

Cosmic Reionization

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Leonid Chuzhoy (U.Chicago), Marcelo Alvarez (CITA/U.Toronto),
Eiichiro Komatsu, Jun Koda, Elizabeth Fernandez, Yi Mao (U. Texas),
Benedetta Ciardi (MPA), Rennan Barkana (Tel Aviv U)**

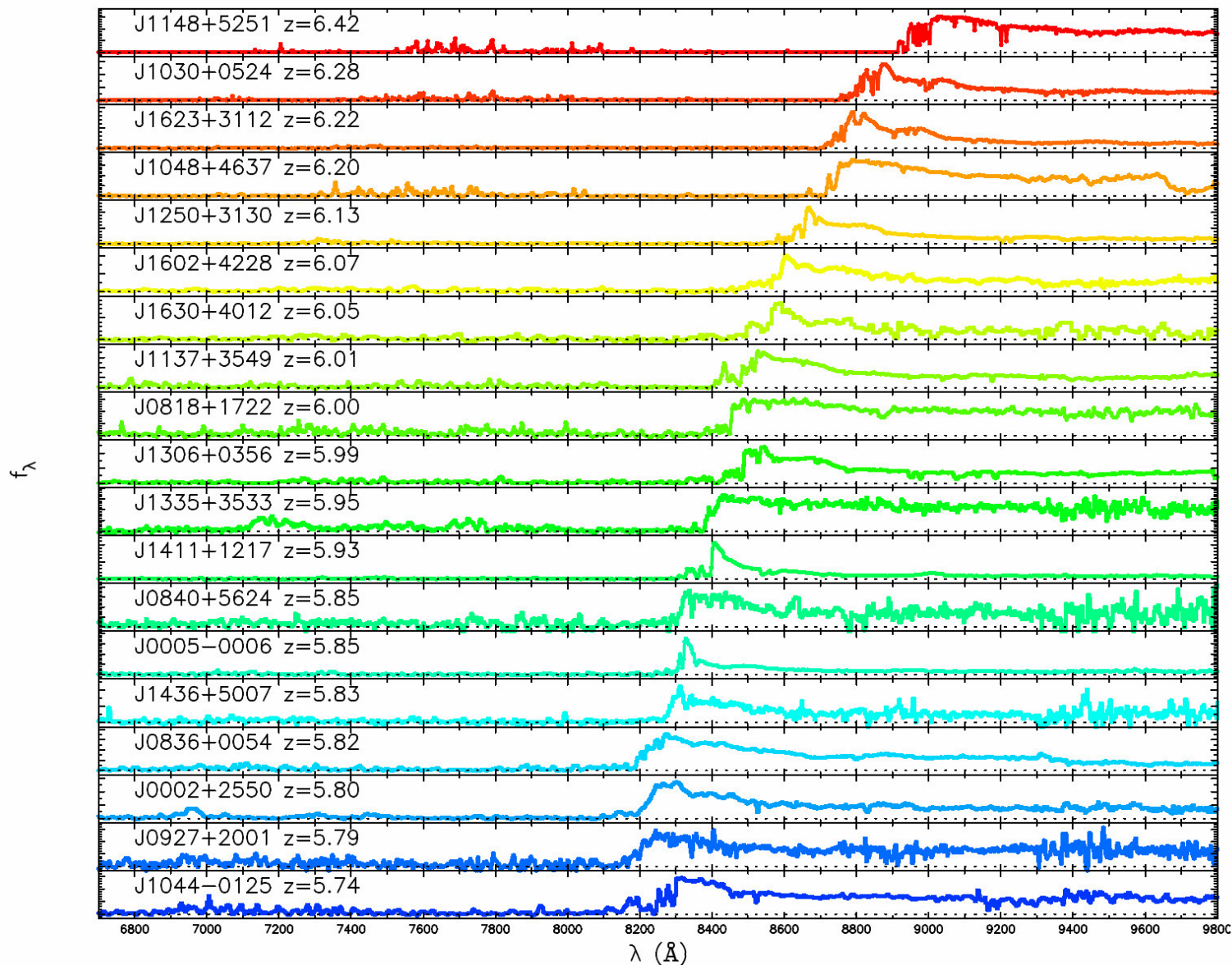
The Epoch of Reionization

- Absorption spectra of quasars have long shown that the intergalactic medium at redshifts $z < 6$ is highly ionized, with a residual neutral H atom concentration of less than 1 atom in 10^4 .
====> universe experienced an “epoch of reionization” before this.

The Epoch of Reionization

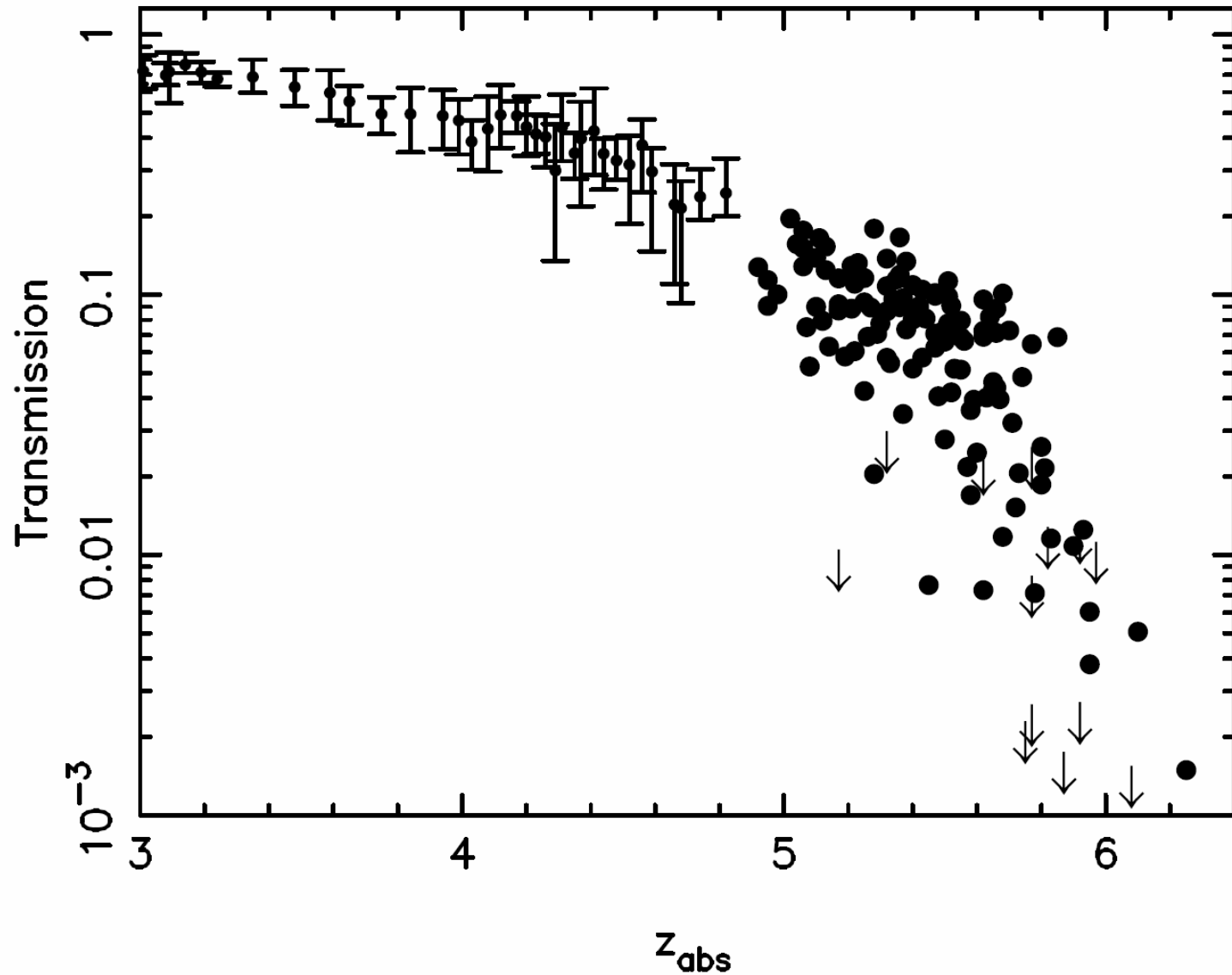
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- Sloan Digital Sky Survey quasars have been observed at $z > 6$ whose absorption spectra show dramatic increase in the H I fraction at this epoch as we look back in time.
====> epoch of reionization only just ended at $z \gtrsim 6$.

SDSS quasars show Lyman α opacity of intergalactic medium rises with increasing redshift at $z = 6 \rightarrow$ IGM more neutral \rightarrow reionization just ending?



Fan et al
(2005)

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Fan et al
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The Epoch of Reionization

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====> epoch of reionization only just ended at $z \gtrsim 6$.
- **The cosmic microwave background (CMB) exhibits polarization which fluctuates on large angular scales; WMAP finds that almost 9% of the CMB photons were scattered by free electrons in the IGM, but only 4% could have been scattered by the IGM at $z < 6$.**
====> **IGM must have been ionized much earlier than $z = 6$ to supply enough electron scattering optical depth**
====> **reionization already substantial by $z \gtrsim 11$**

EoR Probes the First Billion Years of Cosmic Star Formation

Observations of galaxies and quasars as far back in time as currently available suggest that galaxies dominated the production of the ionizing photons necessary to finish reionization by $z > 6$, while quasars were not numerous or luminous enough

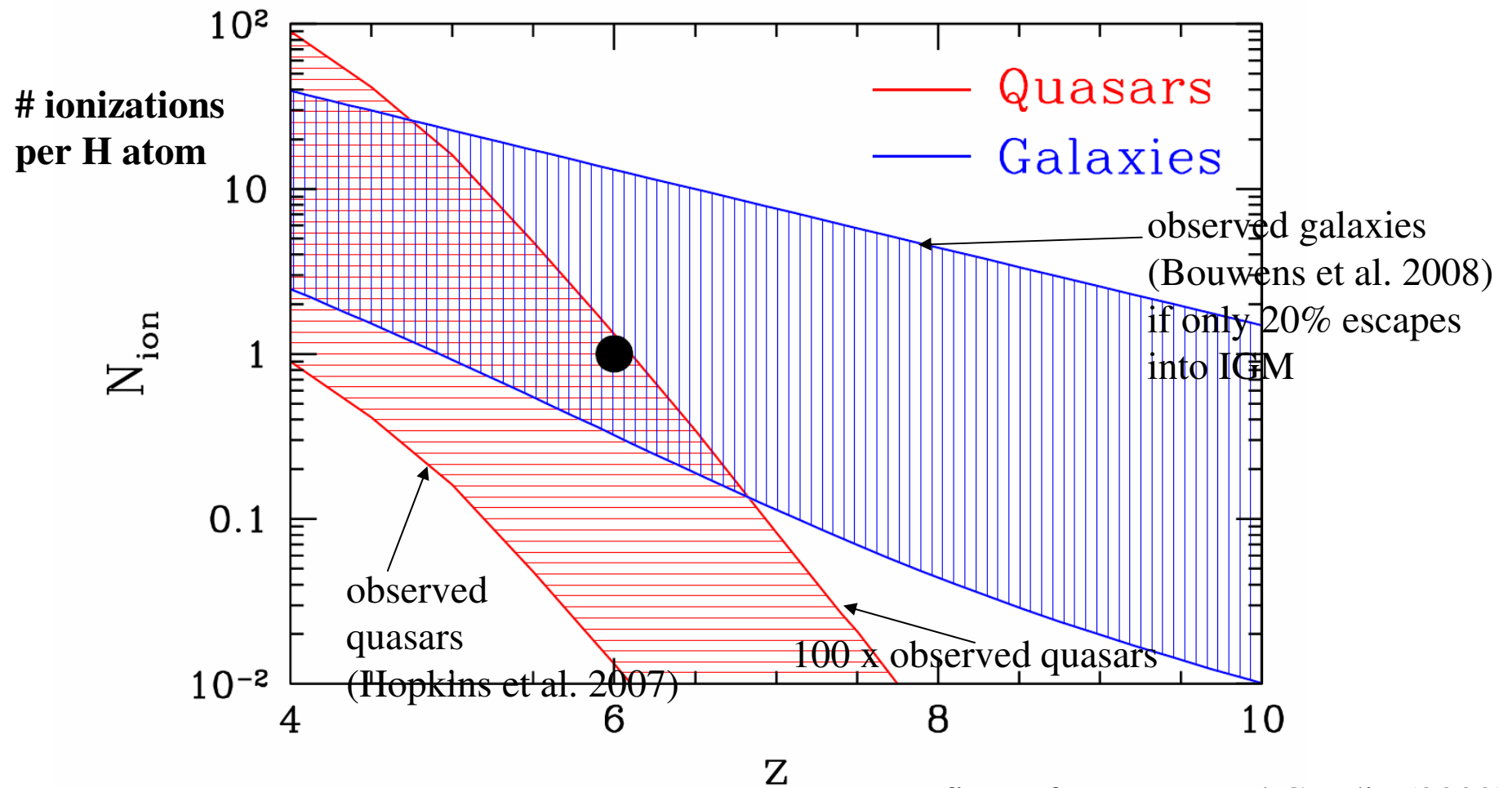
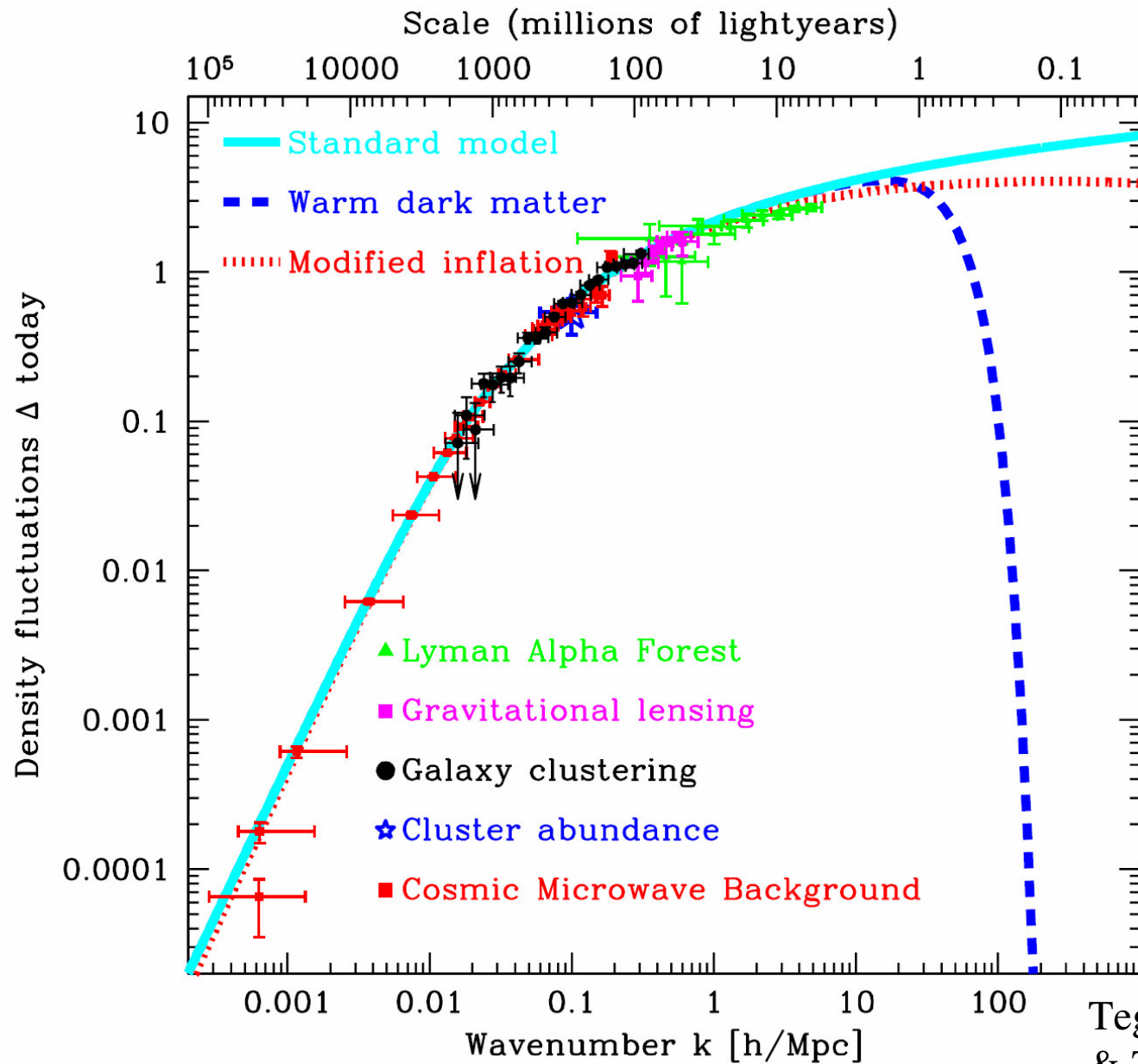


figure from Trac and Gnedin (2009)

EoR Probes the Primordial Power Spectrum Down to Very Small Scales



Tegmark
& Zaldarriaga (2008)

EoR Probes the Nature of Dark Matter

e.g. neutralino DM annihilations can partially ionize and heat the IGM ==>

CMB polarization observations of electron scattering optical depth

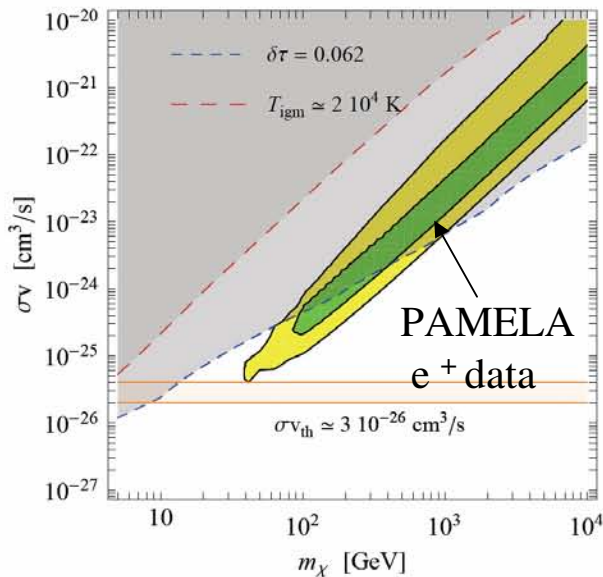
+

quasar absorption line Ly α forest observations of the IGM

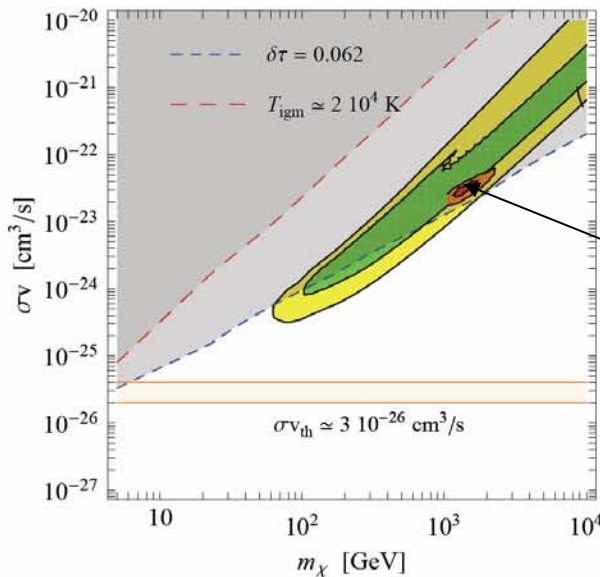
==> limits allowed “ m_χ (DM particle mass) - $\langle\sigma v\rangle$ (annihilation cross section)” parameter space

e.g. recent suggestions of large m_χ + large $\langle\sigma v\rangle$ into leptons suggested to explain anomalous charged-cosmic-ray signals reported by PAMELA, FERMI, & HESS experiments are ruled out; *too many free electrons produced*

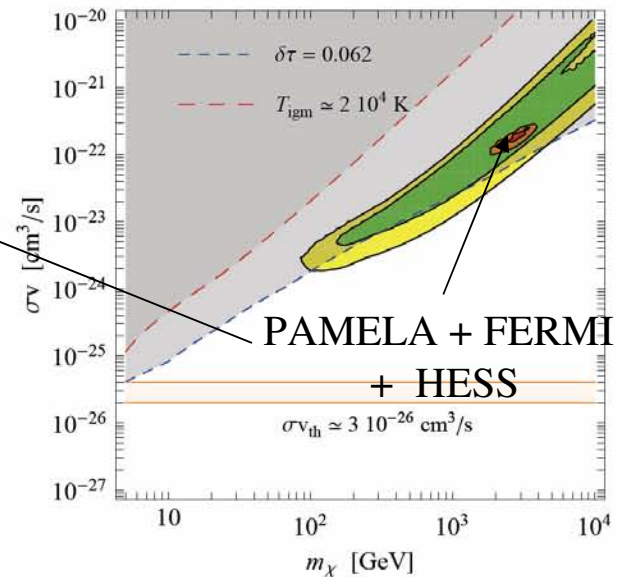
DM DM \rightarrow ee, NFW profile



DM DM \rightarrow $\mu\mu$, NFW profile



DM DM \rightarrow $\tau\tau$, NFW profile



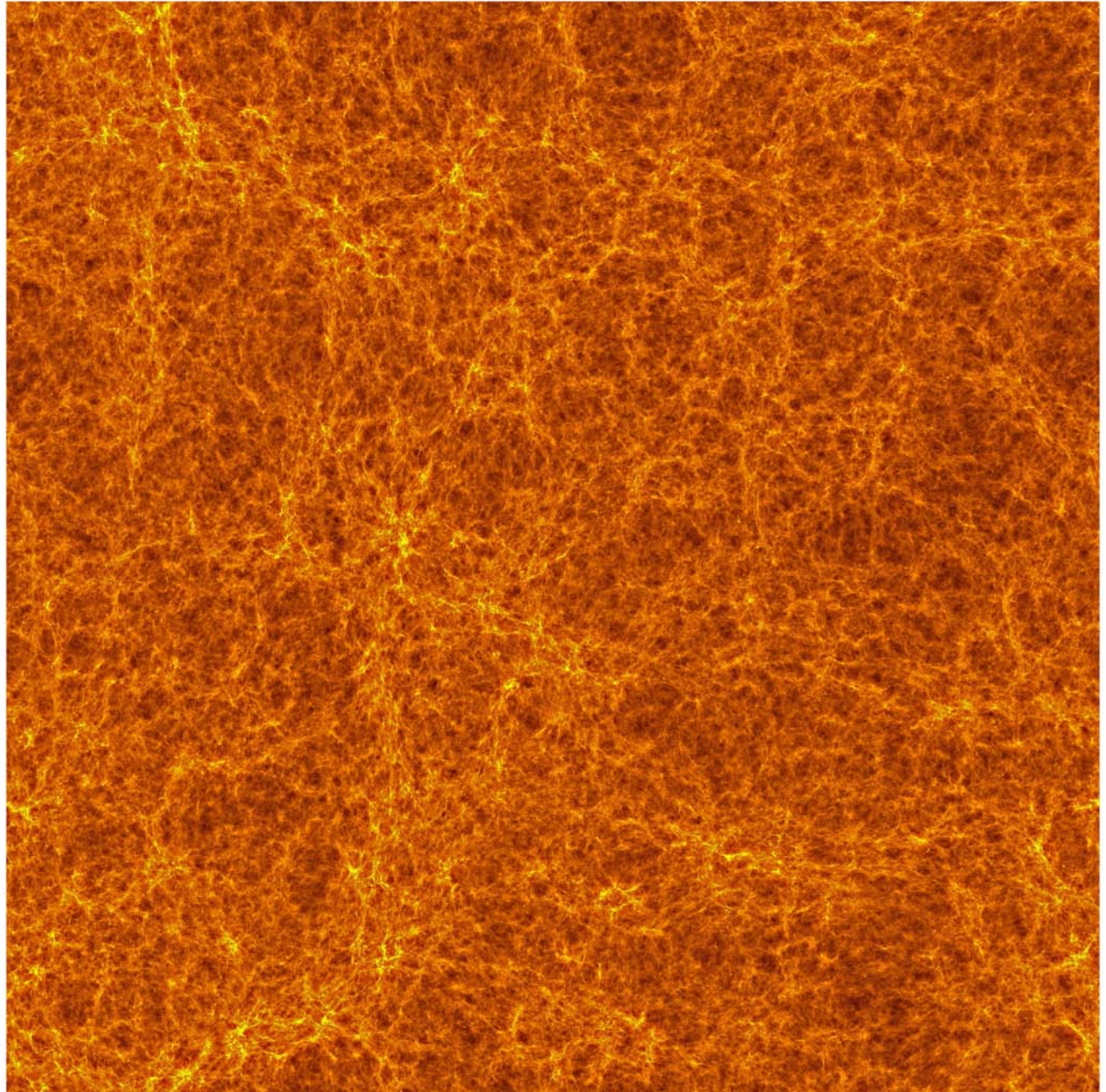
Cirelli, Iocco & Panci (2009)

