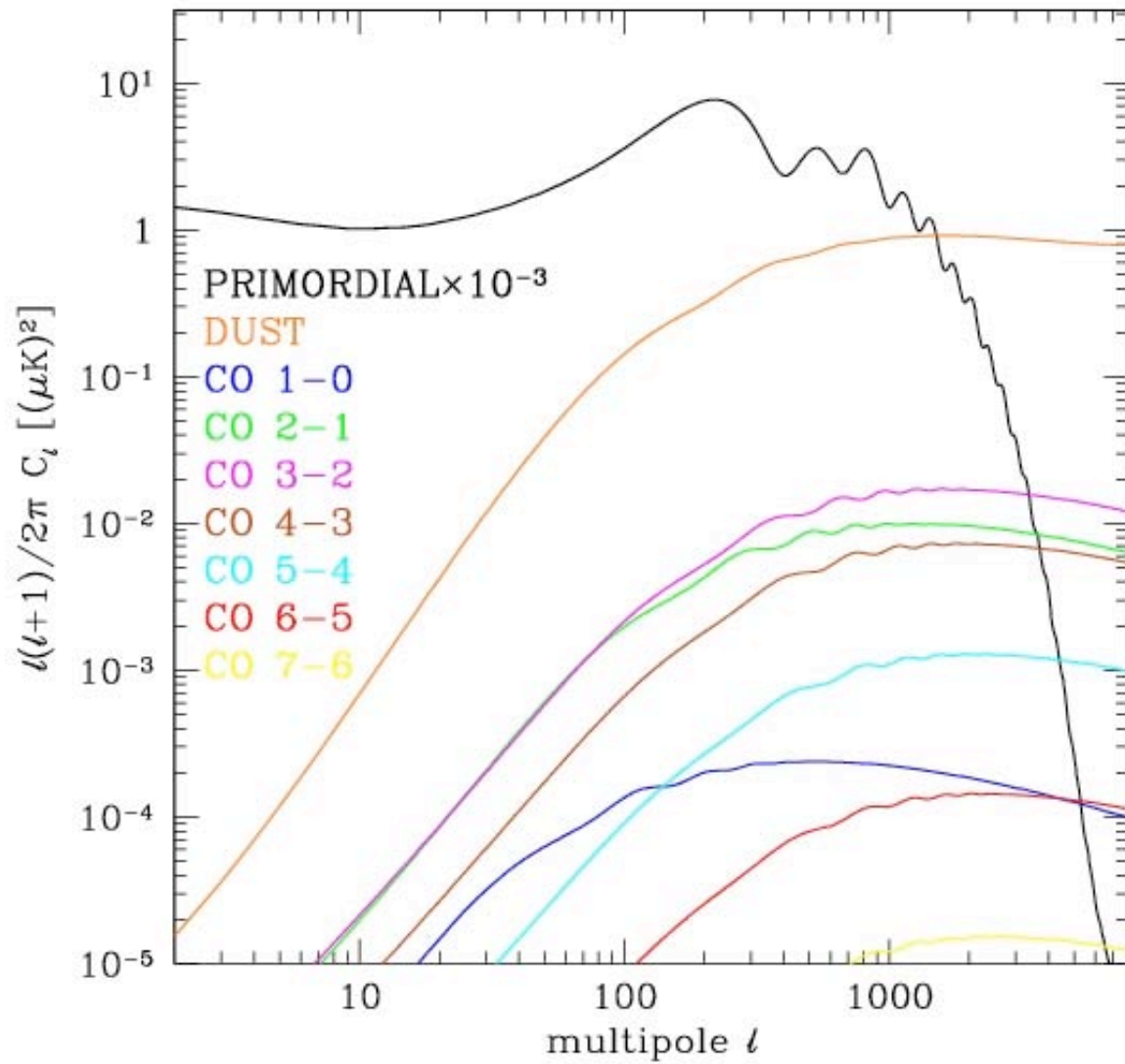
CO lines - $\Delta\nu/\nu_{\text{obs}}=0.2$ - 30 GHz