Structure formation in Λ CDM at $z = 10$

simulation volume
=
(100 h^{-1} Mpc)³,
comoving

1624³ particles on
3248³ cells

Projection of
cloud-in-cell
densities of 20
Mpc slice



Simulating Cosmic Reionization at Large Scales I. : The Geometry of Reionization

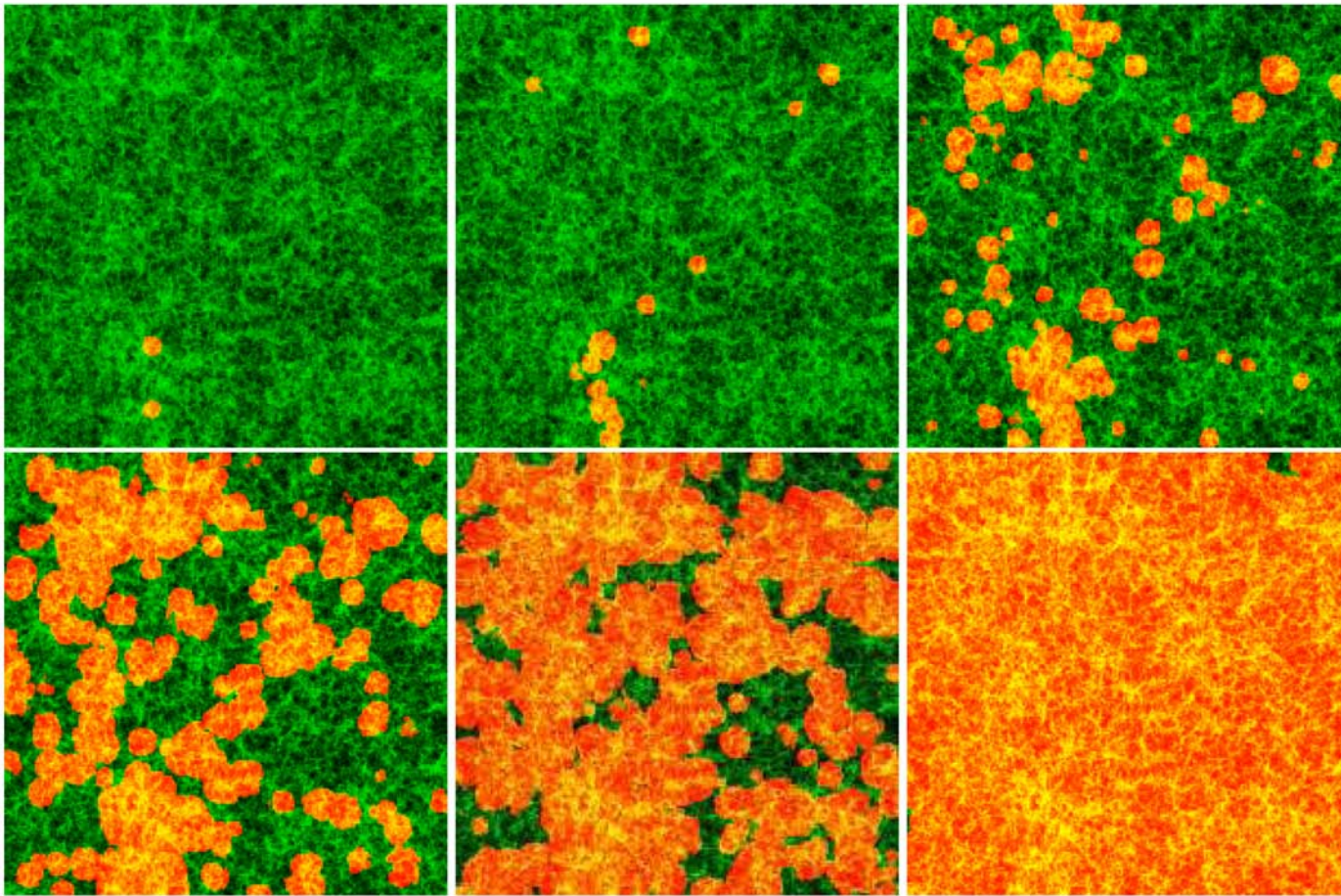
Iliev, Mellema, Pen, Merz, Shapiro & Alvarez (2006)
MNRAS, 369, 1625 (astro-ph/0512187)

$z =$

18.5

16.1

14.5



13.6

12.6

11.3

N-body + Radiative Transfer → Reionization simulation

- N-body simulation yields the density field and sources of ionizing radiation
 - PMFAST code (Merz, et al. 2005) with
 $1624^3 = 4.28$ billion particles, 3248^3 cells,
particle mass = $2.5 \times 10^7 M_{\text{sun}}$ (100 h^{-1} Mpc box),

 - Halo finder “on-the-fly” yields location, mass, other properties of all galaxies,
 $M \geq 2.5 \times 10^9 M_{\text{sun}}$ (100 h^{-1} Mpc box),
- e.g. $N_{\text{halo}} \sim 4 \times 10^5$ by $z \sim 8$ (WMAP1)
 $\sim 3 \times 10^5$ by $z \sim 6$ (WMAP3)

N-body + Radiative Transfer → Reionization simulation

- N-body simulation yields the density field and sources of ionizing radiation
 - **New**: 2nd generation N-body code CUBEP³M, a P³M code, massively paralleled (MPI+Open MP), 3072³ = 29 billion particles, 6144³ cells, particle mass = $5 \times 10^6 M_{\text{sun}}$ (114 h⁻¹Mpc box),
 - Halo finder “on-the-fly” yields location, mass, other properties of all galaxies,
 $M \geq 10^8 M_{\text{sun}}$ (114 h⁻¹Mpc box),
e.g. $N_{\text{halo}} \sim 10^7$ by $z \sim 8$ (WMAP5)

New

N-body + Radiative Transfer → Reionization simulation

- Radiative transfer simulations evolve the radiation field and nonequilibrium ionization state of the gas
 - New, fast, efficient C²-Ray code (Conservative, Causal Ray-Tracing) (Mellema, Ilev, Alvarez, & Shapiro 2006, *New Astronomy*, 11, 374) uses short-characteristics to propagate radiation throughout the evolving gas density field provided by the N-body results, on coarser grid of $\approx (200)^3$ to $(400)^3$ cells, for different resolution runs, from each and every galaxy halo source in the box.

e.g. $N_{\text{halo}} \sim 4 \times 10^5$ by $z \sim 8$ (WMAP1) ($> 2 \times 10^9 M_{\text{sun}}$)
 $\sim 3 \times 10^5$ by $z \sim 6$ (WMAP3) ($> 2 \times 10^9 M_{\text{sun}}$)
 $\sim 10^7$ by $z \sim 8$ (WMAP5) ($> 10^8 M_{\text{sun}}$)

for simulation volumes $\sim (100 h^{-1} \text{ Mpc})^3$

Every galaxy in the simulation volume emits ionizing radiation

- We assume a constant mass-to-light ratio for simplicity:

$$f_{\gamma} = \# \text{ ionizing photons released} \\ \text{by each galaxy per halo baryon} \\ = f_{*} f_{\text{esc}} N_i,$$

where

f_{*} = star-forming fraction of halo baryons,

f_{esc} = ionizing photon escape fraction,

N_i = # ionizing photons emitted per stellar baryon over stellar lifetime

e.g. $N_i = 50,000$ (top-heavy IMF), $f_{*} = 0.2$, $f_{\text{esc}} = 0.2 \rightarrow$
 $f_{\gamma} = 2000$

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 $f_{\gamma} = 2000$

- This yields source luminosity: $dN_{\gamma}/dt = f_{\gamma} M_{\text{bary}} / (\mu m_{\text{H}} t_{*})$,
 t_{*} = source lifetime (e.g. 2×10^7 yrs),
 M_{bary} = halo baryonic mass

C^2 - Ray : A New Method for Photon-Conserving Transport of Ionizing Radiation

Mellema, Iliev, Alvarez & Shapiro (2006) *New Astronomy*, 11, 374

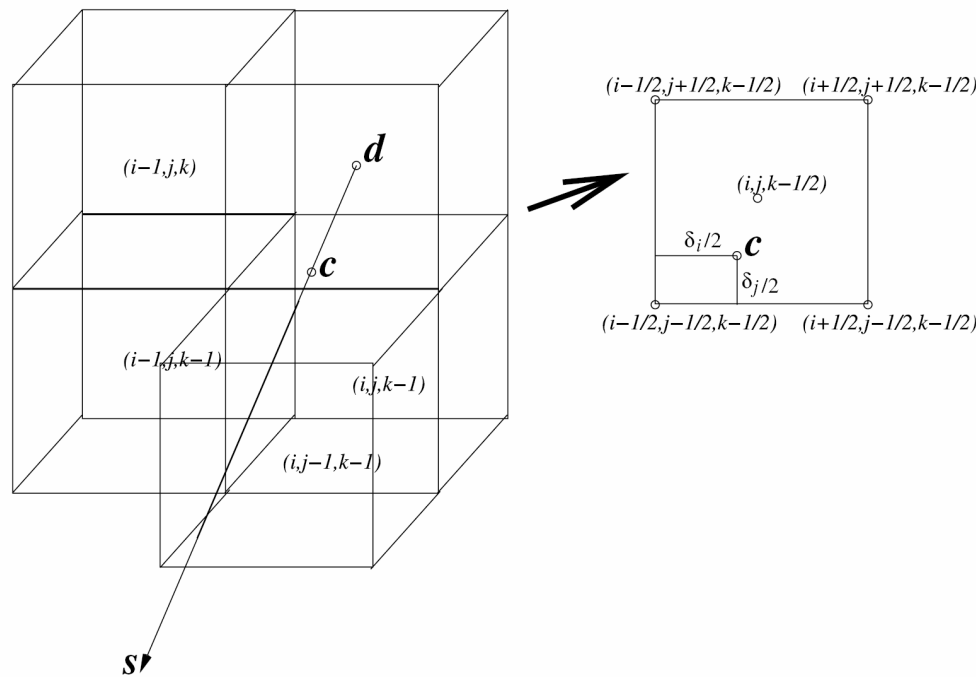
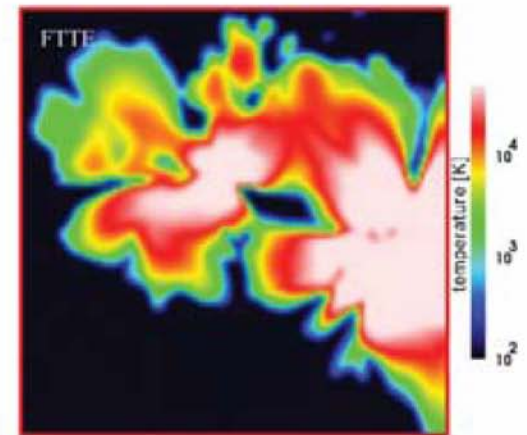
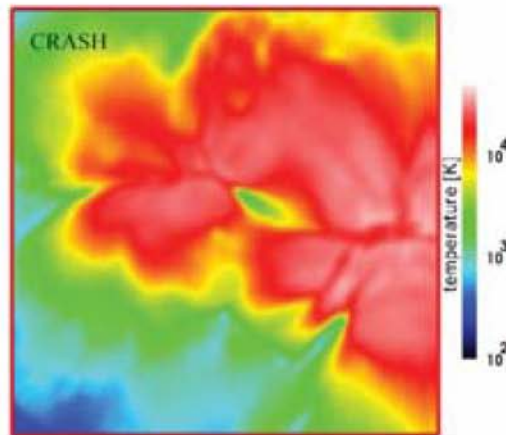
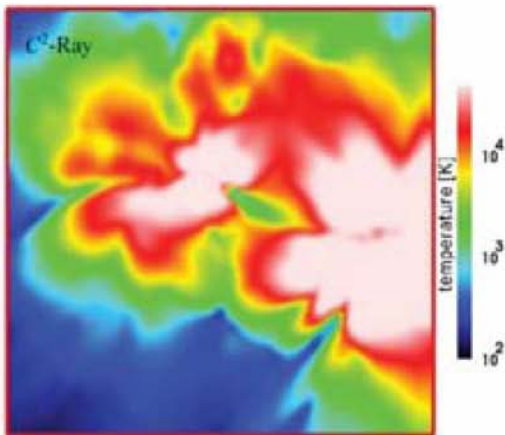


Fig. A.1. Short-characteristics ray tracing.

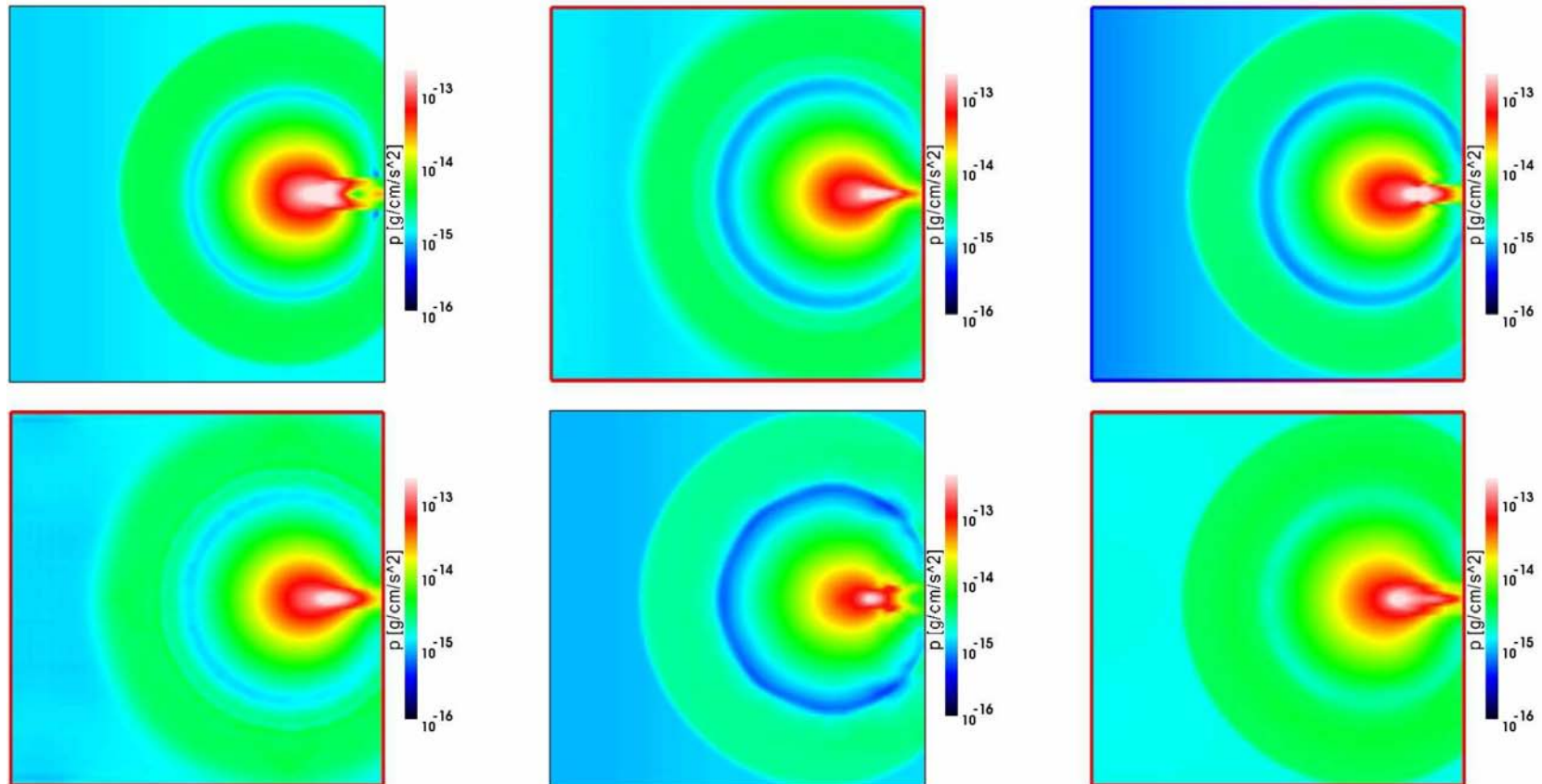
Cosmological Radiative Transfer Codes Comparison Project I. : The Static Density Field Tests

Iliev, Ciardi, ..., Shapiro, ... (2006) MNRAS, 371, 1057



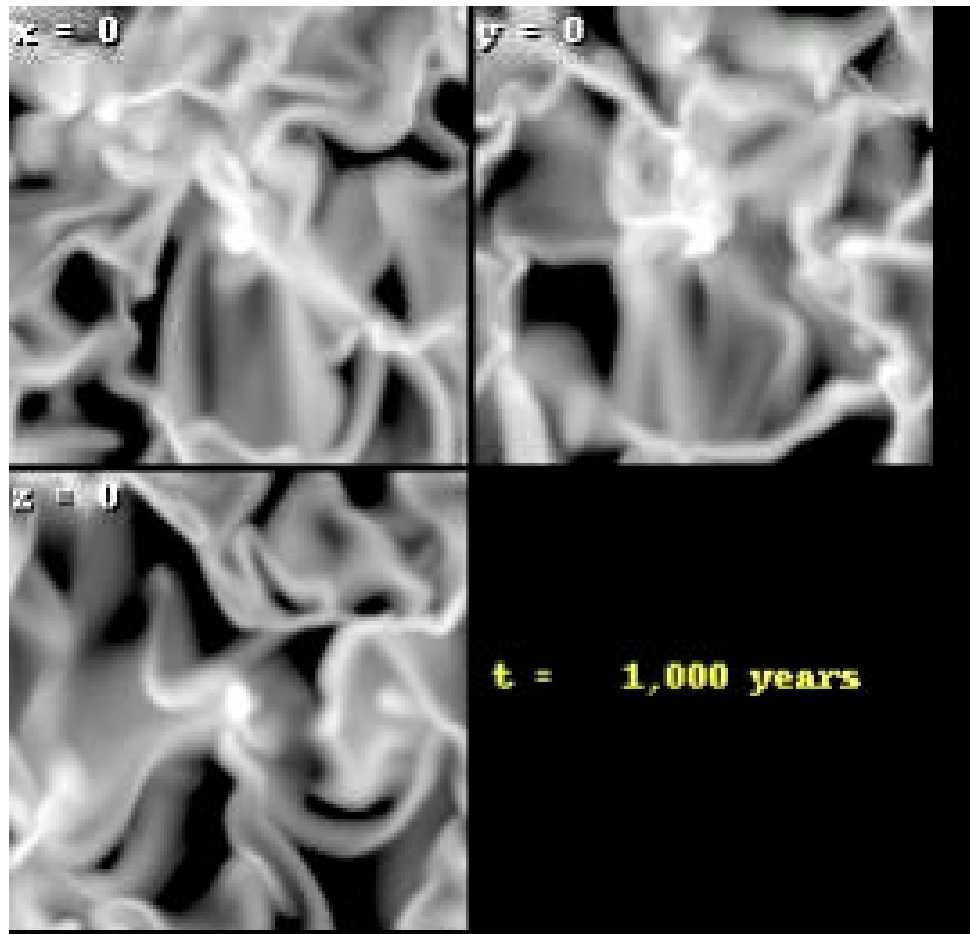
Cosmological Radiative Transfer Comparison Project II. : The Radiation- Hydrodynamic Tests

Iliev, Whalen, ..., Shapiro, ... 2009, MNRAS, in press
(astro-ph/0905.2920)



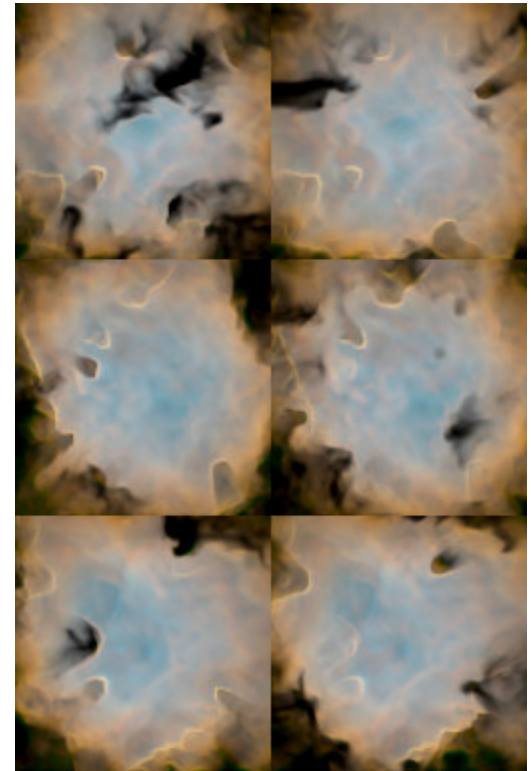
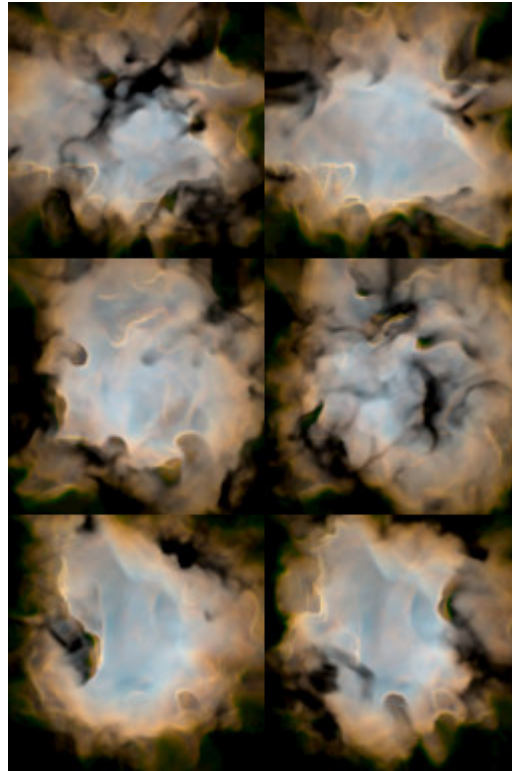
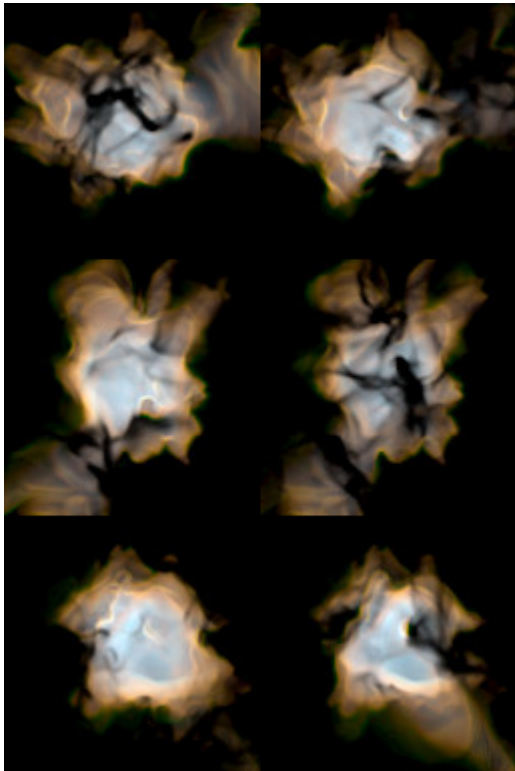
Dynamical H II Region Evolution in Turbulent Molecular Clouds

Mellema, Arthur, Henney, Iliev & Shapiro (2006) *ApJ*,
647, 397 (astro-ph/0512187)



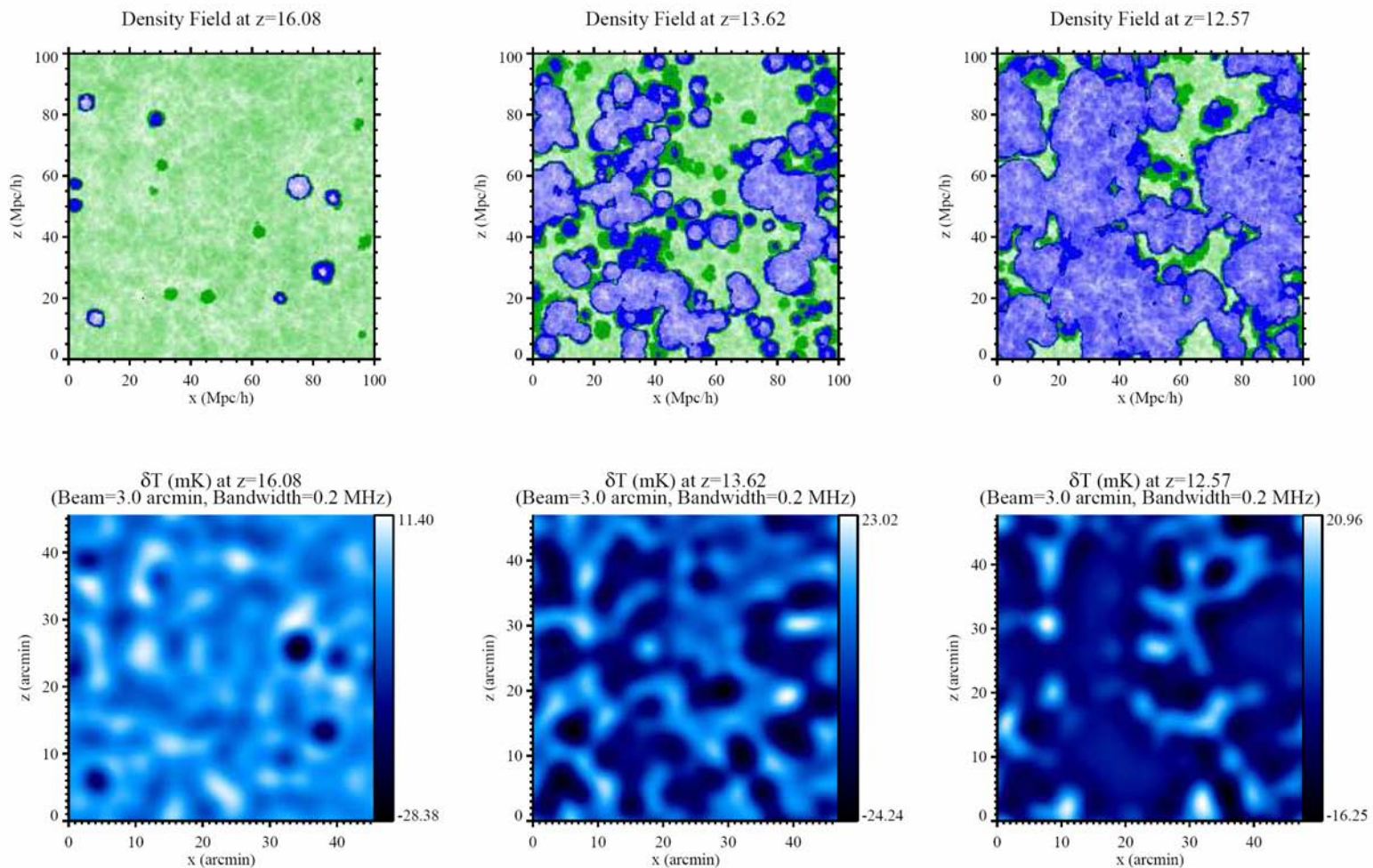
Dynamical H II Region Evolution in Turbulent Molecular Clouds

Mellema, Arthur, Henney, Iliev & Shapiro (2006) *ApJ*,
647, 397 (astro-ph/0512187)



Simulating Reionization at Large Scales II : The 21-cm Emission Features and Statistical Signals

Mellema, Iliev, Pen, & Shapiro (2006), MNRAS, 372,679
(astro-ph/0603518)



Low Frequency Array (LOFAR)

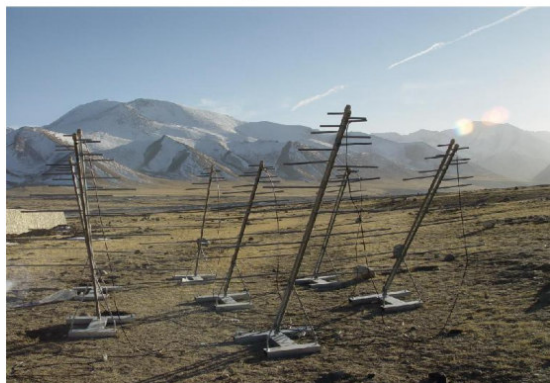


Murchison Widefield Array (MWA)



Primeval Structure Telescope (PAST/21CMA)

Prototype Tests, Ulaistai, Xin Jiang, China



Giant Meterwave Radio Telescope (GMRT)



Square Kilometer Array (SKA)



Q: How large must a reionization simulation volume be to predict 21-cm background fluctuations?

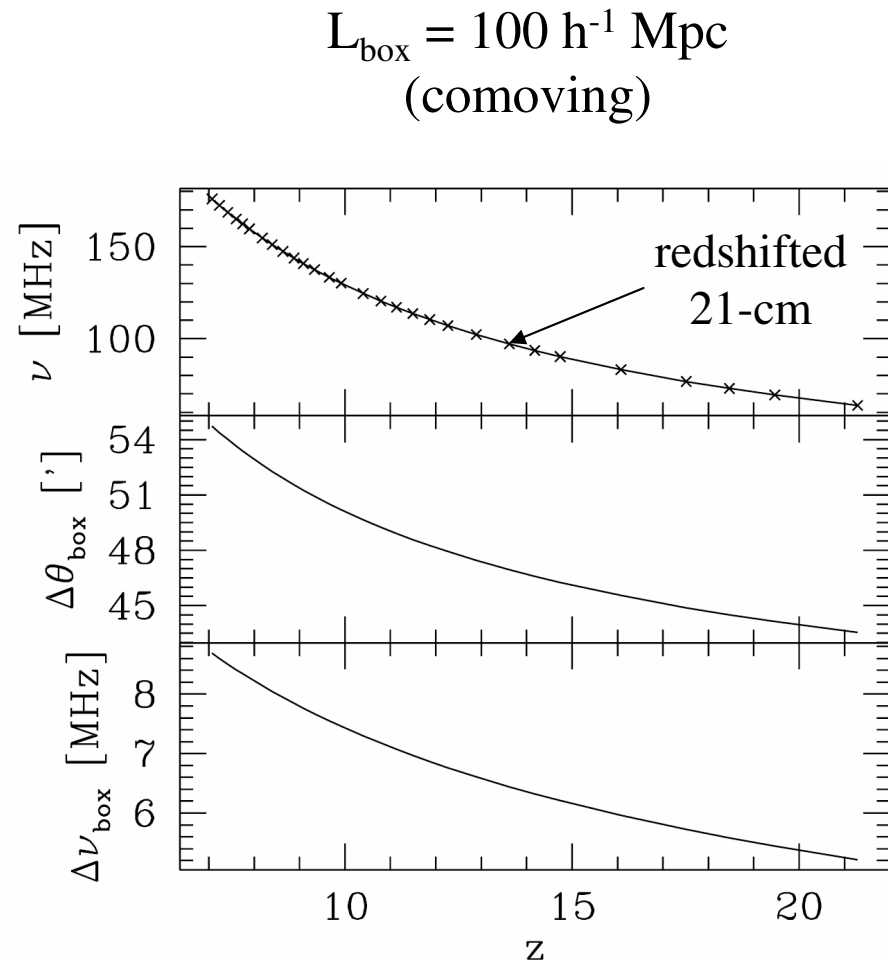
- Simulation volumes must exceed the beamsize and bandwidth of future radio arrays by a large enough factor to make 21-cm predictions meaningful.

e.g. for GMRT, MWA, LOFAR,

$$\Delta\theta_{\text{beam}} \geq 3'$$

$$\Delta\nu_{\text{bandwidth}} > 0.1 \text{ Mhz}$$

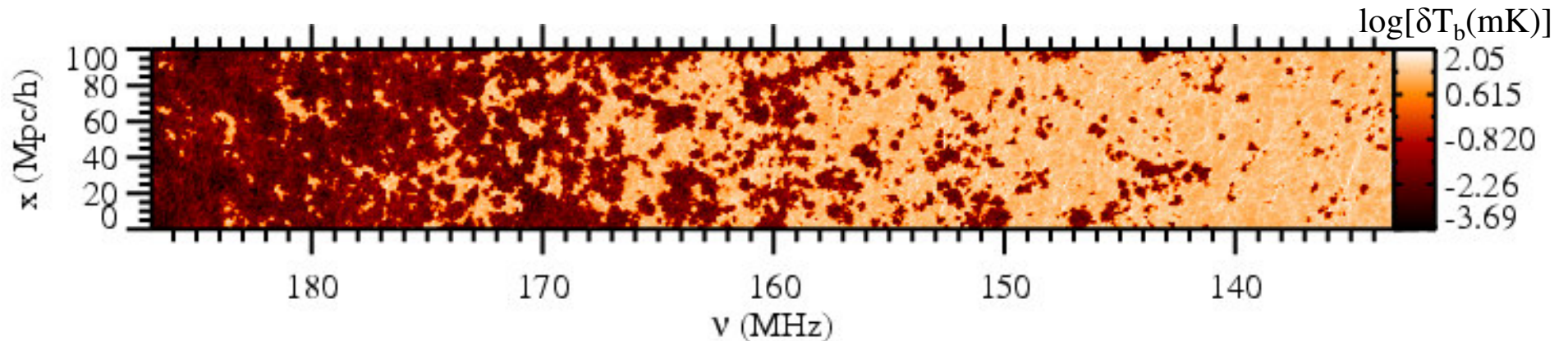
- Our $(143 \text{ Mpc})^3$ comoving volume simulation was the first to be large enough to predict 21-cm fluctuations.



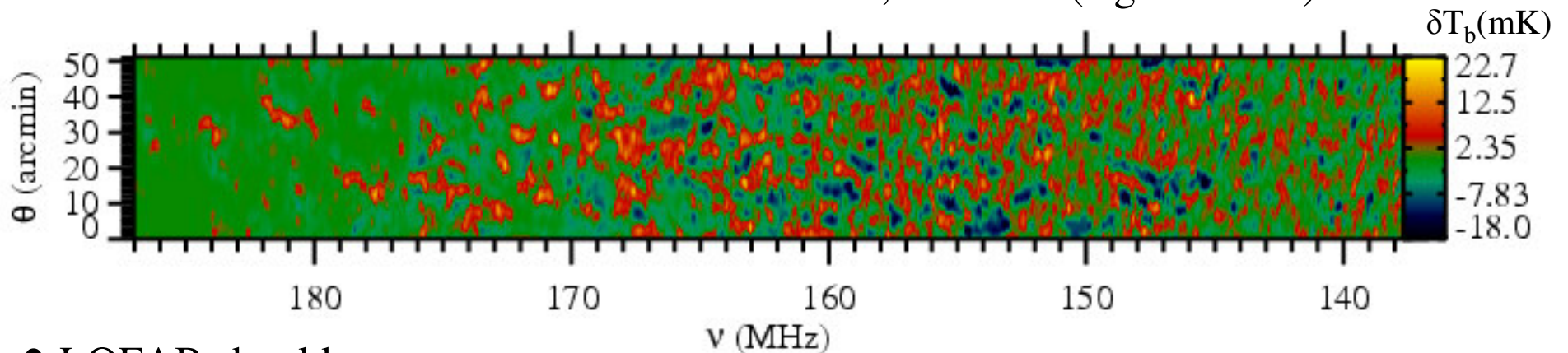
Reionization topology revealed by fluctuations in 21-cm brightness temperature, δT_b , along the line of sight

Iliev, Mellema, Pen, Bond, & Shapiro (2008), MNRAS, 384, 863 (astro-ph/0702099)

- mapping the sky along the LOS: high-resolution cuts in position-redshift space



- beam- and bandwidth-smoothed : 3 arcmin, 0.2 MHz (e.g. LOFAR)



- LOFAR should see
large ionized bubbles!

Case: $f_\gamma = 250$, subgrid clumping factor $C(z)$, WMAP3

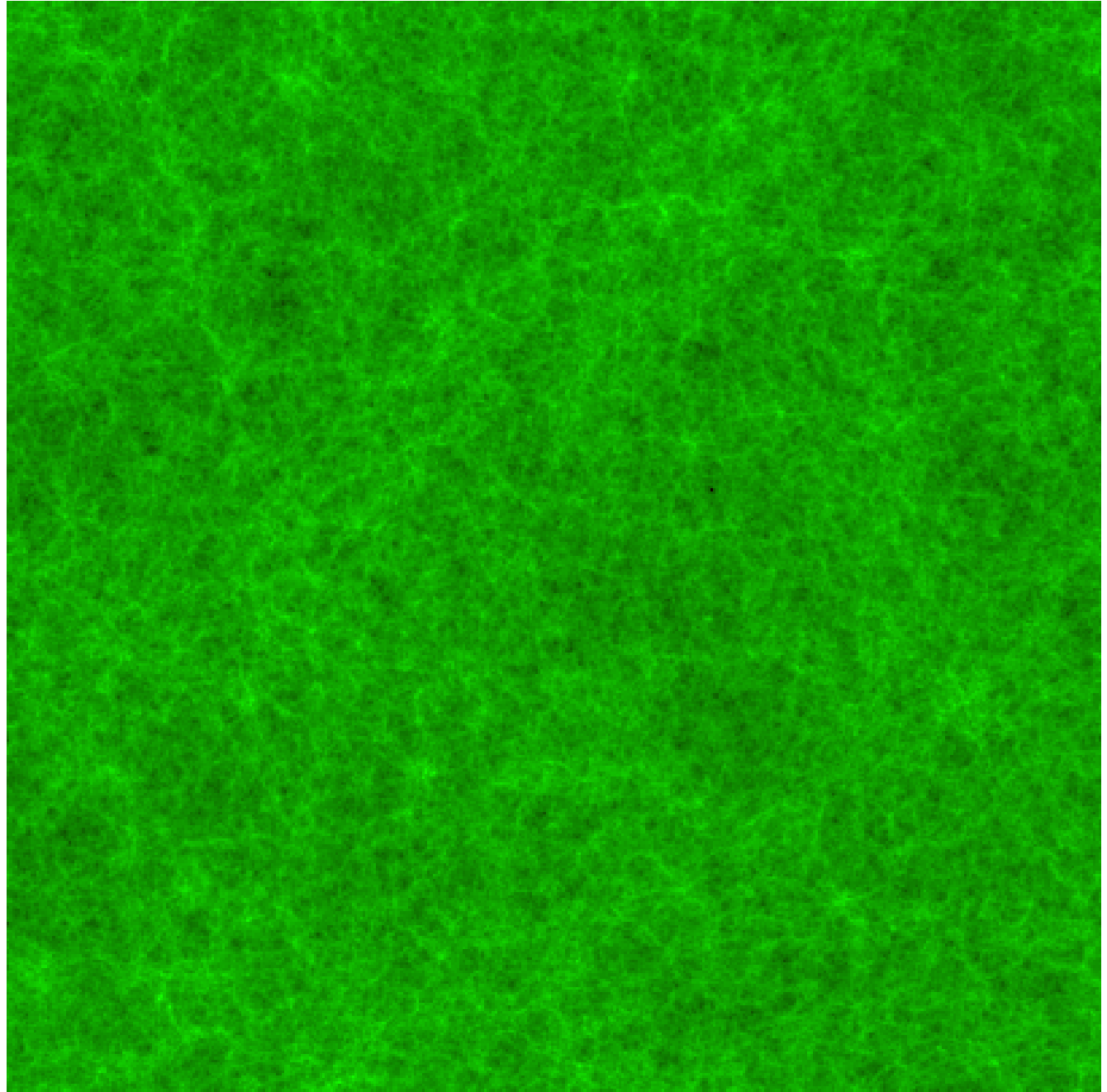
The Geometry of Reionization

for source efficiency
 $f_{\gamma} = 2000$

from $z = 20$ to $z = 12$

a cut through the
simulation volume,
one cell deep

gas density (green in
neutral regions, yellow
in ionized regions),
H II regions (red)
and
sources (dots)



Kinetic Sunyaev-Zel'dovich Effect from Patchy Reionization

- kSZ effect is the CMB temperature anisotropy induced by electron scattering by free electrons moving along the line-of-sight:

$$\frac{\Delta T}{T_{\text{CMB}}} = \int d\eta e^{-\tau_{\text{es}}(\eta)} a n_e \sigma_T \mathbf{n} \cdot \mathbf{v},$$

where η is conformal time,

$$\eta = \int_0^t dt' / a(t')$$

$(\delta T/T_{\text{CMB}})$ Maps of the Kinetic Sunyaev-Zel'dovich Effect from Radiative Transfer Simulations of Patchy Reionization

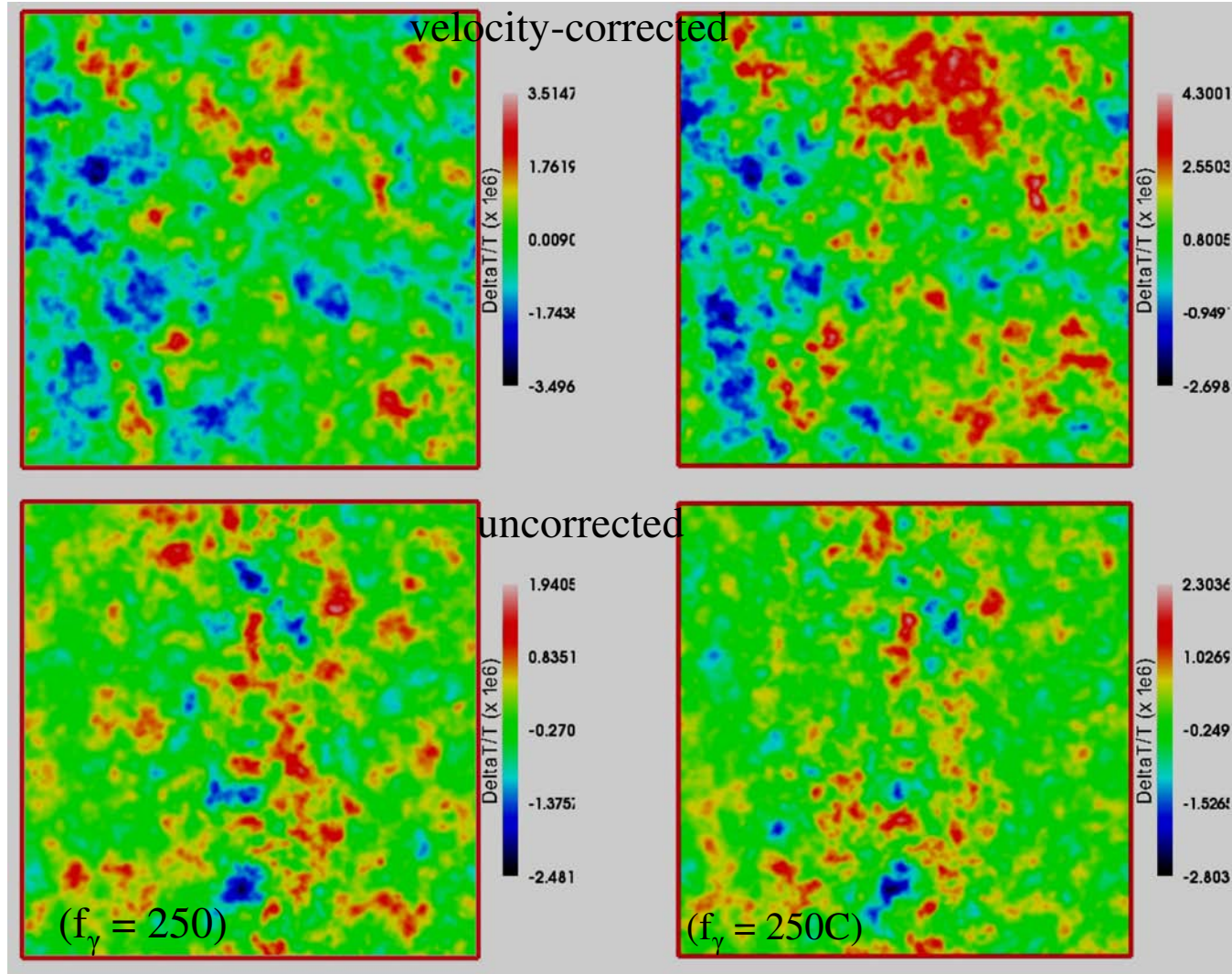
Iliev, Pen, Bond, Mellema & Shapiro (2007), ApJ, 660, 933; (astro-ph/0609592)

Iliev, Mellema, Pen, Bond, & Shapiro (2008), MNRAS, 384, 863; (astro-ph/0702099)

- Box size 100/h
Mpc comoving
→ 50' x 50'

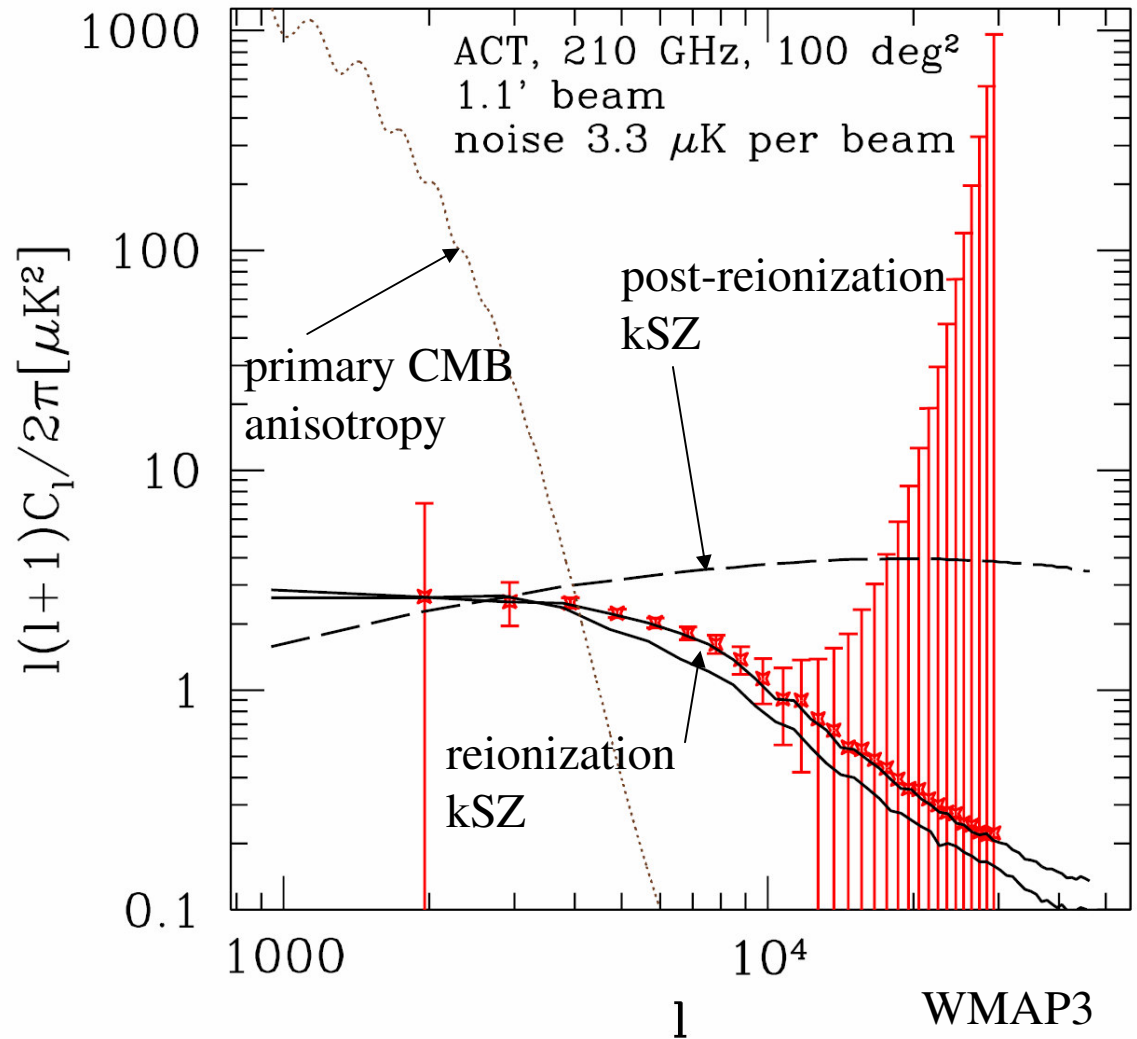
- Even so large a
box is missing
large-scale
power in
velocity field
perturbations

- must correct
for missing
large-scale
velocity
perturbations



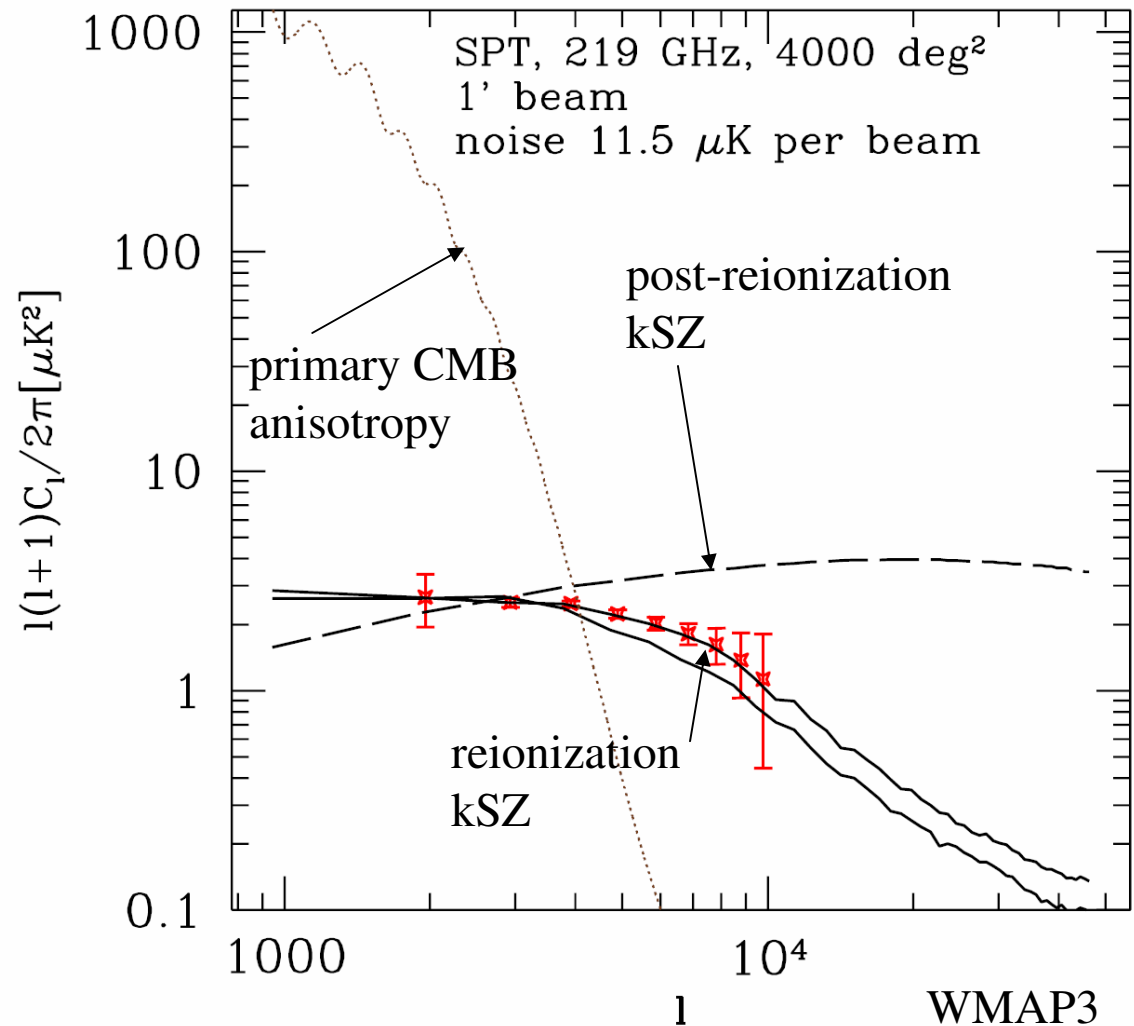
Observability of the kSZ from reionization: sky power spectrum

- Predicted kSZ from reionization simulations and *Atacama Cosmology Telescope* (ACT) expected sensitivity;
- primary CMB and post-reionization kSZ added to noise error bars for reionization signal;
- reionization kSZ signal is comparable to that from post-reionization, so necessary to separate them to extract info on reionization, alone.
- In principle, results show that reionization signal should be observable.