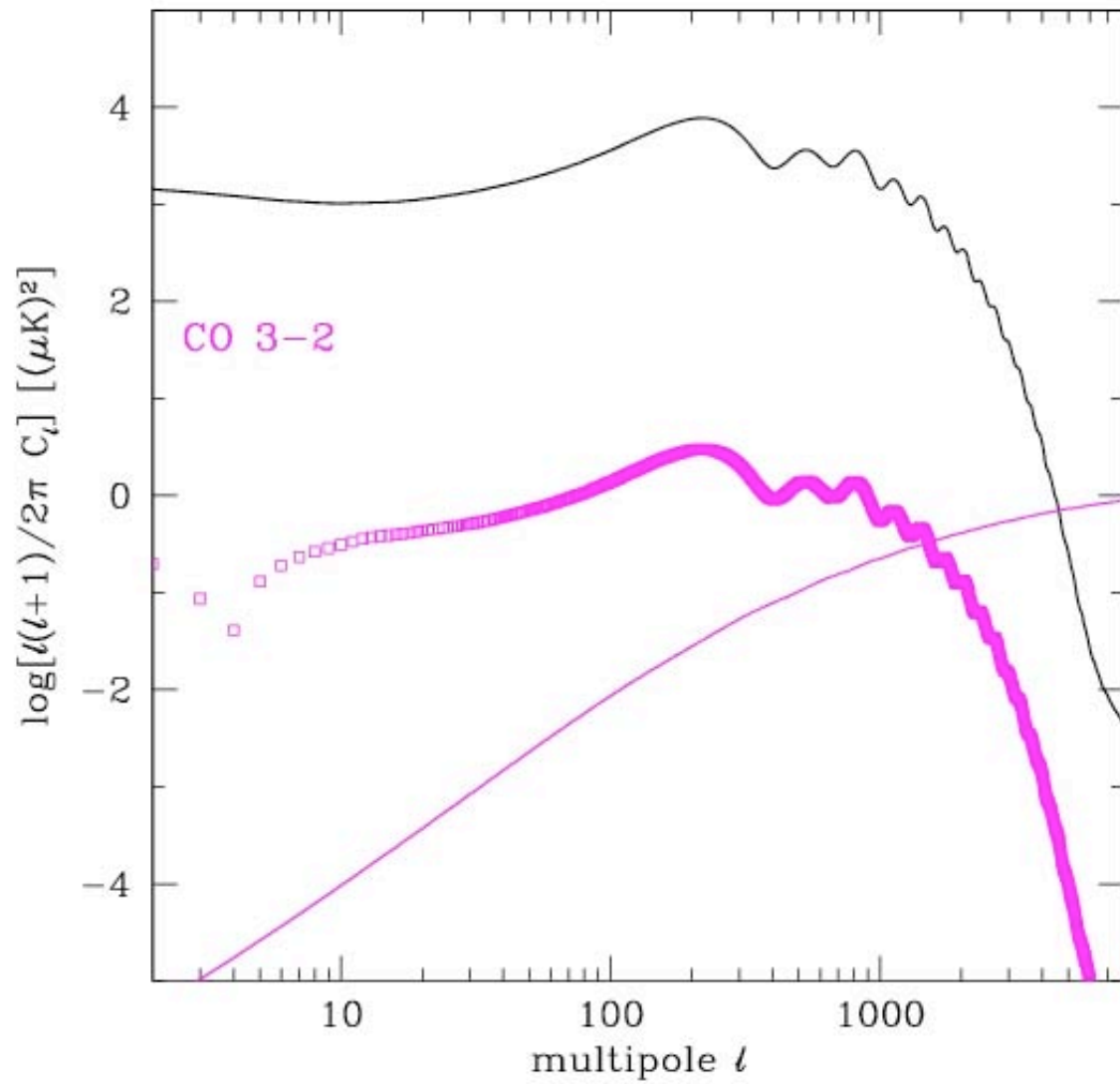
CO lines - $\Delta\nu/\nu_{\text{obs}}=0.2$ - 44 GHz



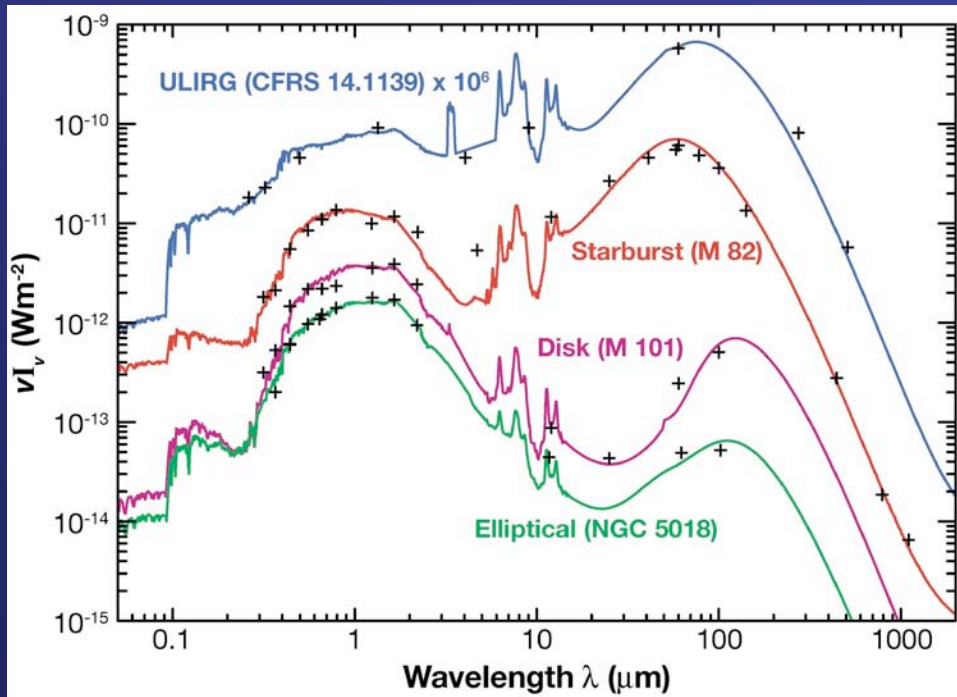
CO lines - $\Delta\nu/\nu_{\text{obs}}=0.2$ - 70 GHz