Observability of the kSZ from reionization: sky power spectrum

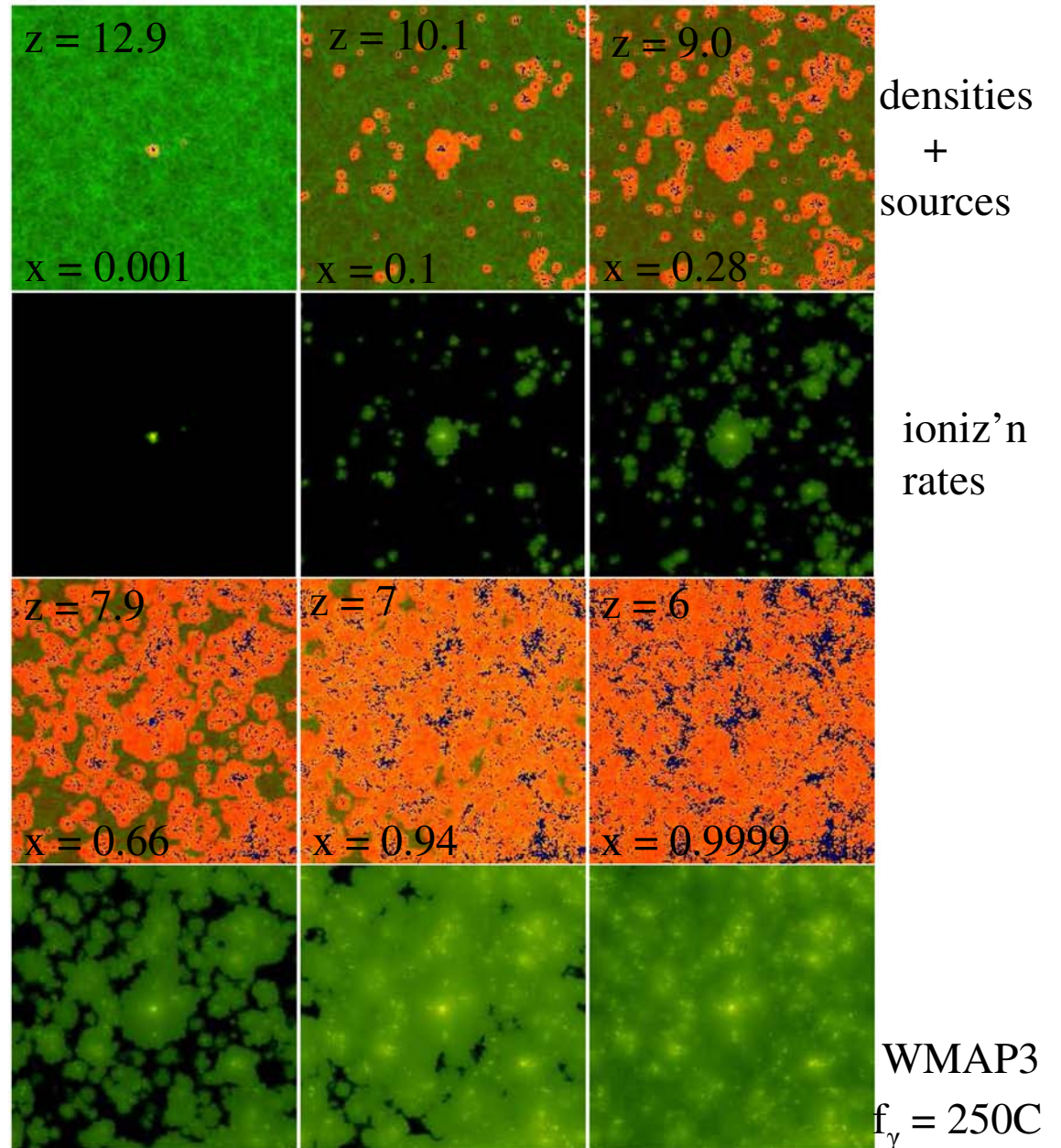
- Predicted kSZ from reionization simulations and *South Pole Telescope* (SPT) expected sensitivity;
- primary CMB and post-reionization kSZ added to noise error bars for reionization signal;
- reionization kSZ signal is comparable to that from post-reionization, so necessary to separate them to extract info on reionization, alone.
- In principle, results show that reionization signal should be observable.



Effect of IGM on Observability of Ly- α Sources During Reionization

(Iliev, Shapiro, McDonald, Mellema, Pen 2008, MNRAS, 391, 63)

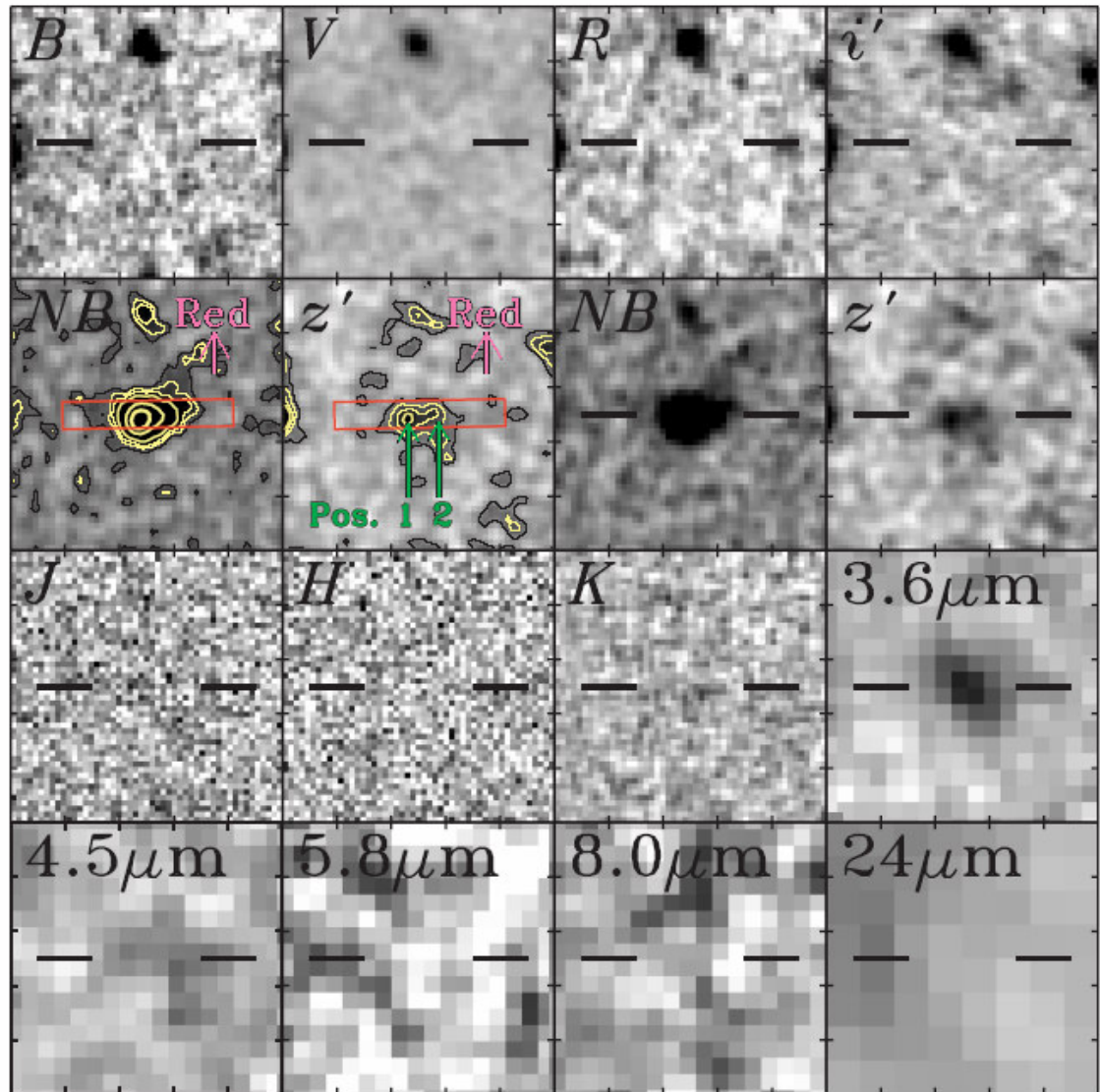
- reionization, centered on most massive halo in $(100h^{-1} \text{ Mpc})^3$:
 $M(z = 6) = 1.5 \times 10^{12} M_{\text{solar}} \rightarrow$
rare ($\sim 5\text{-}\sigma$) density peak \rightarrow
bright Ly- α emitter ;
- H II regions form first around such density peaks and grow continuously, as halos form inside, clustered around peak \rightarrow
Ly- α emitters will be centered on large HII regions ;
(blue dots = source cells)
- Fluctuating GP optical depth inside H II regions can still affect Ly- α source detection, though;



Discovery of a Giant Lyman Alpha Emitter Near the Epoch of Reionization

(Ouchi et al. 2009, ApJ, 696, 1164)

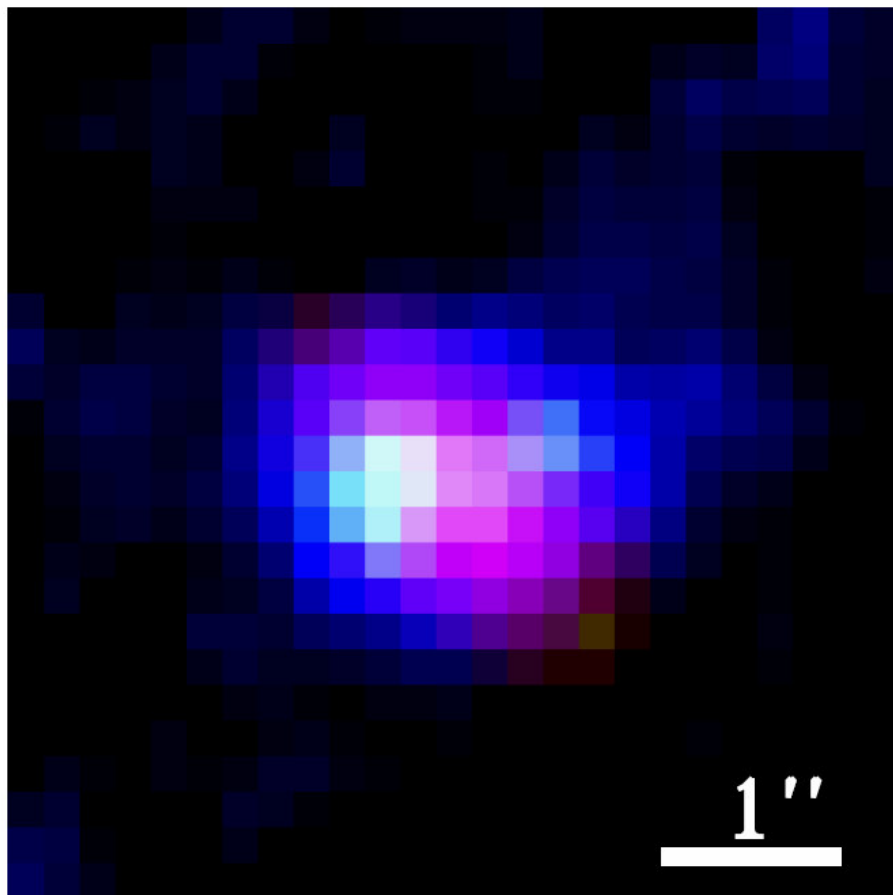
- Optical and IR images of *Himiko* by Subaru, UKIDSS-Ultra Deep Survey, and *Spitzer*/IRAC → Ly α emitter (LAE) at $z = 6.595$
- High luminosity:
 $L(\text{Ly } \alpha) = 3.9 \times 10^{43} \text{ erg s}^{-1}$
- Highest luminosity LAE in survey volume 10^6 Mpc^3



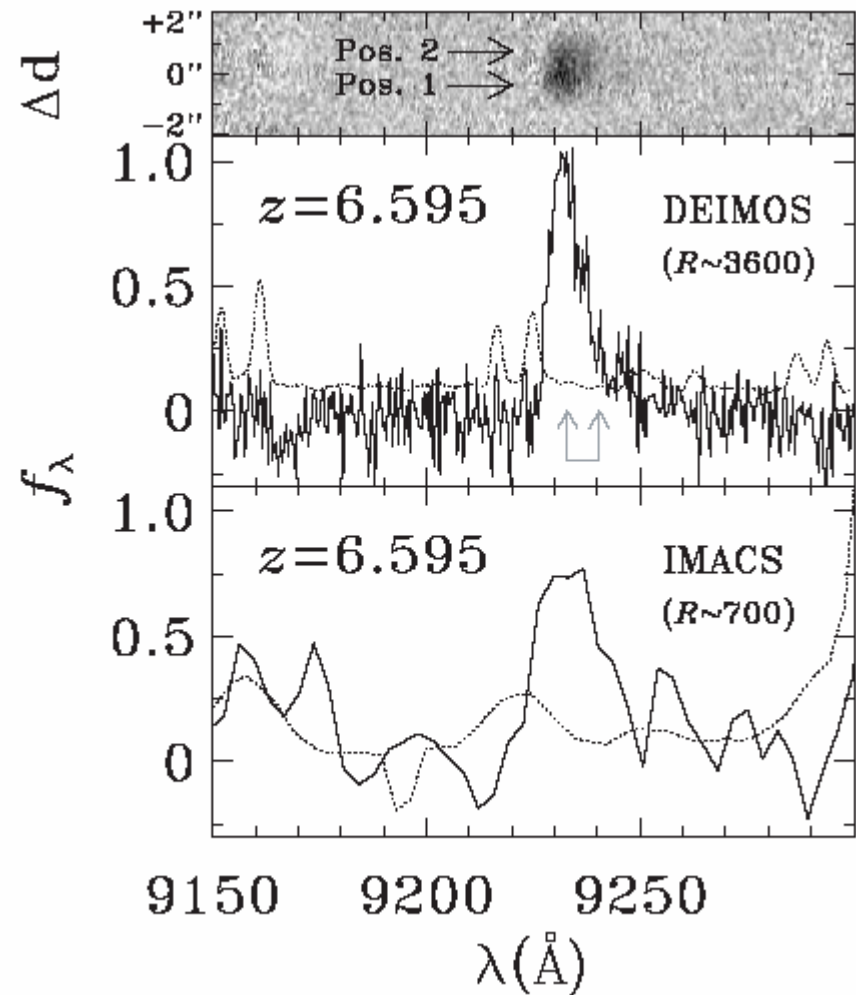
Discovery of a Giant Lyman Alpha Emitter Near the Epoch of Reionization

(Ouchi et al. 2009, ApJ, 696, 1164)

- Composite image shows extended emission 3" (or 17 kpc) across



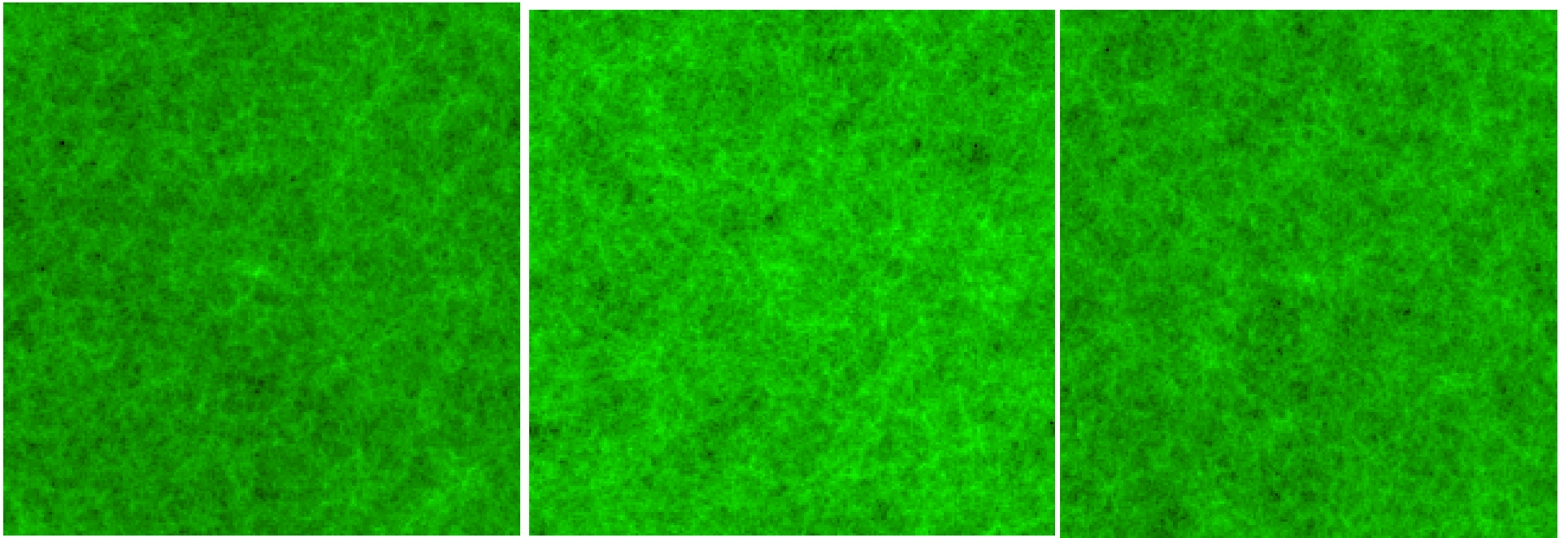
- Keck and Megellan Telescopes confirm Ly α line spectroscopically



Effect of IGM on Observability of Ly- α Sources During Reionization

(Iliev, Shapiro, McDonald, Mellema, Pen 2008, MNRAS, 391, 63)

← 143 Mpc →



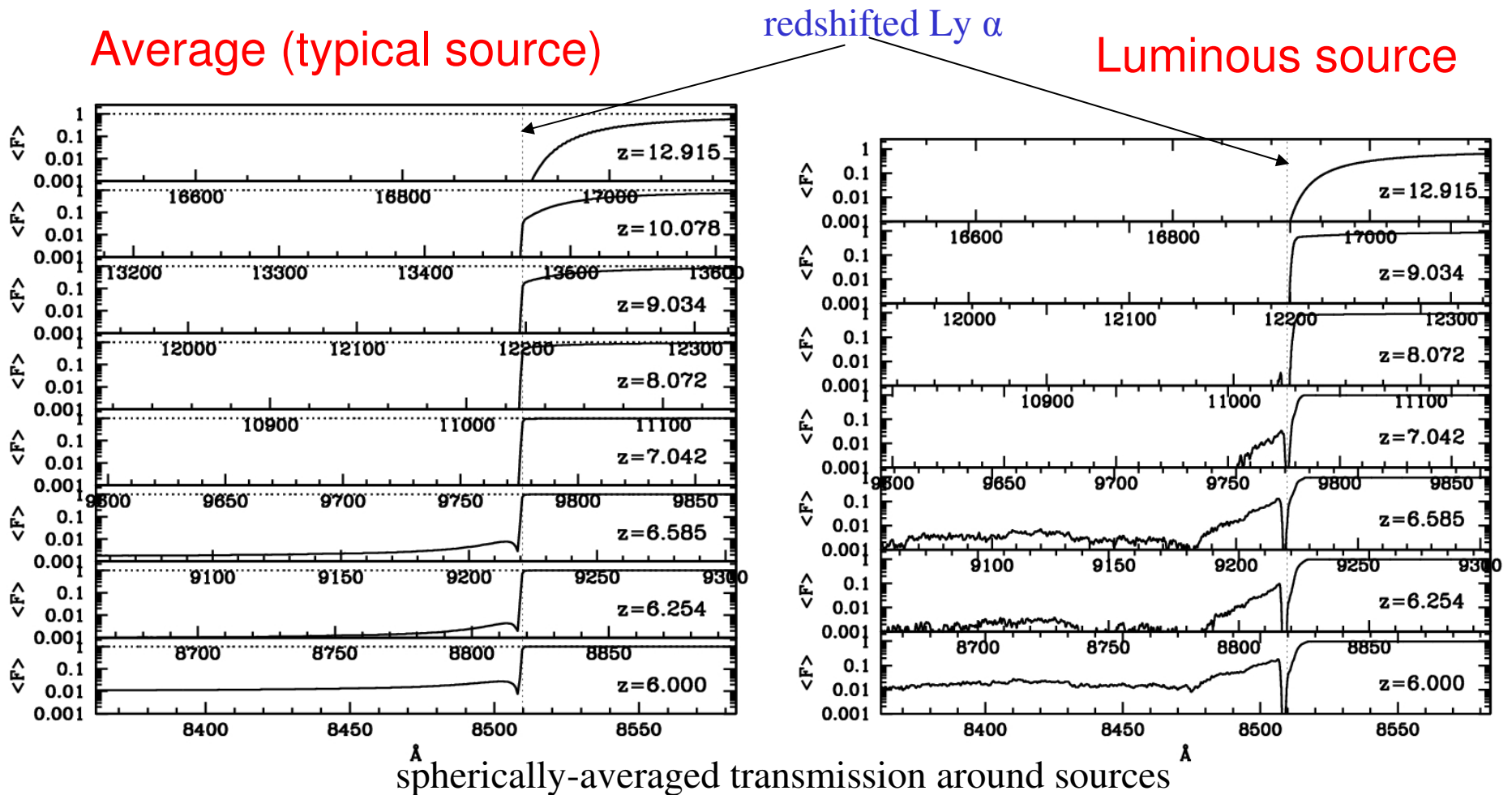
xz

xy

yz

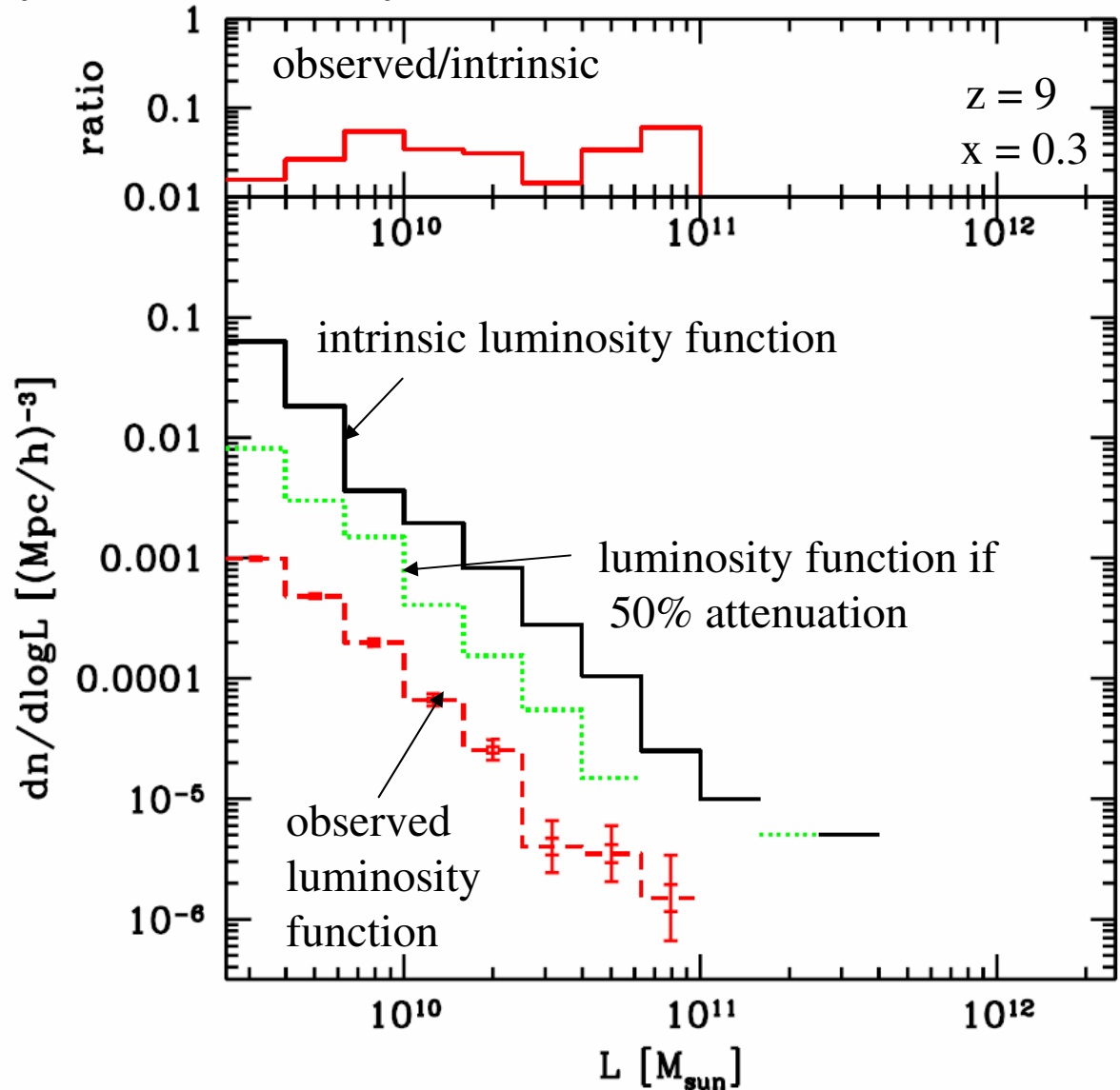
Mean Ly- α transmission vs. redshift

- Strong damping wings at $z > 10$, only minor differences between average and luminous source.
- Some transmission on blue side of line, as IGM slowly becomes transparent; large proximity transmission region for luminous sources.



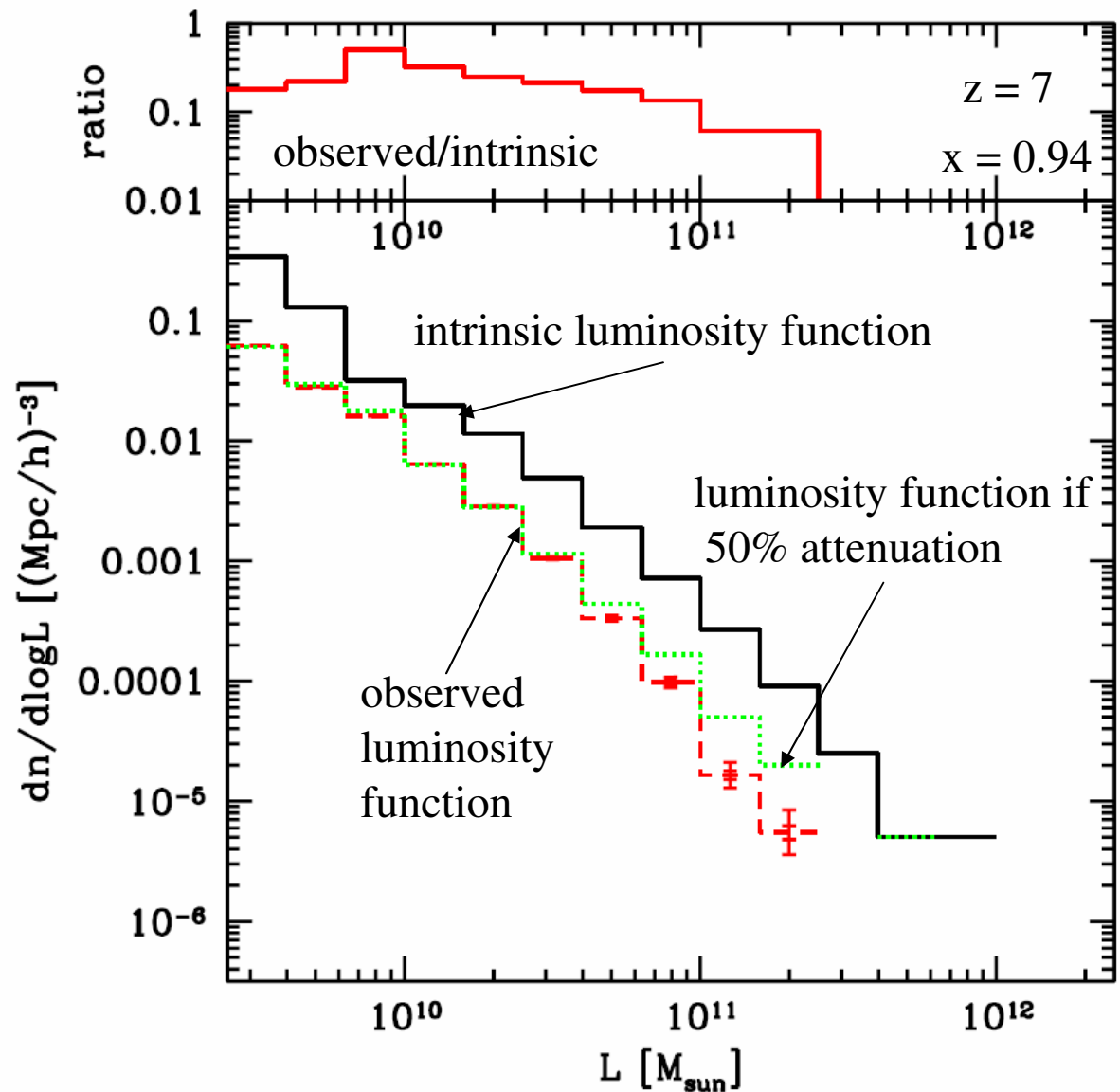
Luminosity Function of High-Redshift Ly- α Sources is Filtered By the Partially Ionized IGM

- Assuming all halos are Ly- α emitters with $M/L = \text{constant}$ and 160 km/s Gaussian line profile;
- At $z = 9$, observed luminosity function is reduced by 0.01 -0.1 from intrinsic luminosity function;
- Attenuation exceeds 50% (i.e. blue-half absorbed, red-half not), since damping wings reduce the red-half, as well..



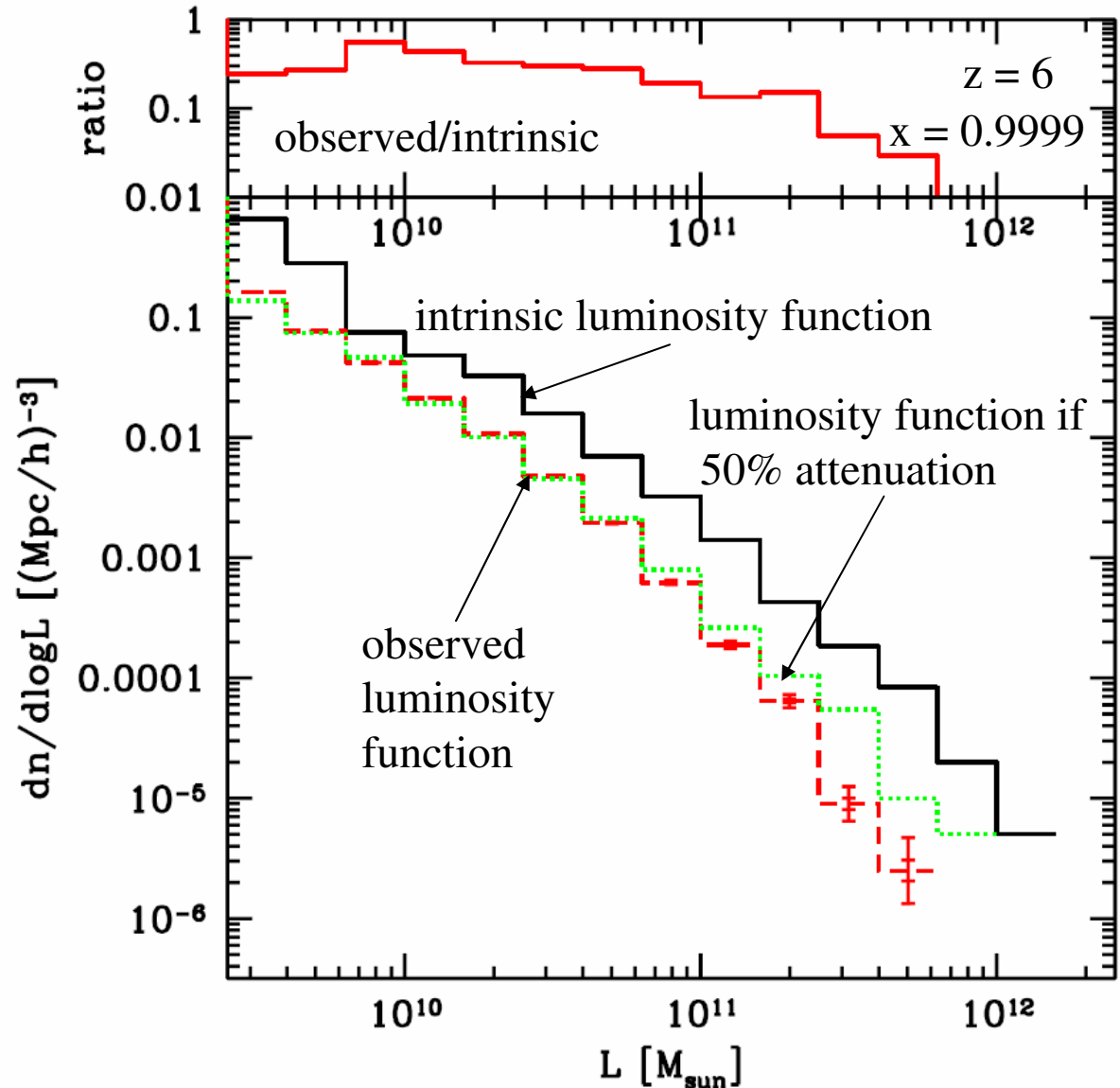
Luminosity Function of High-Redshift Ly- α Sources is Filtered By the Partially Ionized IGM

- Assuming all halos are Ly- α emitters with $M/L = \text{constant}$ and 160 km/s Gaussian line profile;
- At $z = 7$, observed luminosity function is reduced by 0.5 from intrinsic luminosity function, except at bright end where factor is 0.1 or below;
- Attenuation exceeds 50% (i.e. blue-half absorbed, red-half not), now, only at bright end.



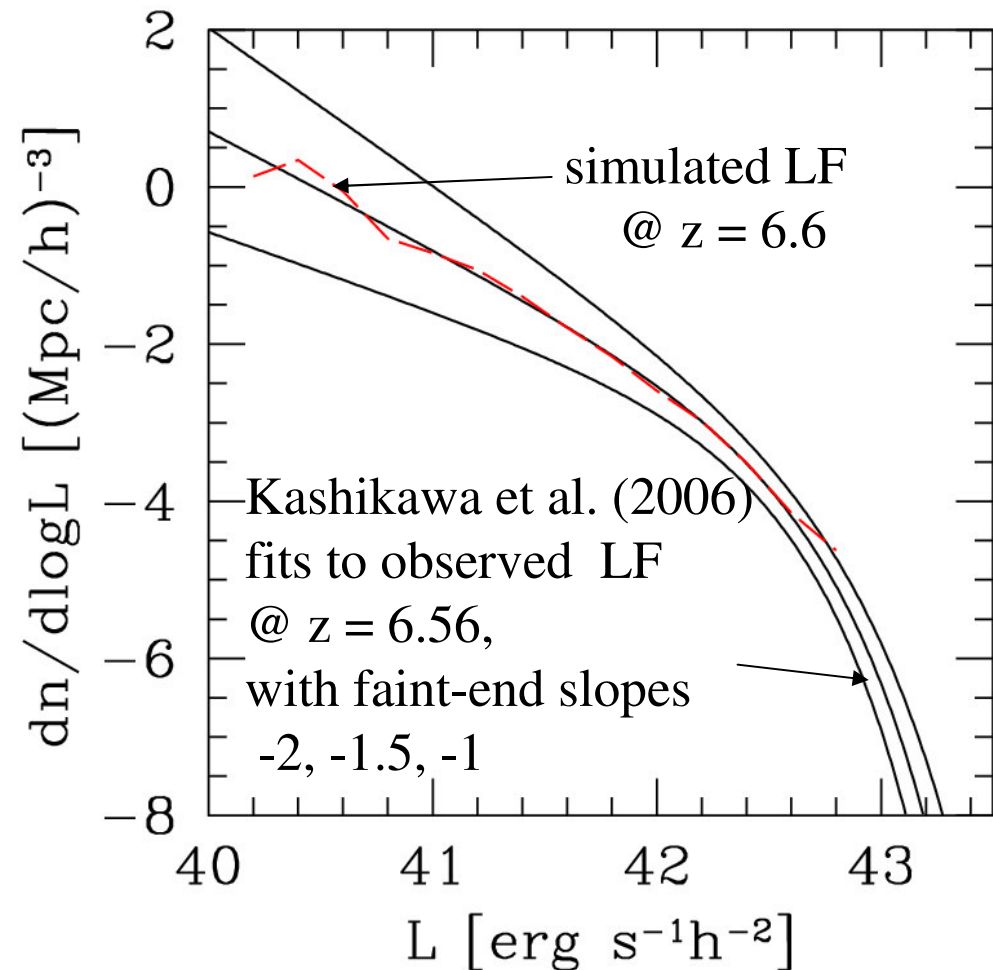
Luminosity Function of High-Redshift Ly- α Sources is Filtered By the Partially Ionized IGM

- Assuming all halos are Ly- α emitters with $M/L = \text{constant}$ and 160 km/s Gaussian line profile;
- Even at $z = 6$, observed luminosity function is reduced by 0.5 from intrinsic luminosity function, except at bright end where factor is 0.1 or below;
- Attenuation exceeds 50% (i.e. blue-half absorbed, red-half not), now, but only at bright end.



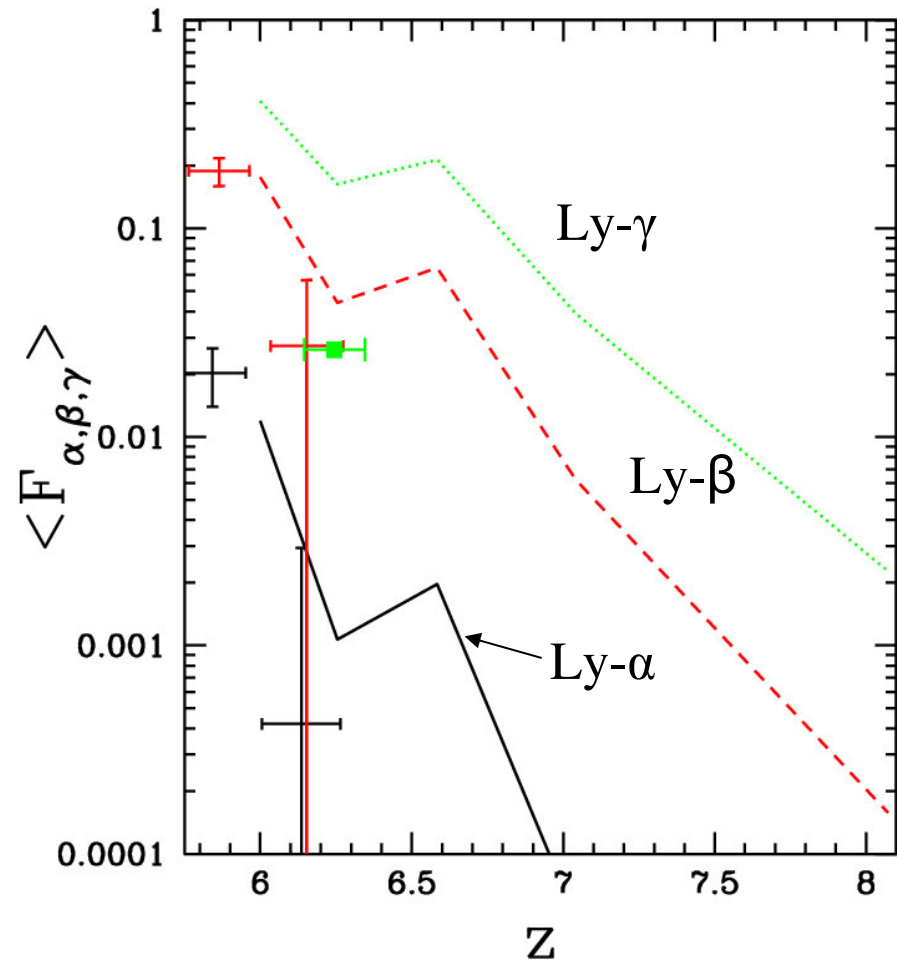
Predicting the observed luminosity function of Ly- α sources at the end of the EOR : simulations vs. observations

- To use our simulations to predict the observed LF, we “tune” the assumed M/L per halo to match the number density of sources in our simulations to the observed one reported by Kashikawa et al. (2006)
- → simulated LF is an excellent match of the shape, for an assumed faint-end slope of -1.5 for the fit to the observations.
- → the **majority** of sources responsible for reionization are **too faint** to be observed at present.



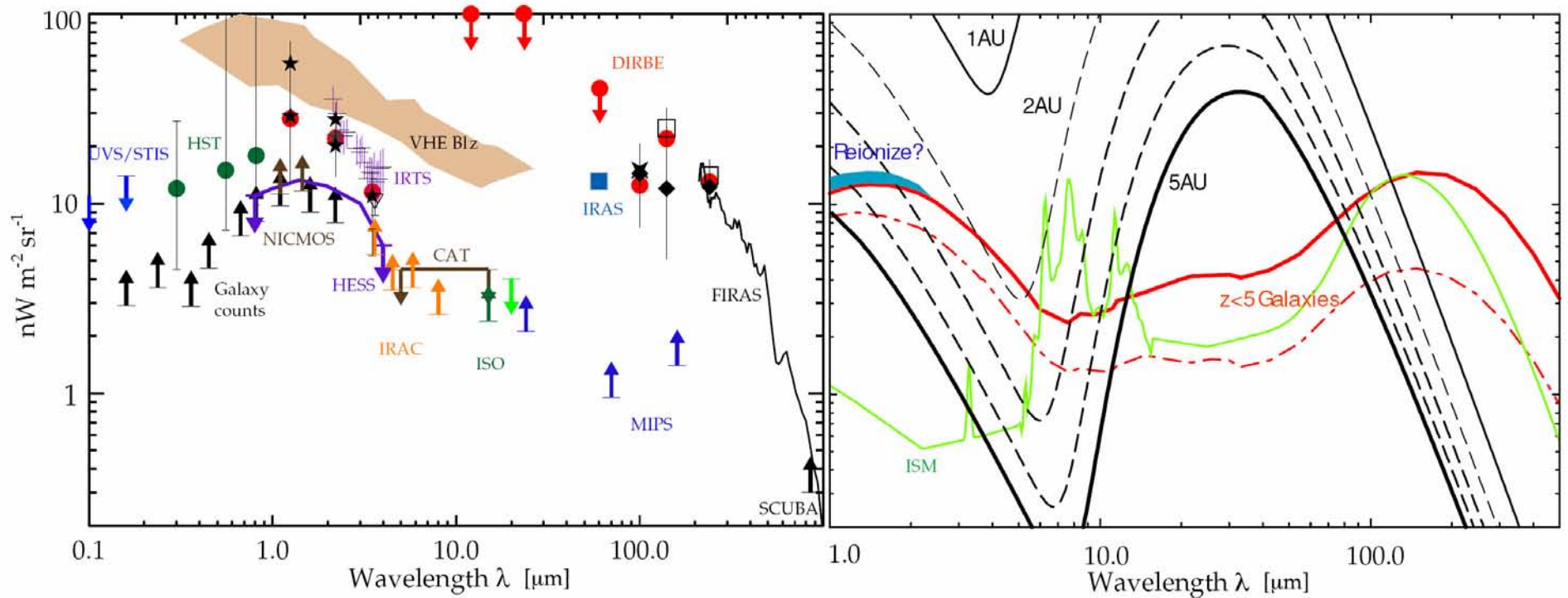
Mean IGM transmission due to Lyman-line resonance scattering at the end of the EOR : the Gunn-Peterson Effect at $z > 6$

- Simulations predict the Ly- α , Ly- β , and Ly- γ opacity of the IGM and its evolution during the EOR \rightarrow compare with the absorption spectra of observed high-redshift quasars to test the theory and the efficiencies assumed for the release of ionizing photons by early galaxies.
- e.g. for this illustrative simulation, EOR ended a bit too early to match the data from Fan et al. (2006) \rightarrow predicts somewhat higher transmission than observed, but captures the observed trend with redshift well \rightarrow lower source efficiencies are required for a better fit.
- Higher- z data can constrain reionization parameters better.



The Cosmic Near-Infrared Background: Fluctuations from the Epoch of Reionization

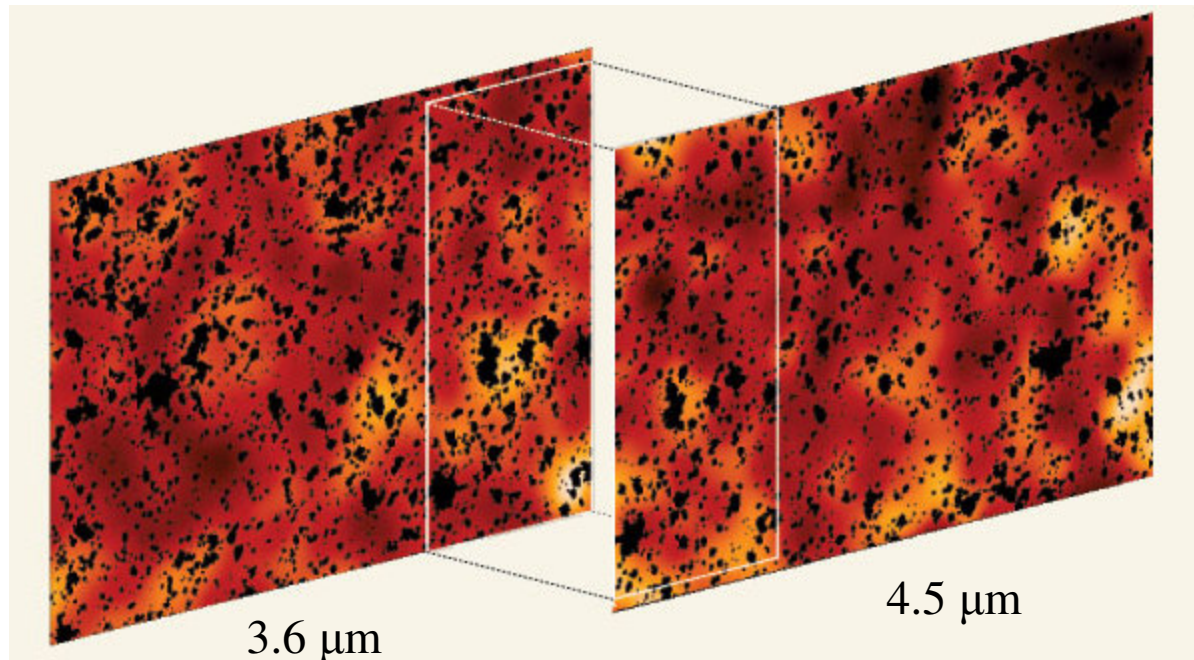
Fernandez, E., Komatsu, E., Iliev, I. T. & Shapiro, P. R. 2009, ApJ, submitted (astro-ph/0906.4552)



Cooray et al. (2009)

The Cosmic Near-Infrared Background: Fluctuations from the Epoch of Reionization

Fernandez, E., Komatsu, E., Iliev, I. T. & Shapiro, P. R. 2009, *ApJ*, submitted (astro-ph/0906.4552)



Near-Infrared Images of the sky in Hubble Deep Field North, by NASA's *Spitzer* Satellite, in two, partially-overlapping fields of view, with point sources removed and regions near bright sources masked

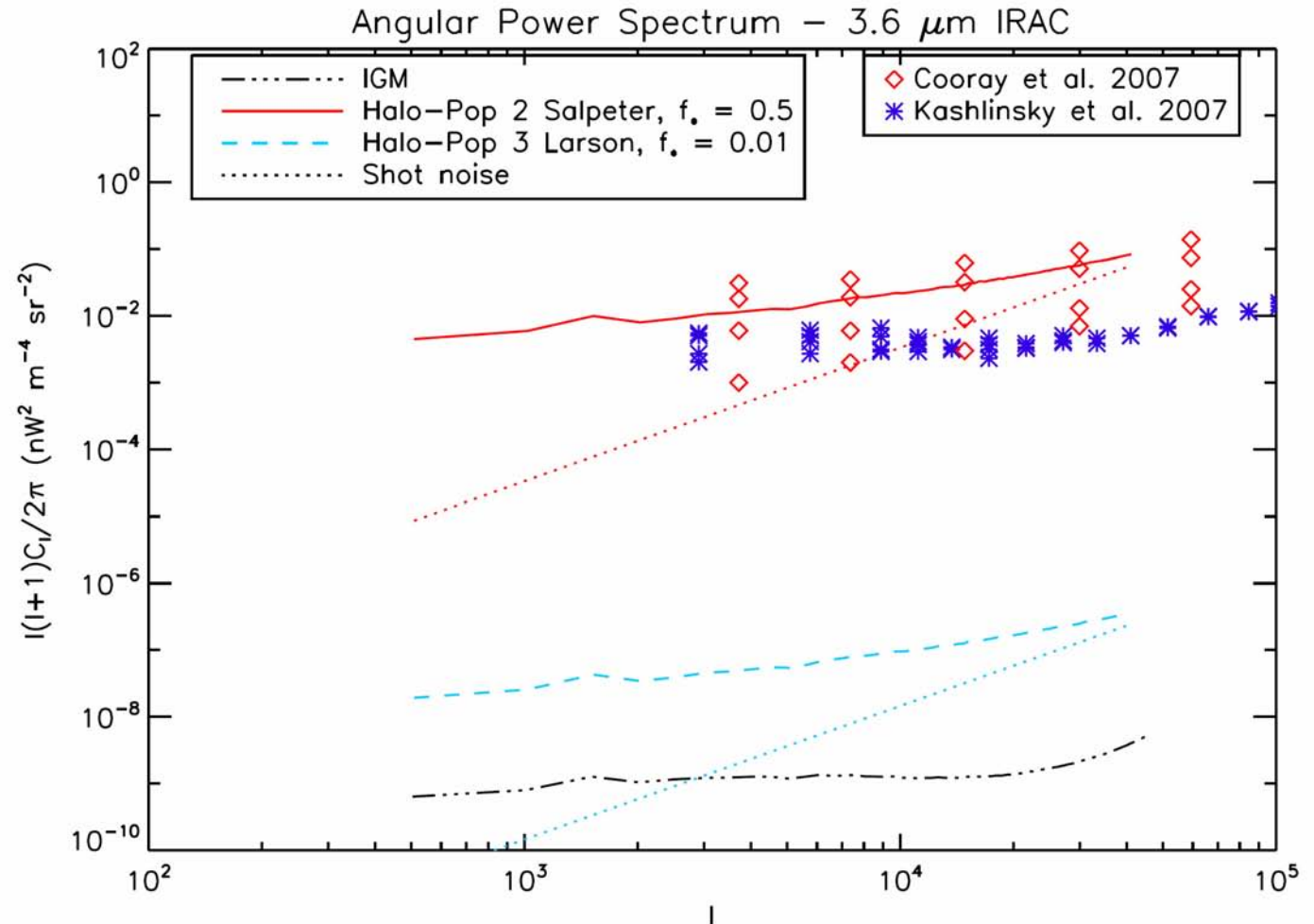
Kashlinsky et al. (2006)

The Cosmic Near-Infrared Background: Fluctuations from the Epoch of Reionization

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- For a given reionization history, the product of f_* and f_{esc} are fixed.