CO lines scattering - 70 GHz - $\Delta\nu/\nu_{\text{obs}}=10^{-3}$



Dust spectrum



Lagache et al. (2005)

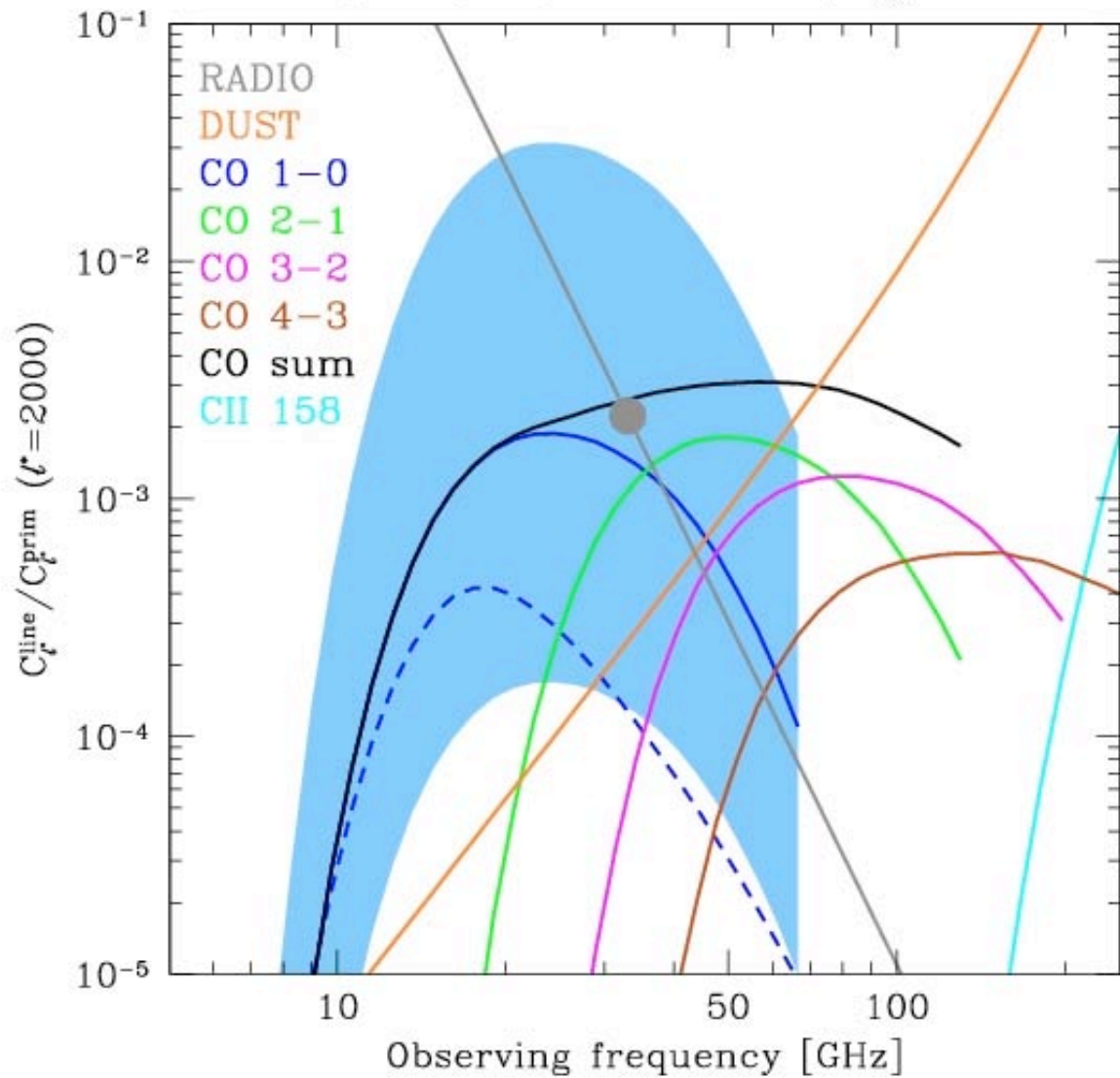
Sub-mm region of the SED of starbursts and ULIRGs is well fitted by a **graybody** (Blain et al. 02, Chapman et al. 05)