- NIRB fluctuations are maximized by high f_* and low f_{esc}

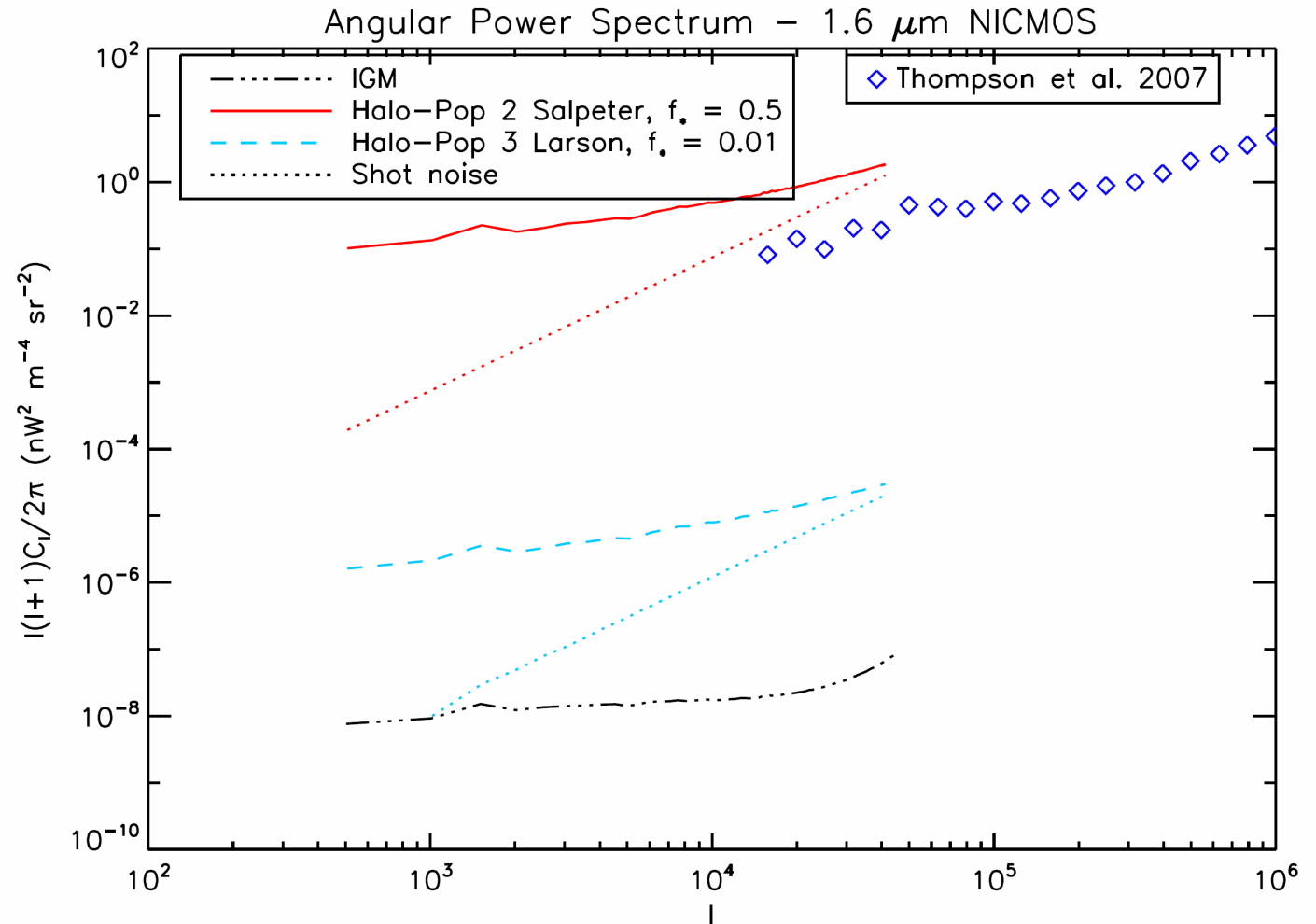


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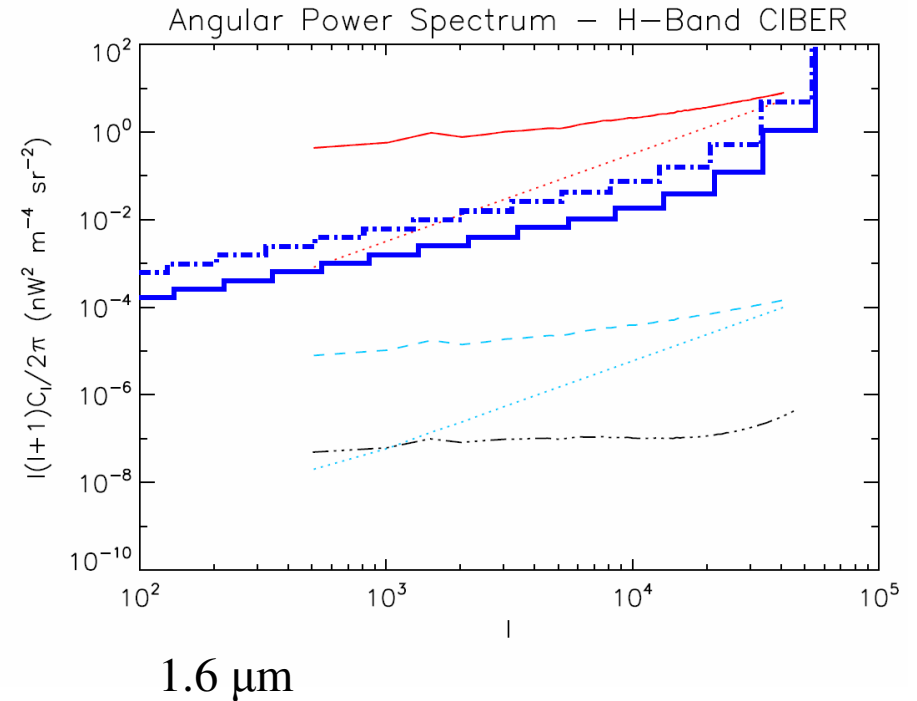
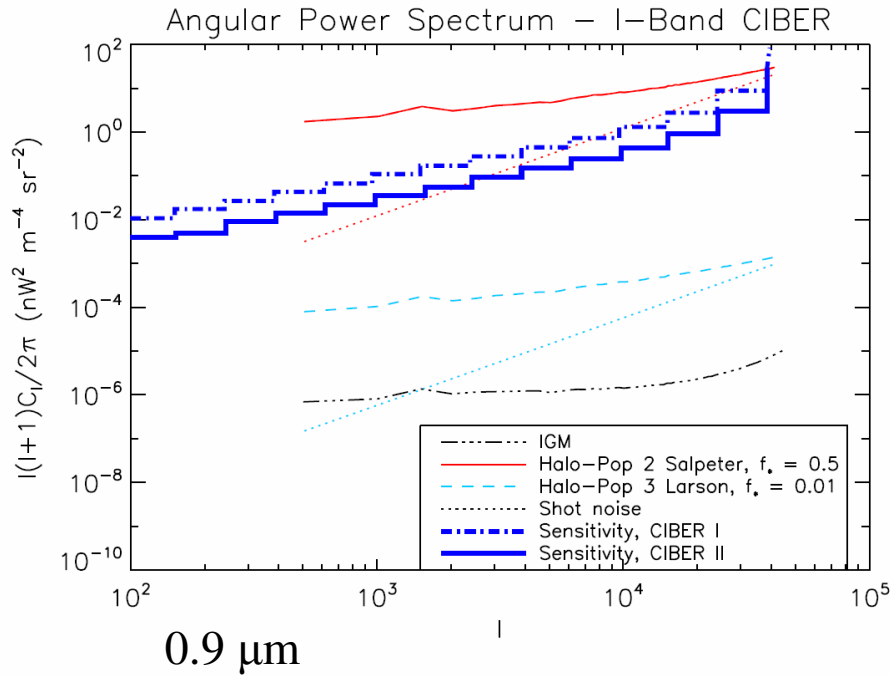
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The Cosmic Near-Infrared Background: Fluctuations from the Epoch of Reionization

Fernandez, E., Komatsu, E., Iliev, I. T. & Shapiro, P. R. 2009, ApJ, submitted
(astro-ph/0906.4552)



CIBER sensitivity from
Cooray et al. (2009)

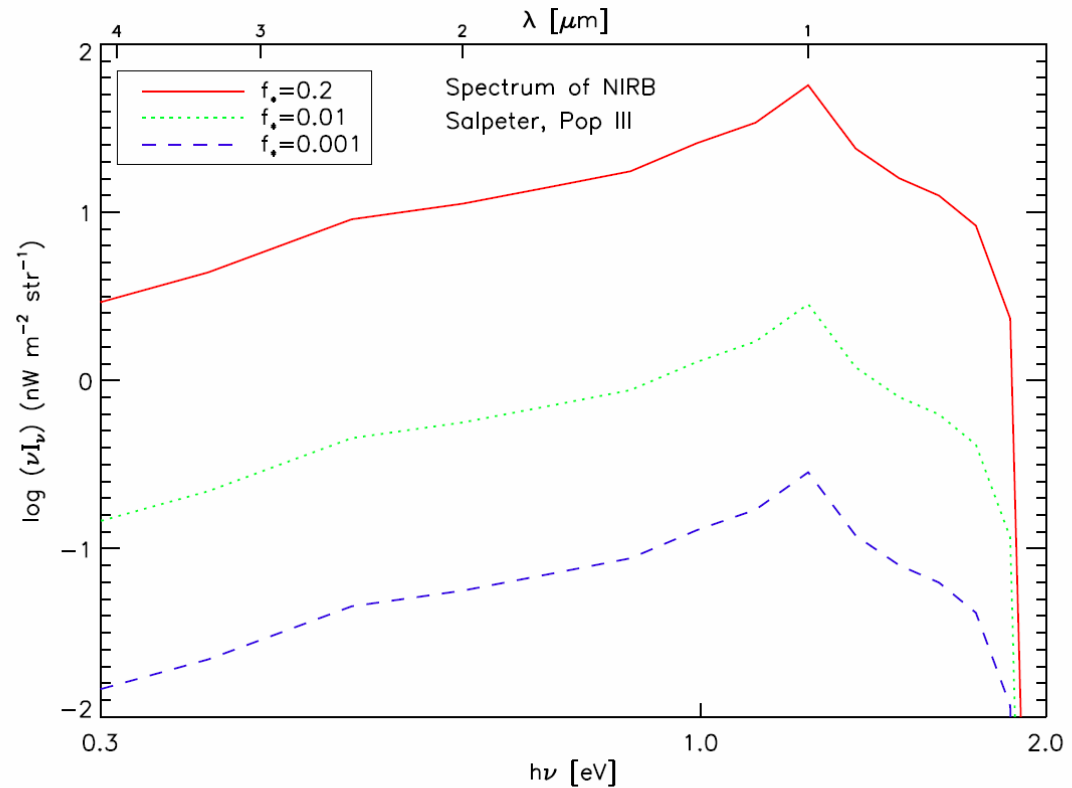
- a range of possible reionization efficiencies makes the NIRB background fluctuations detectable by future experiments like CIBER

The Cosmic Near-Infrared Background: Fluctuations from the Epoch of Reionization

Fernandez, E., Komatsu, E., Iliev, I. T. & Shapiro, P. R. 2009, ApJ, submitted (astro-ph/0906.4552)

if mean NIRB intensity
< 3 nW m⁻² str⁻¹
(Thompson et al. 2007)

====> high star
formation efficiency
(e.g. $f_* > 0.2$) makes
mean NIRB too high,
regardless of NIRB
fluctuation constraints



Self-Regulated Reionization

Iliev, Mellema, Shapiro, & Pen (2007), MNRAS, 376, 534; (astro-ph/0607517)

- Jeans-mass filtering →

low-mass source halos

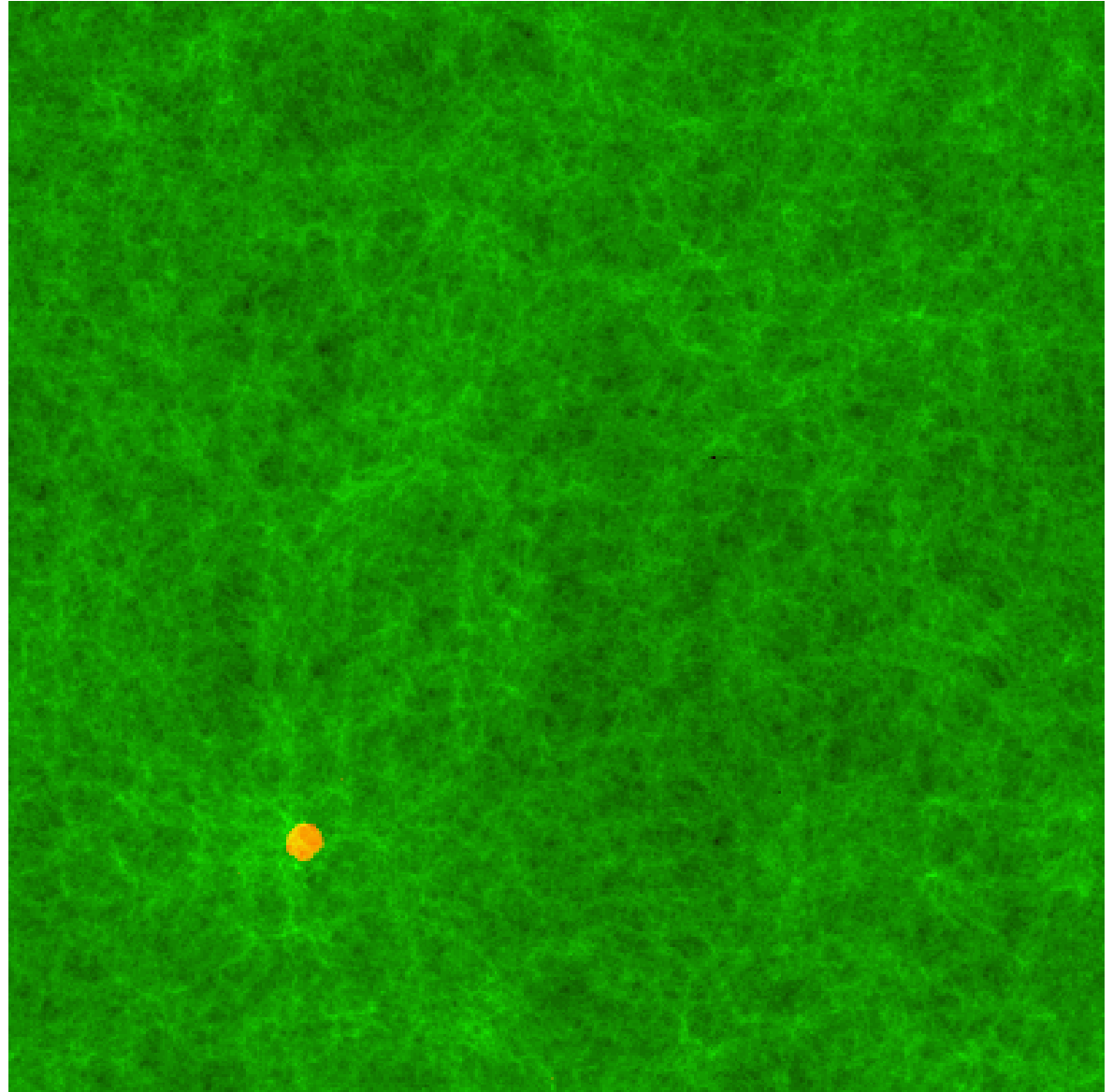
($M < 10^9 M_{\text{solar}}$) cannot form
inside H II regions ;

- 35/h Mpc box, 406^3 radiative
transfer simulation, WMAP3,
 $f_{\gamma} = 250$;

- resolved all halos with

$M > 10^8 M_{\text{solar}}$ (i.e. all
atomically-cooling halos),
(blue dots = source cells);

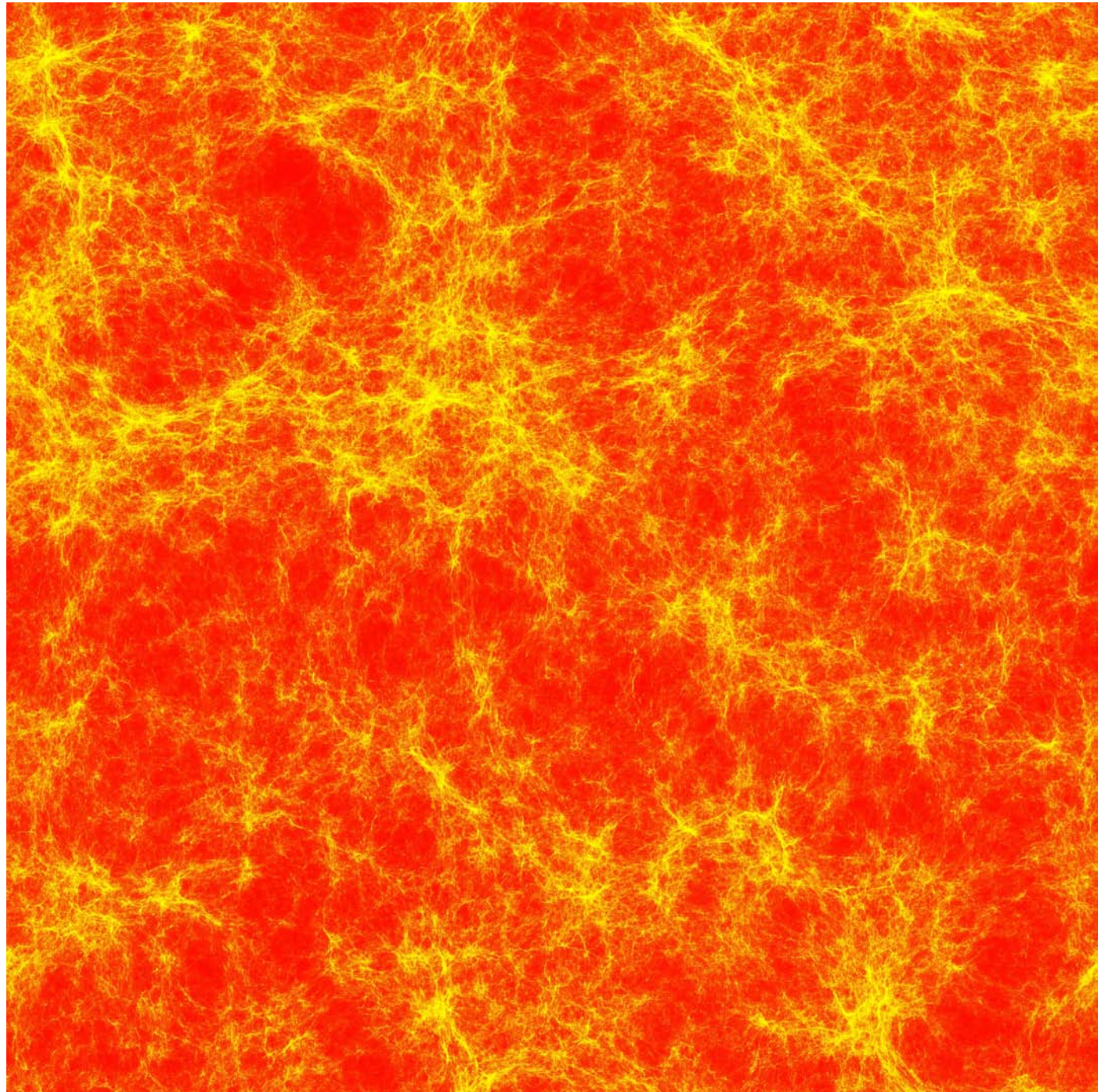
- Evolution: $z=21$ to $z_{\text{ov}} = 7.5$.



New, Large-Scale Simulations of Self-Regulated Reionization

Iliev, Mellema, Pen,
Shapiro, and Merz
(2008), in press
(astro-ph/0806.2887);
Shapiro, Iliev,
Mellema, Pen, &
Merz (2008), AIP
1035, 68
(astro-ph/0806.3091)

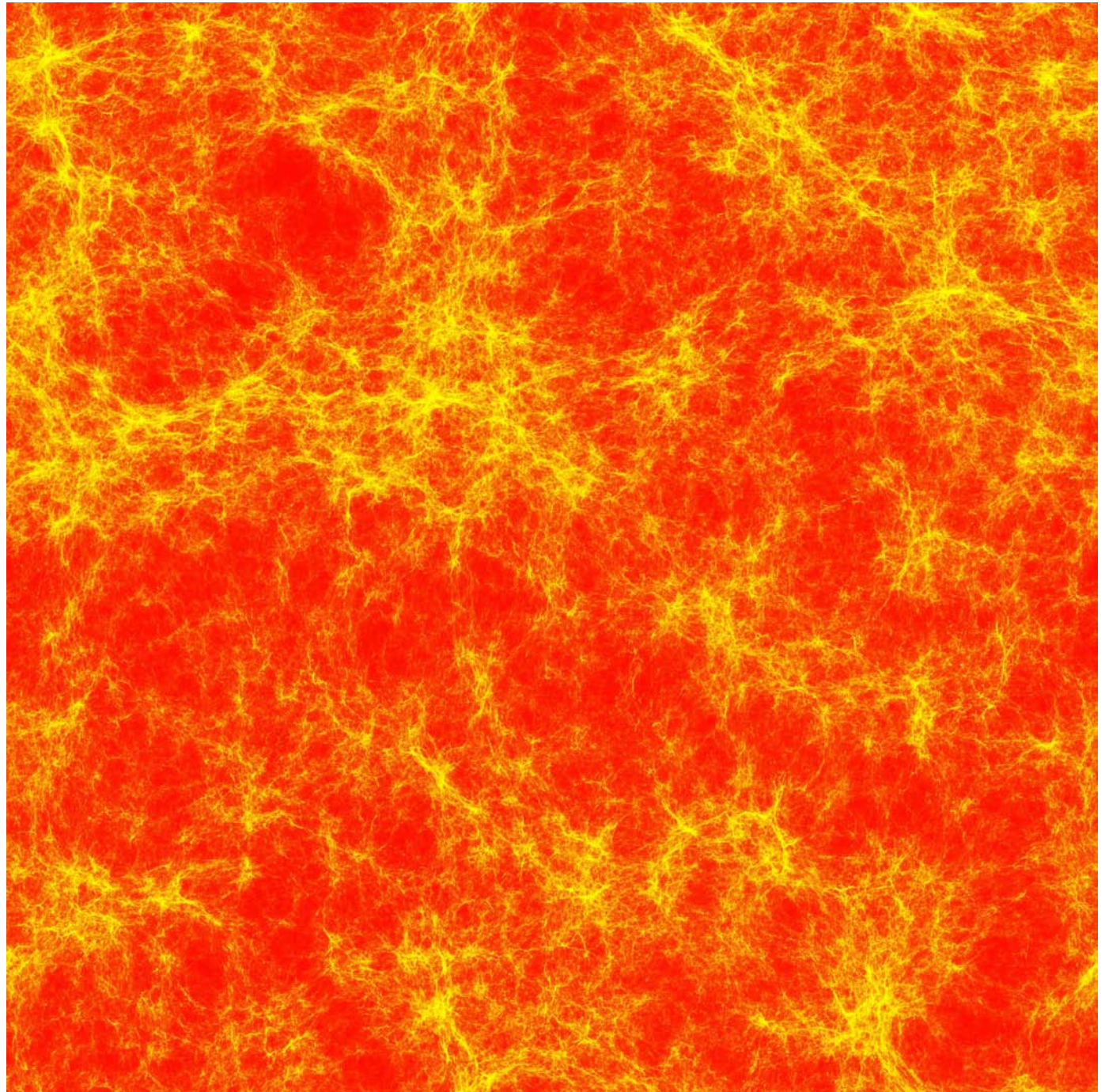
CubeP³M N-body
 Λ CDM sim with
 3072^3 (29 billion)
particles,
 6144^3 cells,
box size = 160 Mpc;
particle mass =
5 million solar
masses



New, Large-Scale Simulations of Self-Regulated Reionization

Iliev, Mellema, Pen,
Shapiro, and Merz
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CubeP³M N-body
 Λ CDM sim with
 3072^3 (29 billion)
particles,
resolves halos above
 10^8 solar masses

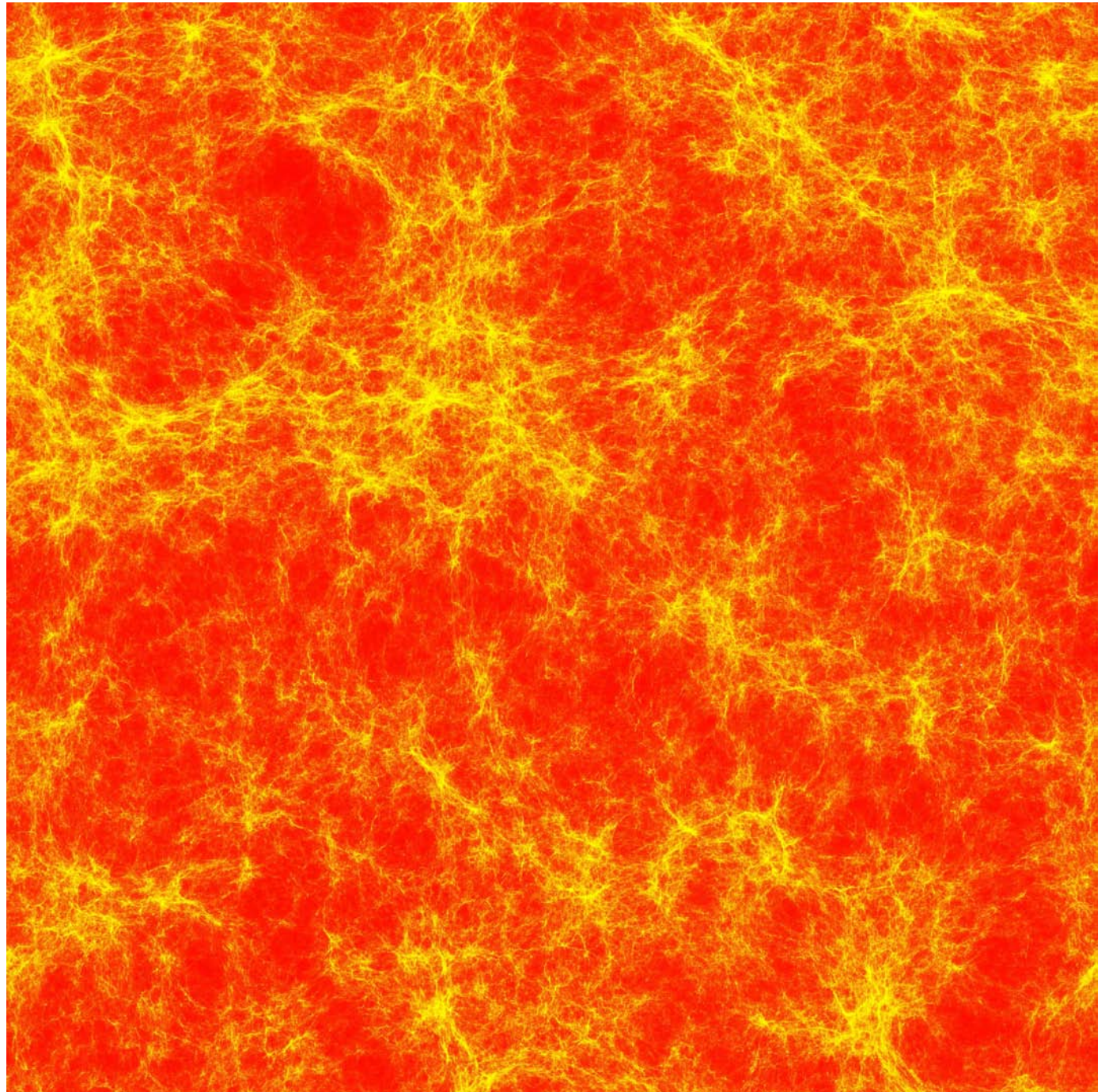


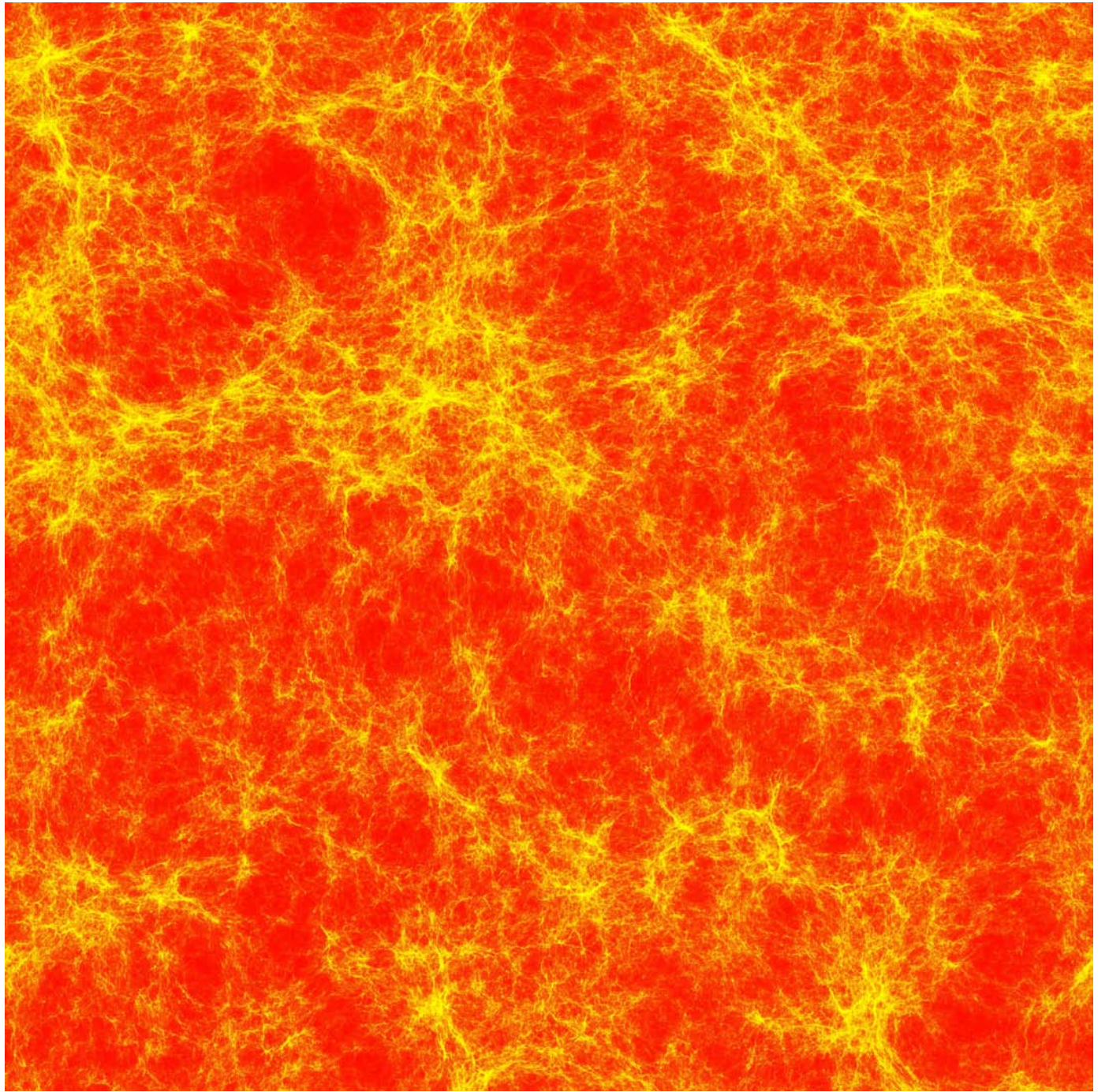
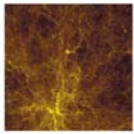
New, Large-Scale Simulations of Self-Regulated Reionization

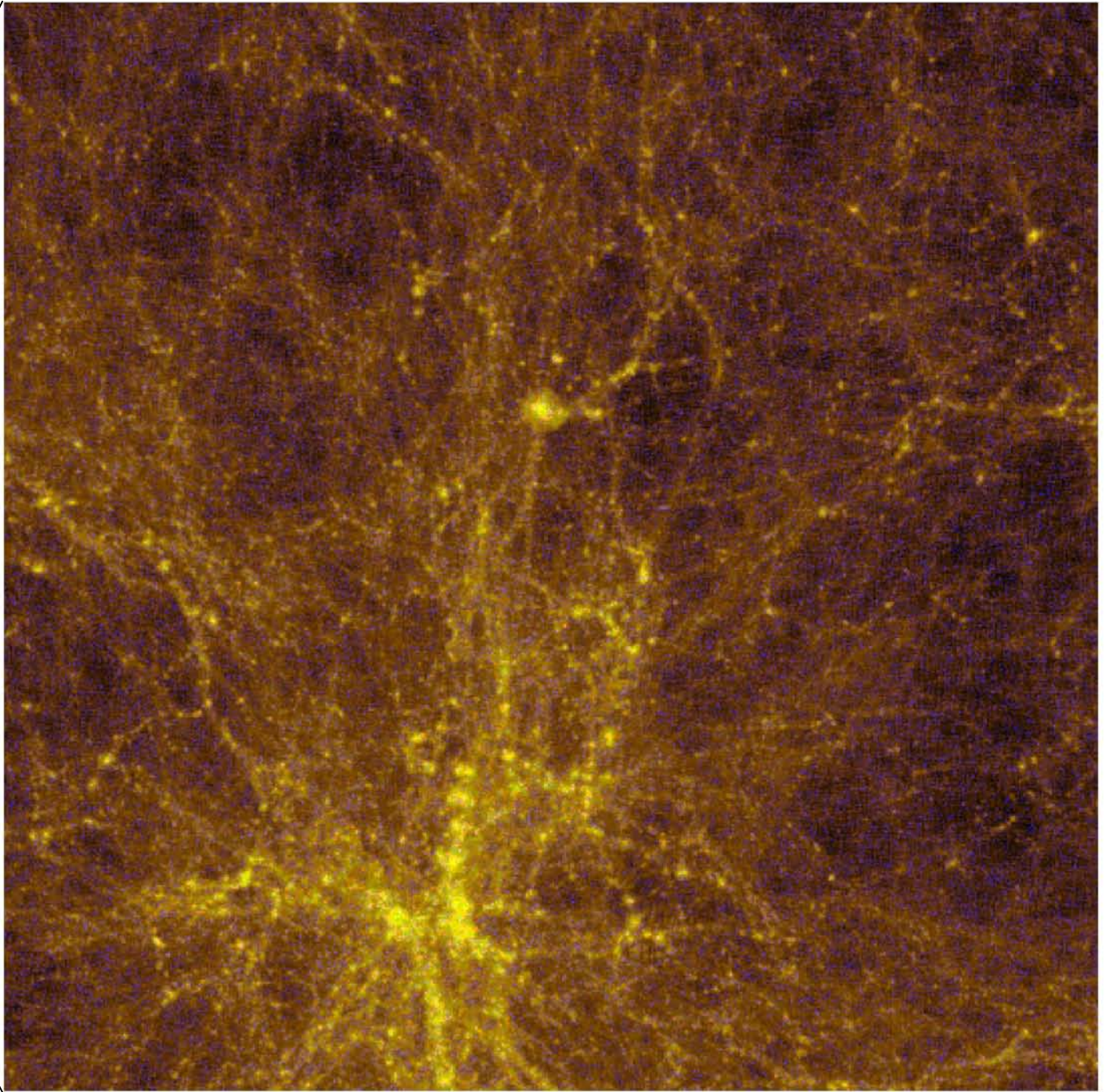
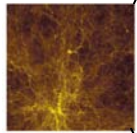
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Shapiro, Iliev,
Mellema, Pen, &
Merz (2008), AIP
1035, 68
(astro-ph/0806.3091)

CubeP³M N-body

- new Texas Sun
Constellation Linux
Cluster, *Ranger*,
2048 cores, 159,000
SUs (cores x hours)

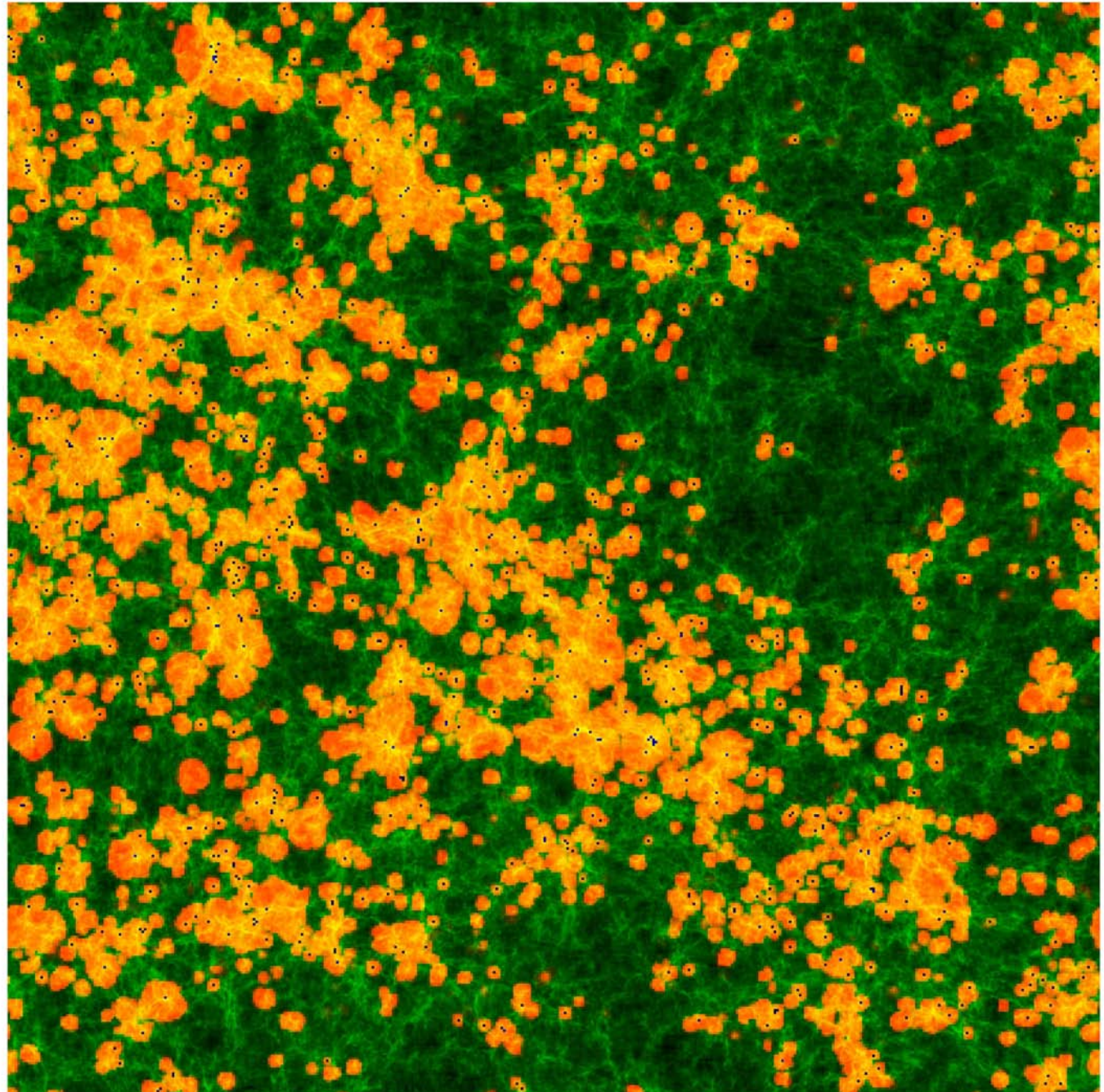






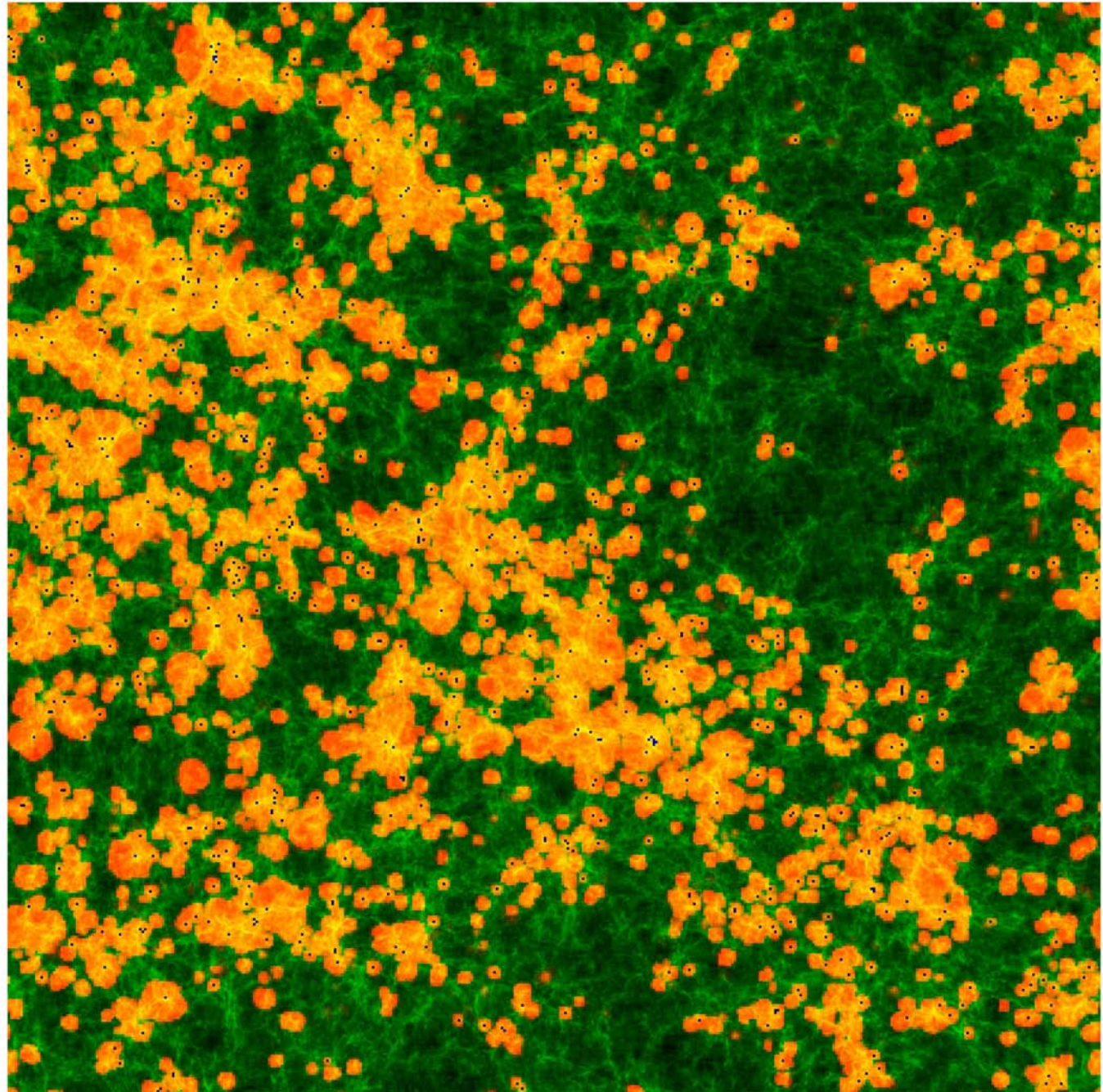
C²Ray radiative transfer

- RT grid 432^3
cells
- box size = 90
Mpc



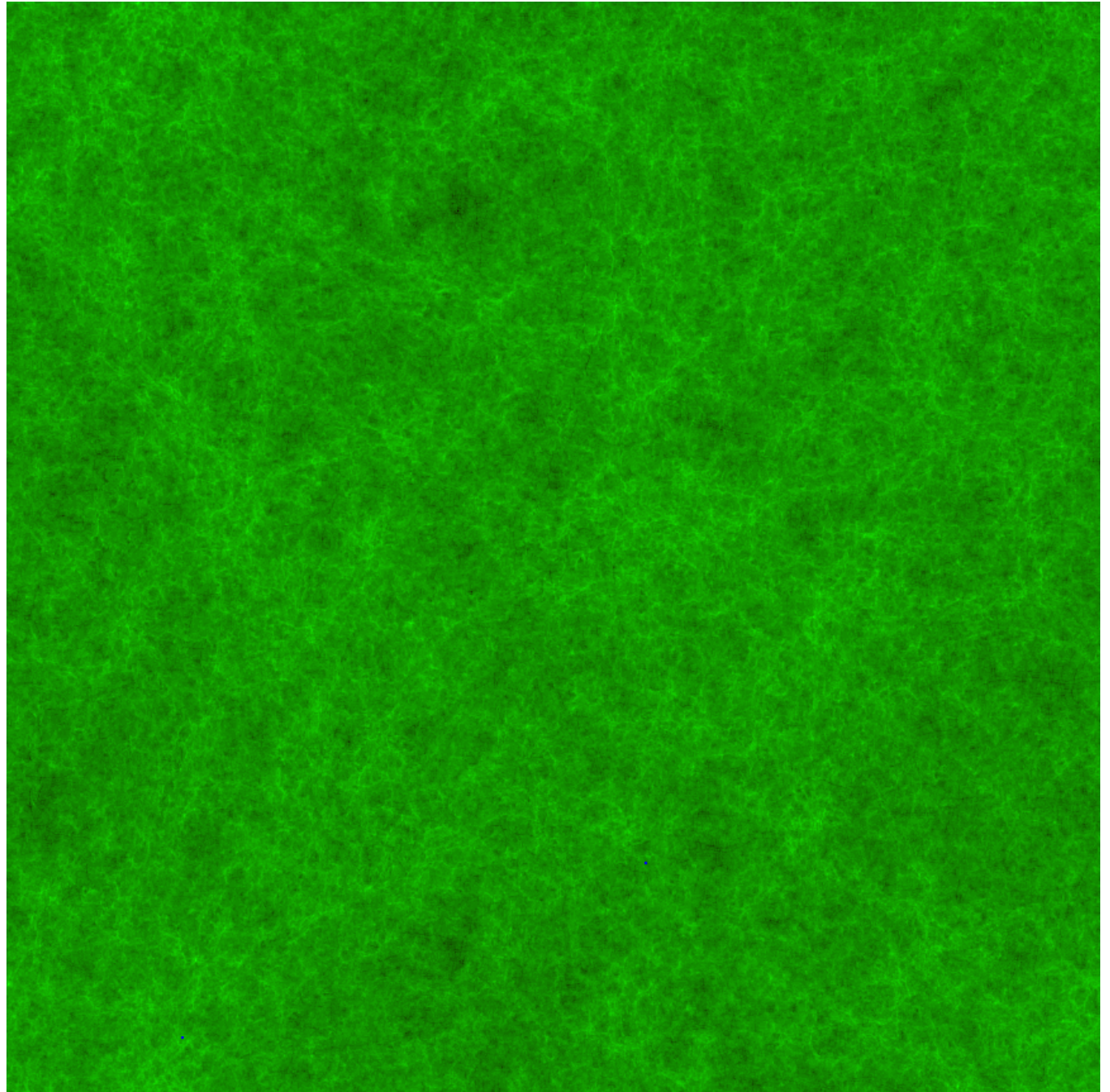
C²Ray radiative transfer

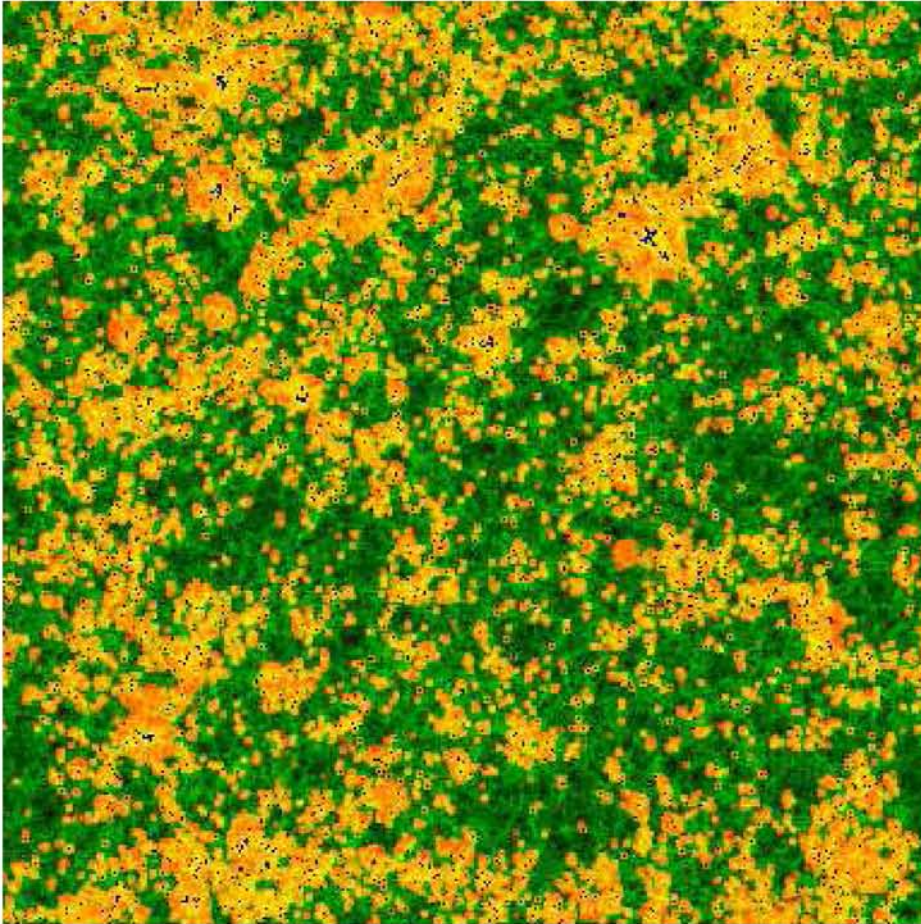
- RT grid 432^3
cells
- box size = 90
Mpc
- *Ranger*,
Texas Sun
Constellation
Linux Cluster,
700,000 SUs
(cores x hours),
up to 10,000
cores