$$L_\nu \propto \nu^\beta B_\nu(T_{dust})$$

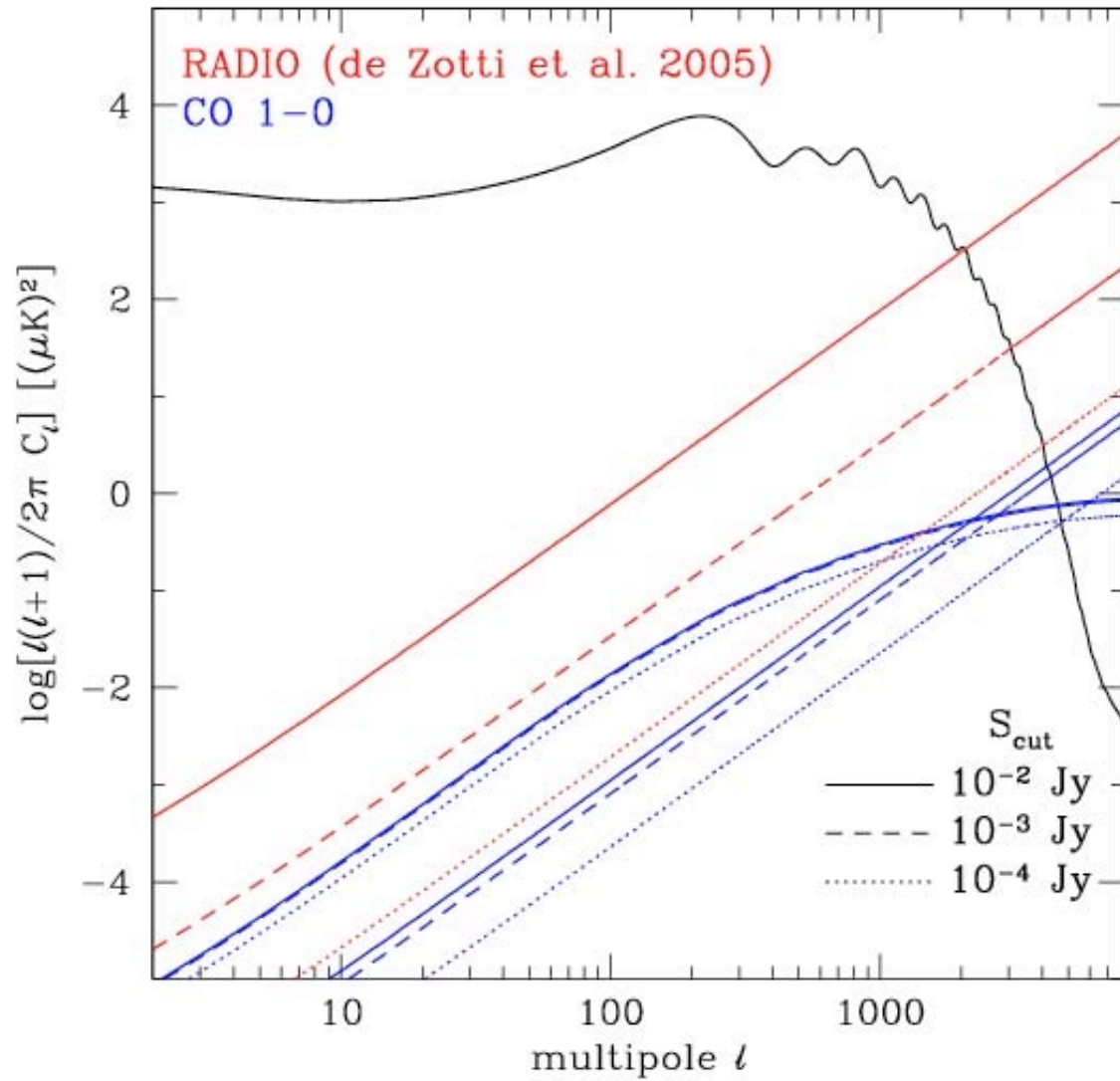
Bulk of the emission is due to low-temperature-dust and peaks around 100 μm

Two free parameters: **dust temperature** T_{dust} and **emissivity index** β , to be calibrated with the observations (**SCUBA**)

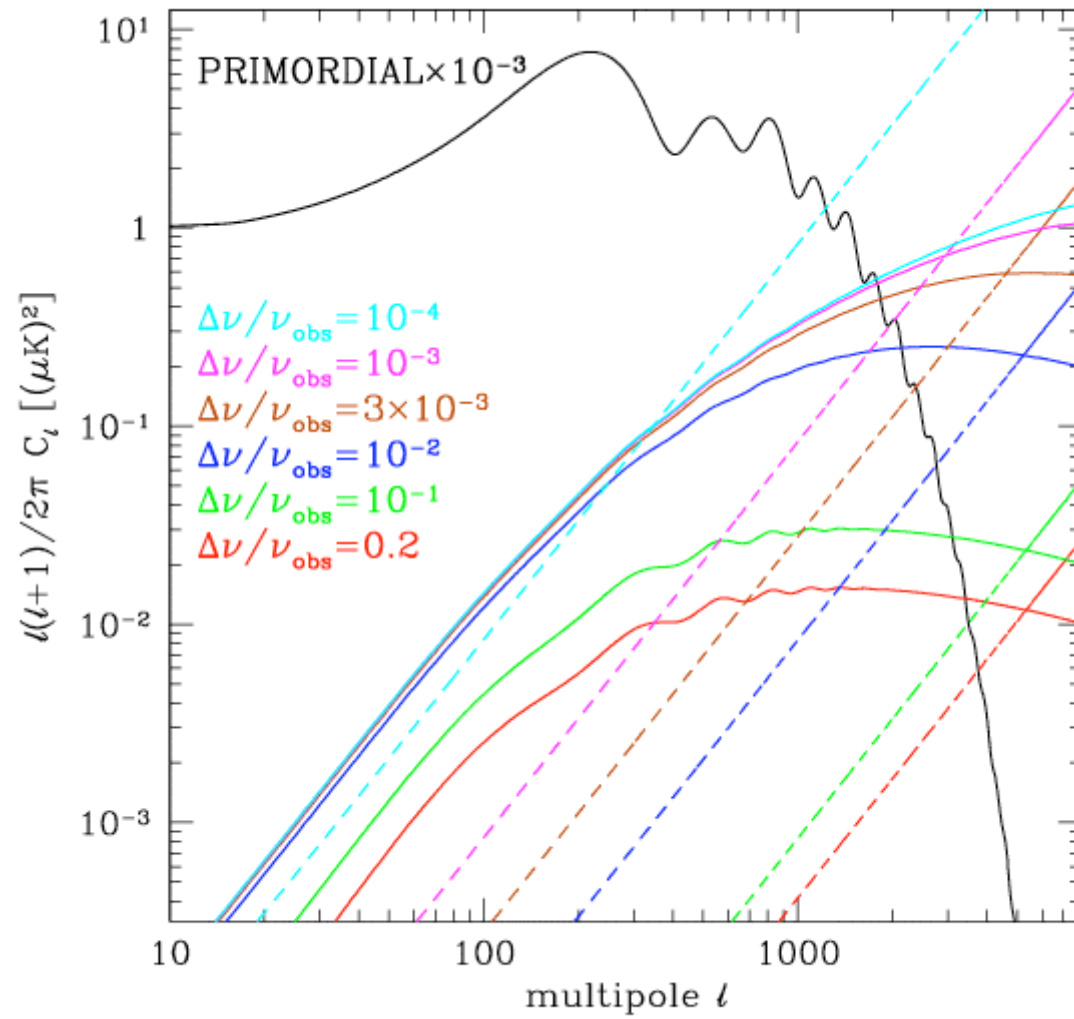
Frequency dependence - $\Delta\nu/\nu_{\text{obs}}=10^{-3}$



Radio source - $\Delta\nu/\nu_{\text{obs}} = 10^{-3} - 33 \text{ GHz}$



Spectral resolution - CO 1-0 line - 30 GHz



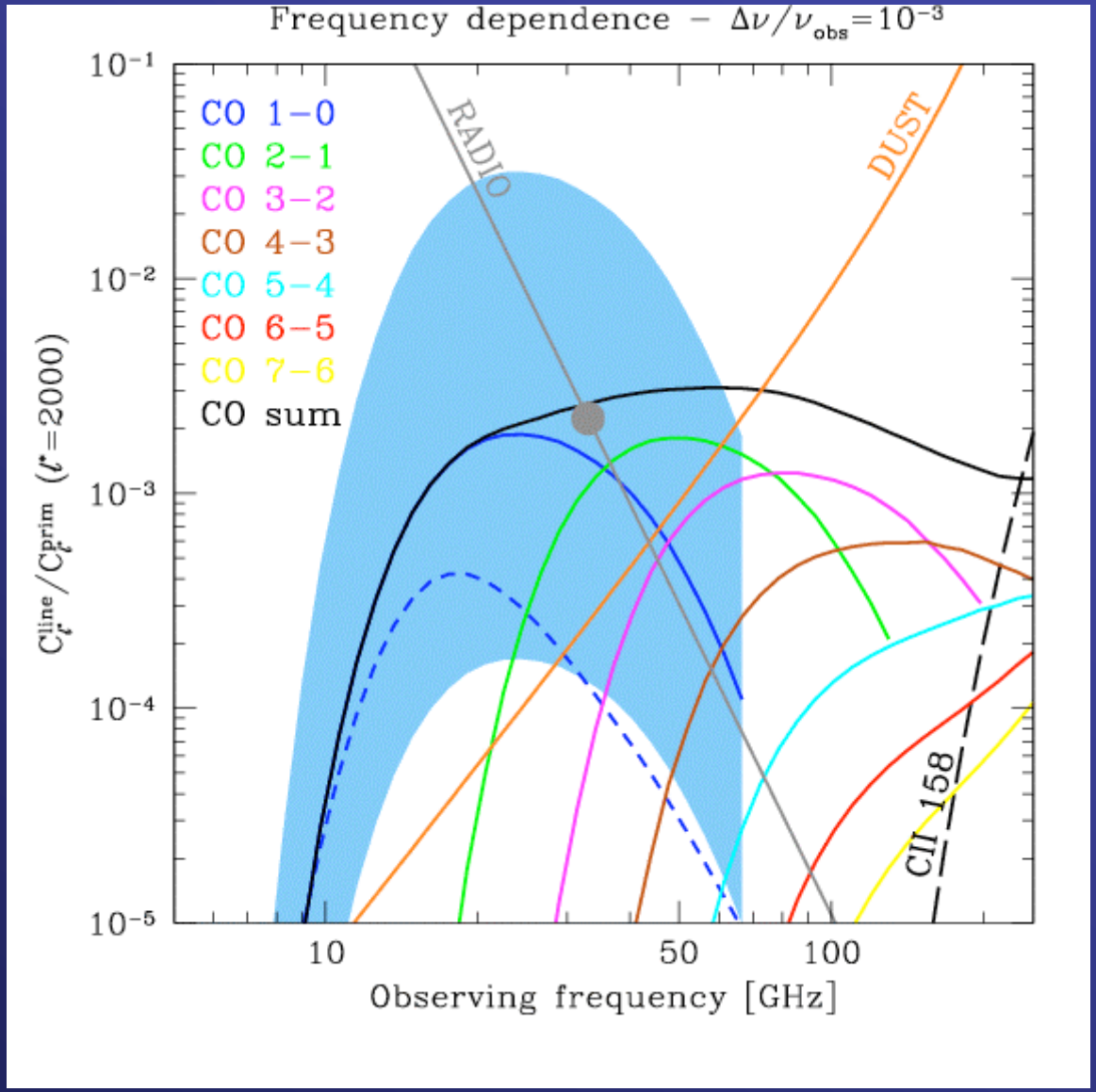


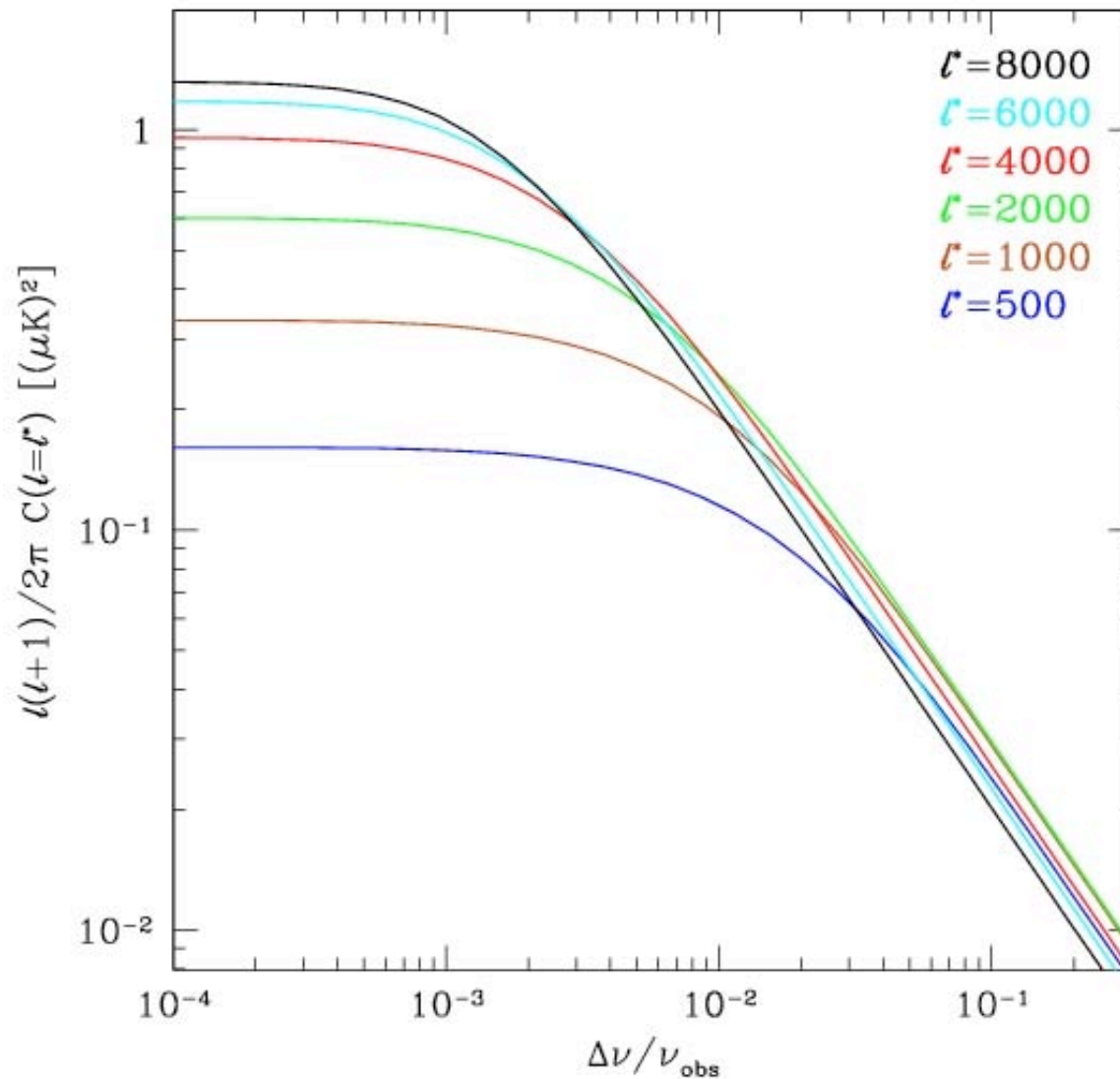
Table 1. Luminosities and \mathcal{R} ratios (in boldface) of the CO lines for the local and low-redshift sample of galaxies. The star formation rates are obtained applying the Kennicutt (1998) relation to the far-infrared (8 – 1000 μm) luminosity given in the *IRAS revised bright galaxies sample* (Sanders et al. 2003). Distances are from the same catalog: conversion from proper to luminosity distance has been computed using the cosmology-corrected redshift in the Nasa Extragalactic Database (NED). For our Galaxy we use the value of star formation rate given by Cox (2000). The third line in M82 shows the ratio between dust and CO lines luminosities, computed assuming $\Delta\nu/\nu \sim 10^{-3}$ and using the spectrum given by Lagache et al. (2005). References: (a) Weiß et al. (2005b), (b) Gao et al. (2001), (c) Wright et al. (1991), (d) Baan et al. (2008).

Object name	Type	d_L	SFR	CO (1-0)	CO (2-1)	CO (3-2)	CO (4-3)	CO (5-4)	CO (6-5)	CO (7-6)
		[Mpc]	[M_\odot/yr]	115.3 GHz	230.5 GHz	345.8 GHz	461.0 GHz	576.3 GHz	691.5 GHz	806.7 GHz
				[L_\odot]	[L_\odot]	[L_\odot]	[L_\odot]	[L_\odot]	[L_\odot]	[L_\odot]
M82	SB	3.63	10.1	3.7×10^4 a	2.8×10^5 a	7.1×10^5 a	9.8×10^5 a	9.7×10^5 a	9.6×10^5 a	9.0×10^5 a
				3.7×10^3	2.8×10^4	7.0×10^4	9.7×10^4	9.6×10^4	9.5×10^4	8.9×10^4
		M82 $L_{\text{dust}}/L_{\text{CO}}$ ratio:			0.005	0.009	0.02	0.05	0.12	0.24
Antennae	SB	21.8	11.8	1.8×10^5 b						
				1.5×10^4						
Milky Way		-	3.0		8.5×10^4 c	1.2×10^5 c	1.3×10^5 c	9.5×10^4 c		
					2.8×10^4	4.0×10^4	4.3×10^4	3.2×10^4		
IRAS 01077-1707		145	72.9	2.9×10^5 d						
				4.0×10^3						
IRAS 01364-1042		207	98.4	2.3×10^5 d						
				2.3×10^3						
IRAS 04454-4838		220	105	1.6×10^5 d						
				1.5×10^3						
IRAS 08520-6850		205	98.4	1.3×10^5 d						
				1.3×10^3						
IRAS 09111-1007		243	171	3.8×10^5 d						
				2.2×10^3						
IRAS 14348-1447	SB	388	341	6.9×10^5 d						
				2.0×10^3						
IRAS 14378-3651		315	242	2.7×10^5 d						
				1.1×10^3						
IRAS 18293-3413		80.6	110	5.3×10^5 d	2.2×10^6 d					
				4.8×10^3	2.0×10^4					
IRAS 19115-2124		216	127	5.1×10^5 d						
				4.0×10^3						
IRAS 20550+1656		136	127	1.7×10^5 d	4.8×10^5 d					
				1.3×10^3	3.8×10^3					
IRAS 22491-1808	SB	350	220	3.0×10^5 d						
				1.4×10^3						

Table 2. Luminosities and \mathcal{R} ratios (in boldface) of the CO lines for the high-redshift sample of galaxies. Luminosities of the lensed sources (marked with \star) are corrected with the magnification factors given in Greve et al. (2005) and Solomon & Vanden Bout (2005). Where more than one observation were available we assumed an average between them. References: (a) Kneib et al. (2004), (b) Weiß et al. (2005a), (c) Sheth et al. (2004), (d) Kneib et al. (2005), (e) Solomon & Vanden Bout (2005), (f) Greve et al. (2005), (g) Hainline et al. (2006), (h) Tacconi et al. (2006), (i) Kovács et al. (2006), (j) Neri et al. (2003), (k) Genzel et al. (2003), (l) Frayer et al. (1999), (m) Downes & Solomon (2003), (n) Takata et al. (2006), (o) Greve et al. (2003), (p) Andreani et al. (2000), (q) Frayer et al. (2008), (r) Baker et al. (2004), (s) Planesas et al. (1999), (t) Brown & Vanden Bout (1991), (u) Solomon et al. (1992b), (v) Solomon et al. (1992a), (w) Downes et al. (1995), (x) Barvainis et al. (1994), (y) Wilner et al. (1995), (z) Barvainis et al. (1997), (aa) Weiß et al. (2003), (ab) Hainline et al. (2004), (ac) Barvainis et al. (2002), (ad) Carilli et al. (2002), (ae) Cox et al. (2002), (af) Guilleaume et al. (1999), (ag) De Breuck et al. (2003a), (ah) De Breuck et al. (2003b), (ai) Papadopoulos et al. (2000).

Object name	Type	z	SFR [M_{\odot}/yr]	CO (1-0)	CO (2-1)	CO (3-2)	CO (4-3)	CO (5-4)	CO (6-5)	CO (7-6)
				115.3 GHz [L_{\odot}]	230.5 GHz [L_{\odot}]	345.8 GHz [L_{\odot}]	461.0 GHz [L_{\odot}]	576.3 GHz [L_{\odot}]	691.5 GHz [L_{\odot}]	806.7 GHz [L_{\odot}]
SMM J16359 \star	SB	2.52	500 a			6.0×10^6 b-d	1.1×10^7 b	1.7×10^7 b	1.6×10^7 b	1.4×10^7 b
SMM J02396 \star	AGN	1.06	975 e		7.9×10^6 f 8.1×10^6		2.2×10^7	3.4×10^7	3.2×10^7	2.8×10^7
SMM J13120	AGN	3.41	810 g	1.0×10^7 g 1.2×10^7			1.7×10^6 f 2.1×10^6			
SMM J16366	SB	2.45	1455 i			7.7×10^7 fh 5.3×10^7				3.3×10^8 h 2.3×10^8
SMM J16371	SB+AGN	2.38	877 i			4.0×10^7 f 4.6×10^7				
SMM J22174	SB	3.10	1800 e			5.1×10^7 f 2.8×10^7				
SMM J04431 \star	SB+AGN	2.51	450 e			1.4×10^7 hj 3.1×10^7				2.4×10^7 h 5.3×10^7
SMM J09431 \star	SB+AGN	3.35	1200 g	$< 2.5 \times 10^6$ g $< 2.1 \times 10^7$			8.9×10^7 hj 7.4×10^7			
SMM J16368	SB+AGN	2.38	897 i			9.3×10^7 hj 1.0×10^8				3.7×10^8 hj 4.1×10^8
SMM J02399 \star	SB+AGN	2.80	500 k			6.6×10^7 k,l 1.3×10^8				
SMM J14011 \star	SB	2.56	360 e			2.4×10^7 lm 6.7×10^7				6.9×10^7 m 1.9×10^8
SMM J123549	SB+AGN	2.20	1163 n			5.6×10^7 h 4.8×10^7			1.6×10^8 h 1.4×10^8	
ERO J16450	SB	1.44	1539 o	3.2×10^6 o 2.1×10^7	1.5×10^7 d 9.7×10^6			3.6×10^7 p 2.3×10^7		
GOODS J123634	SB	1.22	950 q		2.6×10^7 q 2.8×10^7					
MS 1512 \star	LBG	2.73	15 e			5.9×10^6 r 4.0×10^6				
Q 0957 \star	QSO	1.41	900 e		7.6×10^6 s 8.5×10^6					
IRAS F10214 \star	QSO	2.29	540 e			2.2×10^7 t-w 4.0×10^7			4.2×10^7 v 7.7×10^7	
CLOVERLEAF \star	QSO	2.56	810 e			4.8×10^7 x-za 5.9×10^7	1.2×10^8 z 1.5×10^8	1.7×10^8 z 2.1×10^8		4.6×10^8 z 5.7×10^8
RX J0911 \star	QSO	2.80	345 e			7.1×10^6 ab 2.1×10^7				
SMM J04135 \star	QSO	2.84	3600 e			2.3×10^8 ab 6.4×10^8				
MG 0751 \star	QSO	3.20	435 e				3.2×10^7 ac 7.2×10^7			
PSS J2322 \star	QSO	4.11	1800 e	2.6×10^6 ad 1.4×10^7	2.5×10^7 ad 1.4×10^7	2.3×10^8 ae 1.3×10^8		2.5×10^8 ae 1.4×10^8		
BRI 0952 \star	QSO	4.43	360 e					4.3×10^7 af 1.5×10^8		
B3 J2330	HzRG	3.09	1950 e				1.1×10^8 ag 5.6×10^7			
TN J0121	HzRG	3.52	1050 e				1.3×10^8 ah 1.2×10^8			
6C 1909	HzRG	3.54	1470 e				1.7×10^8 ai 1.2×10^8			

Spectral resolution - CO 1-0 line - 30 GHz



Several measurements with different spectral resolution will permit to separate narrow line contribution from foregrounds with continuum spectrum