Self-Regulated Reionization in Λ CDM

- 90 Mpc box

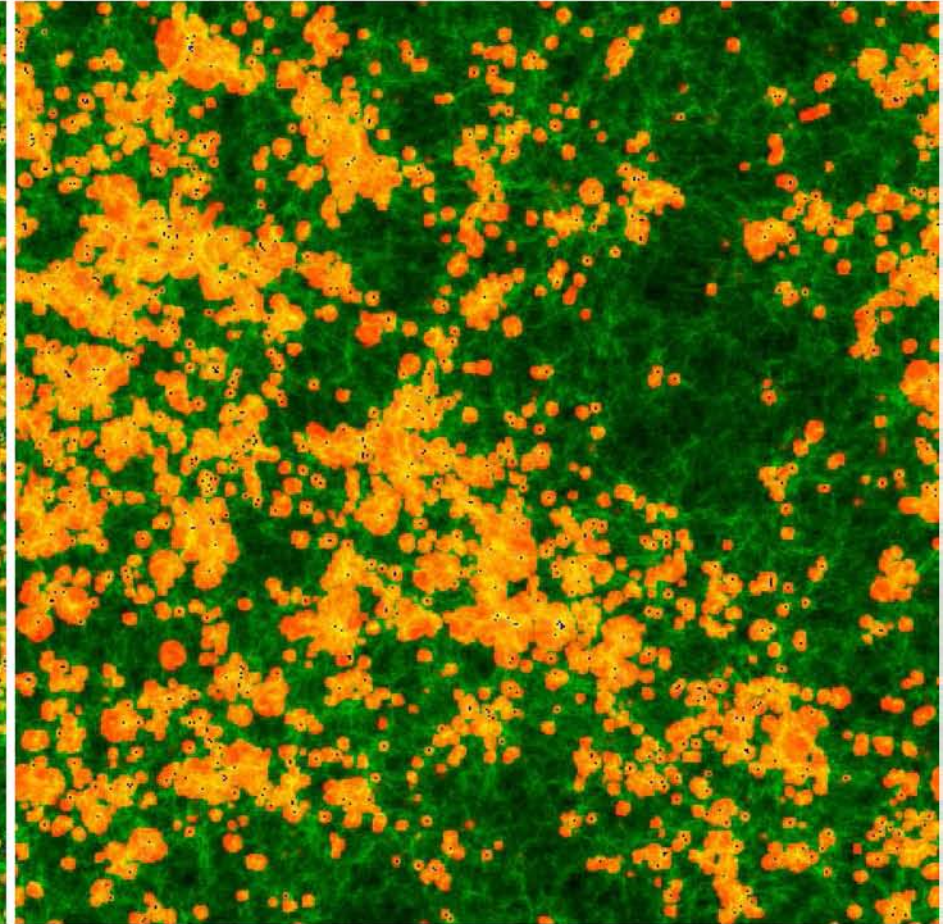




160 Mpc box

$$z = 11.6$$

when mass-weighted mean ionized fraction of universe $x_m = 0.3$

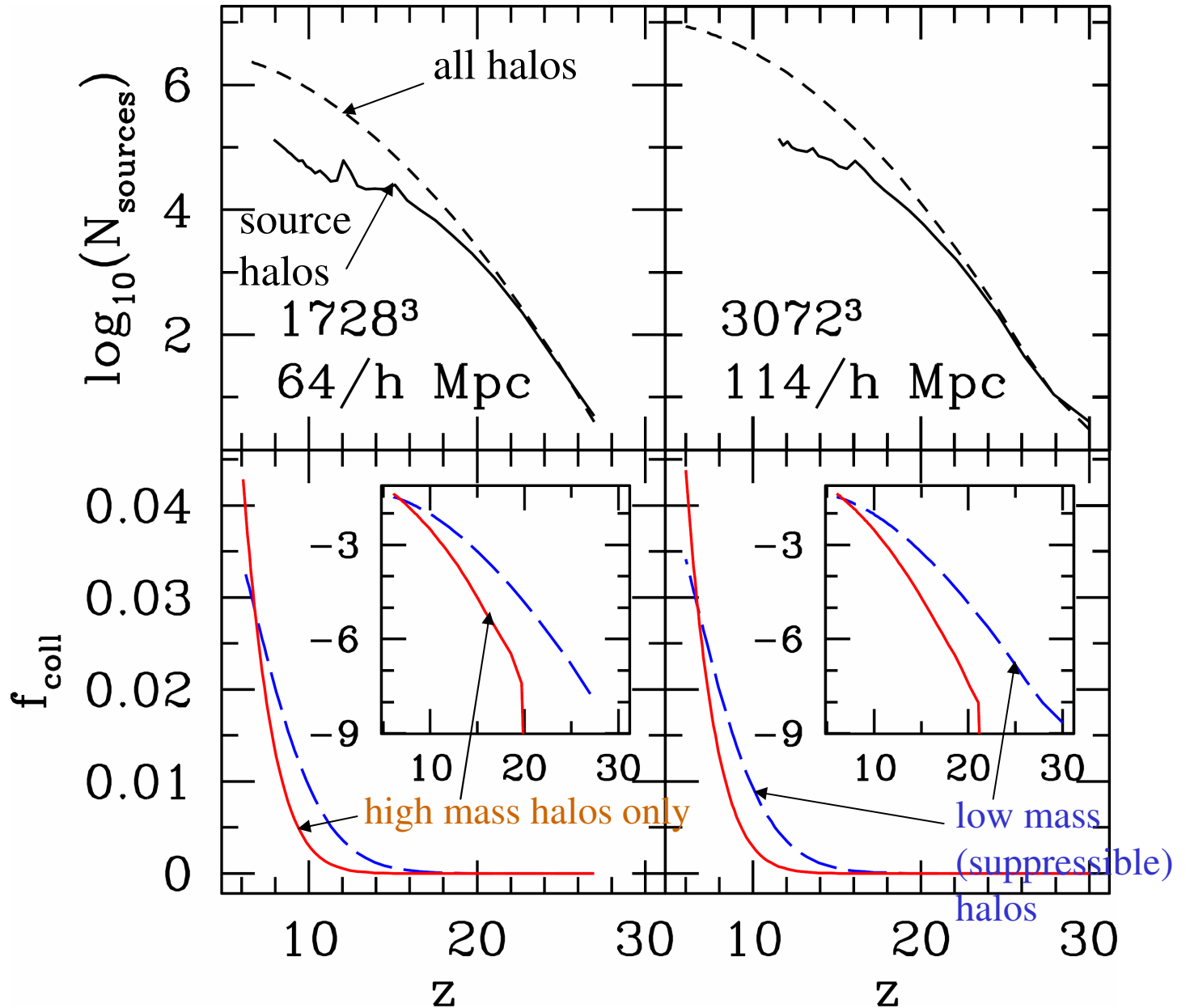


90 Mpc box

$$z = 11.9$$

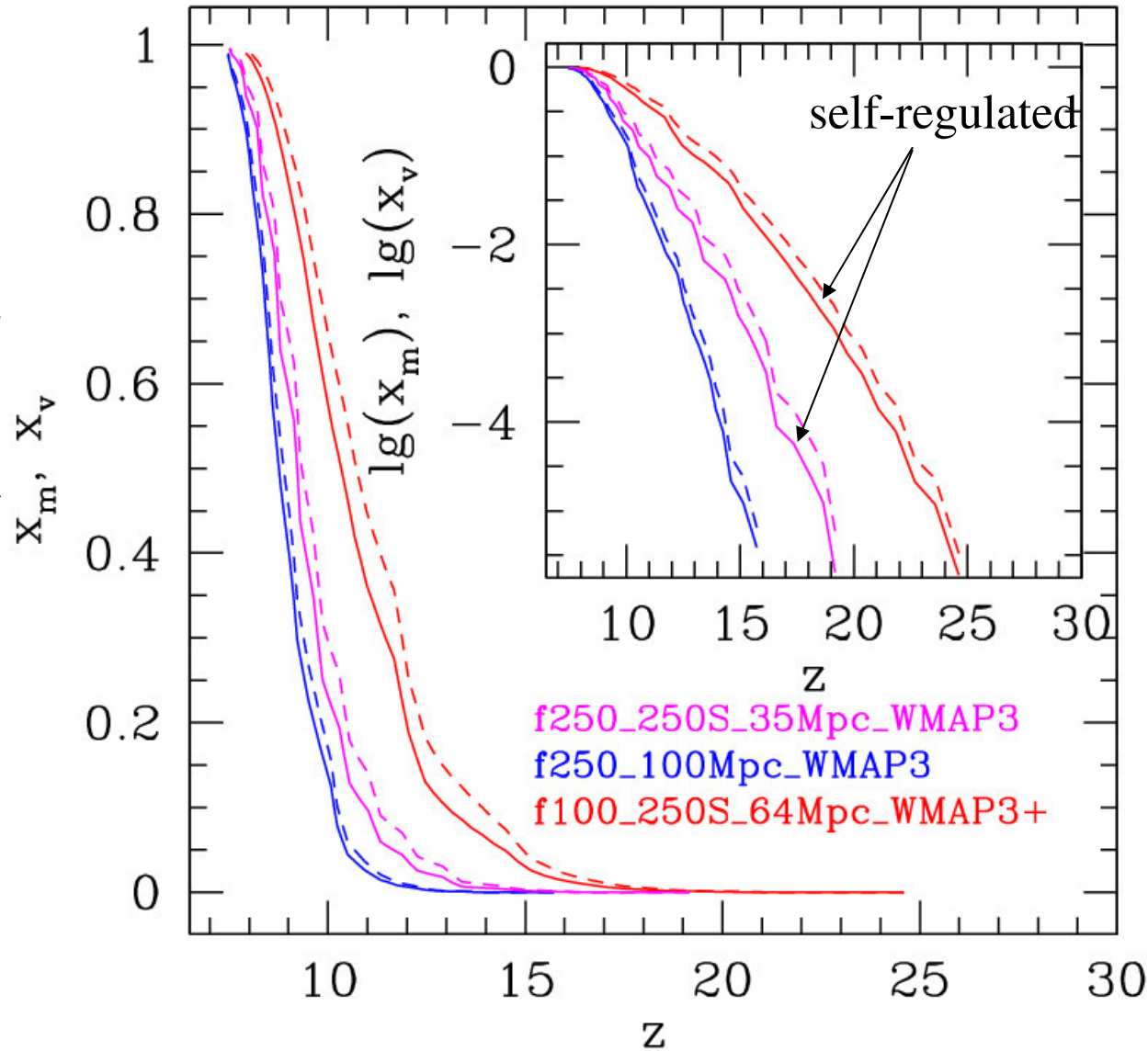
Self-regulated halo mass function

- Jeans-mass filtering suppresses formation of sources in small-mass halos which form inside H II regions
- clustering of small-mass halos around density peaks enhances this effect → suppression is strongly **biased**



Evolution of the Mean Ionized Fraction of the Universe

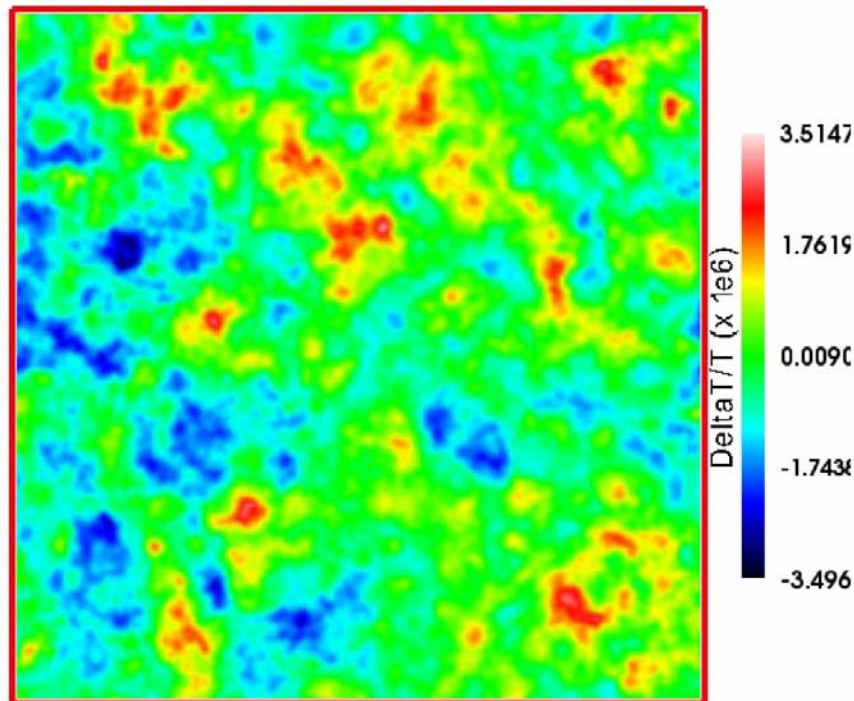
- self-regulated reion. (i.e. small-mass halos resolved) starts earlier, but ends about the same time →
- high-mass halos dominate the end of EOR



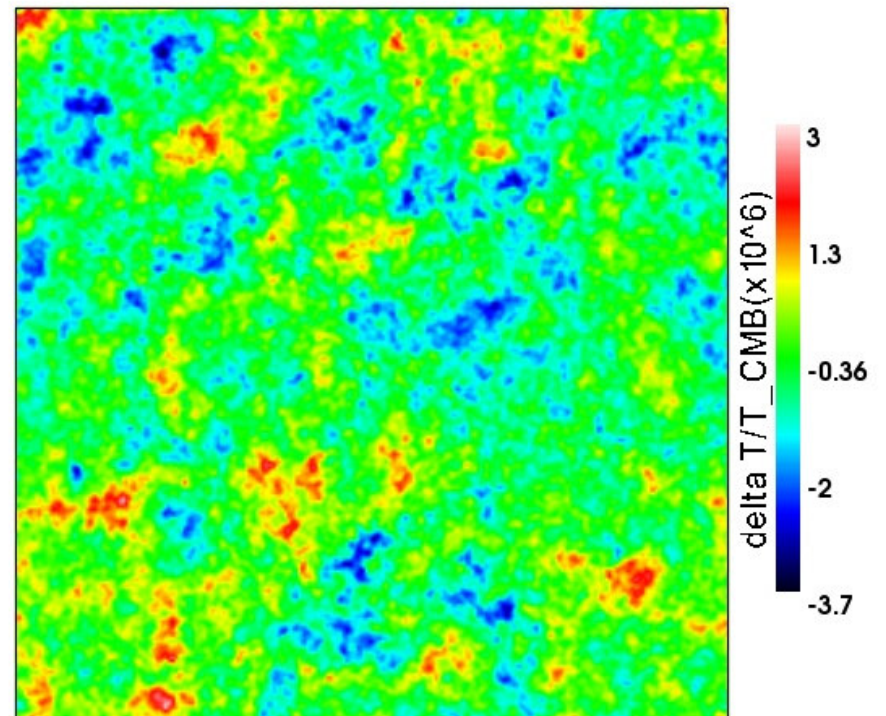
$(\delta T/T_{\text{CMB}})$ Maps of the Kinetic Sunyaev-Zel'dovich Effect from Radiative Transfer Simulations of Patchy Reionization: Effect of Self-Regulation

- Box size 100/h Mpc comoving
→ 50' x 50'
- Source halos $> 10^9$ solar masses

- Box size 114/h Mpc comoving
→ $1^\circ \times 1^\circ$
- Source halos $> 10^8$ solar masses
- Self – Regulated Reionization



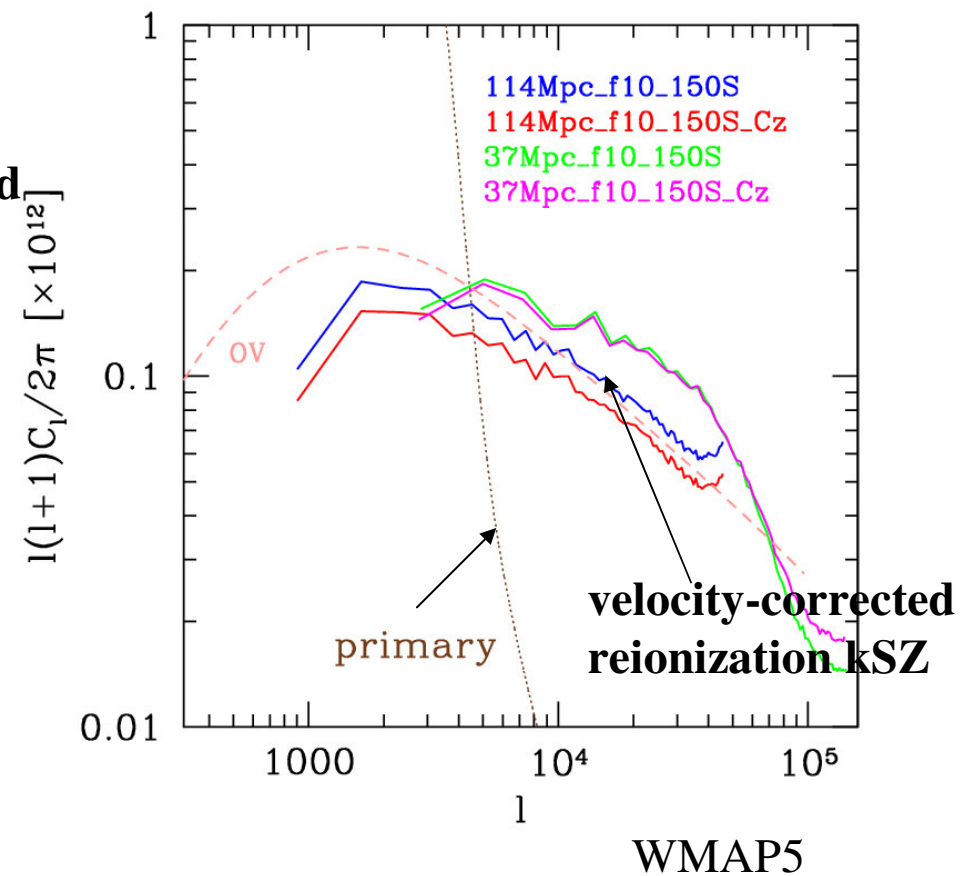
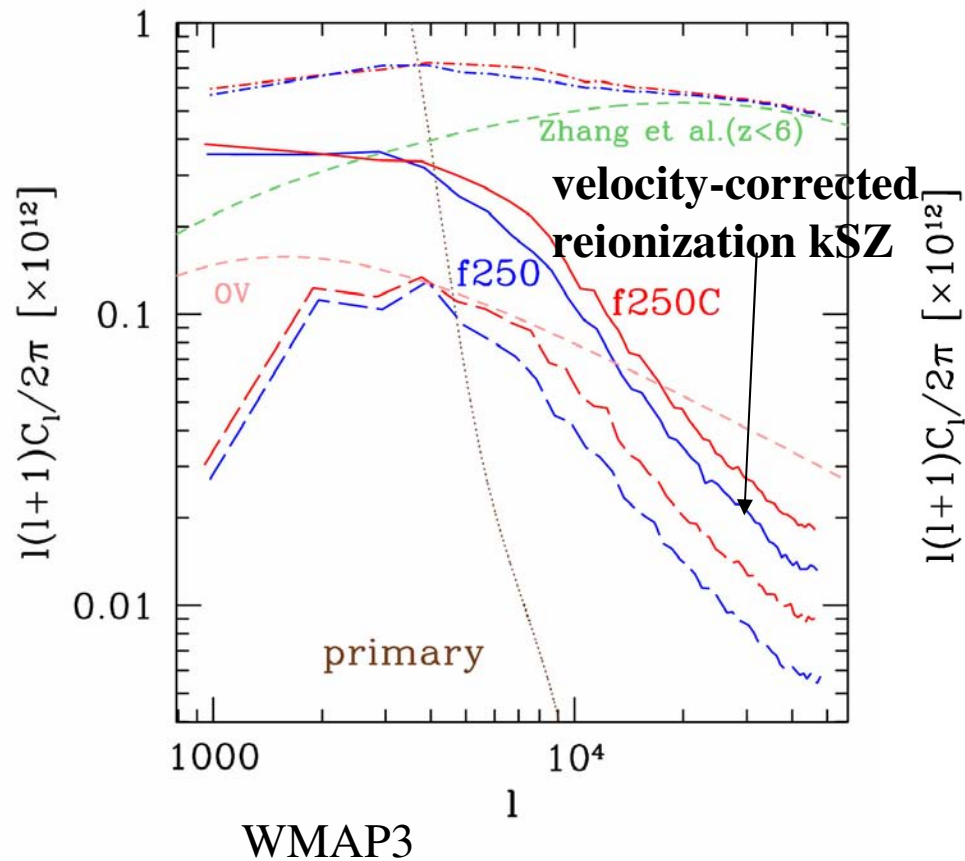
WMAP3



WMAP5

kSZ CMB Anisotropy Signal: Sky Power Spectra of $\delta T_{\text{kSZ}} / T_{\text{CMB}}$: Effect of Self-Regulation

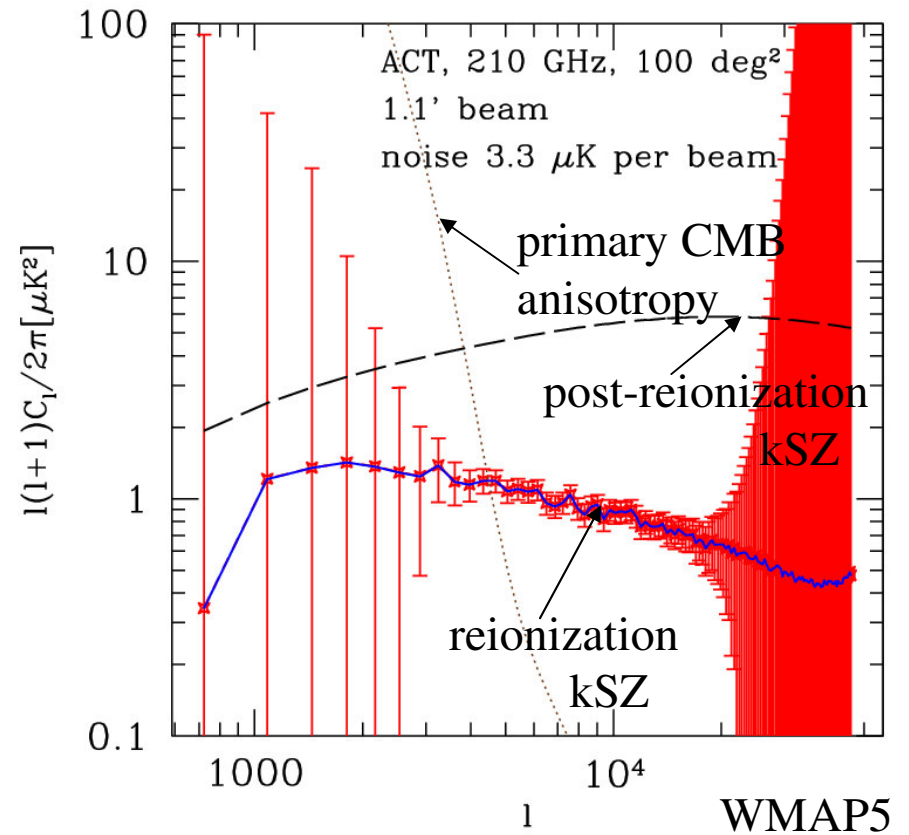
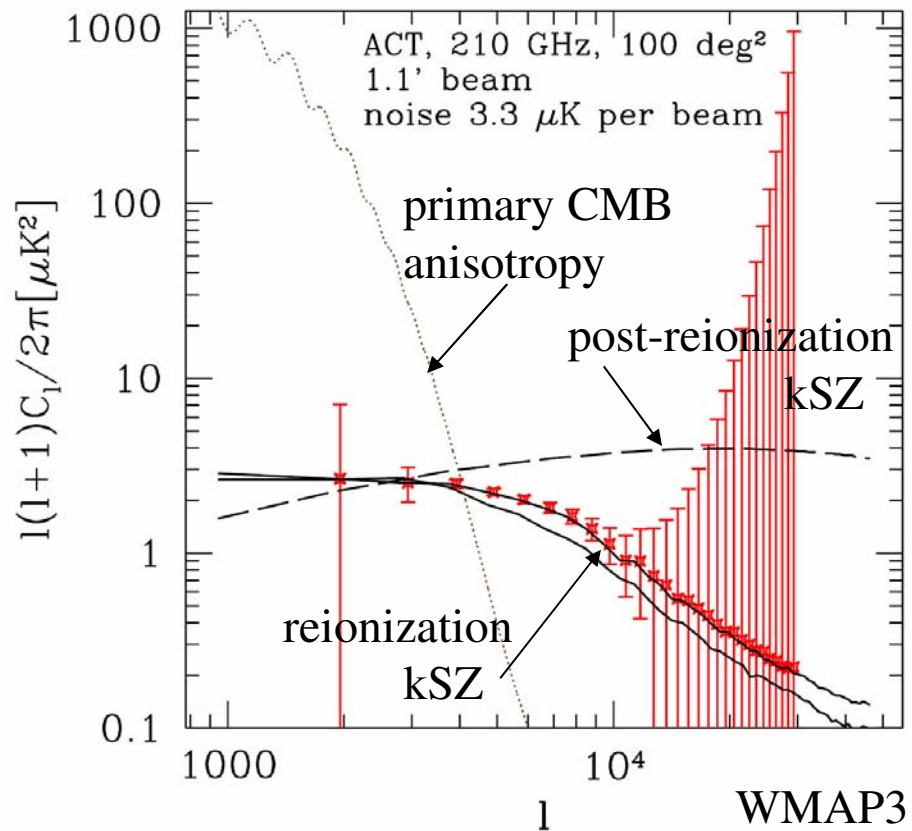
- Source halos $> 10^9$ solar masses
- Source halos $> 10^8$ solar masses
- Self-regulated reionization



Observability of the kSZ from reionization: sky power spectrum

Effect of Self-Regulation

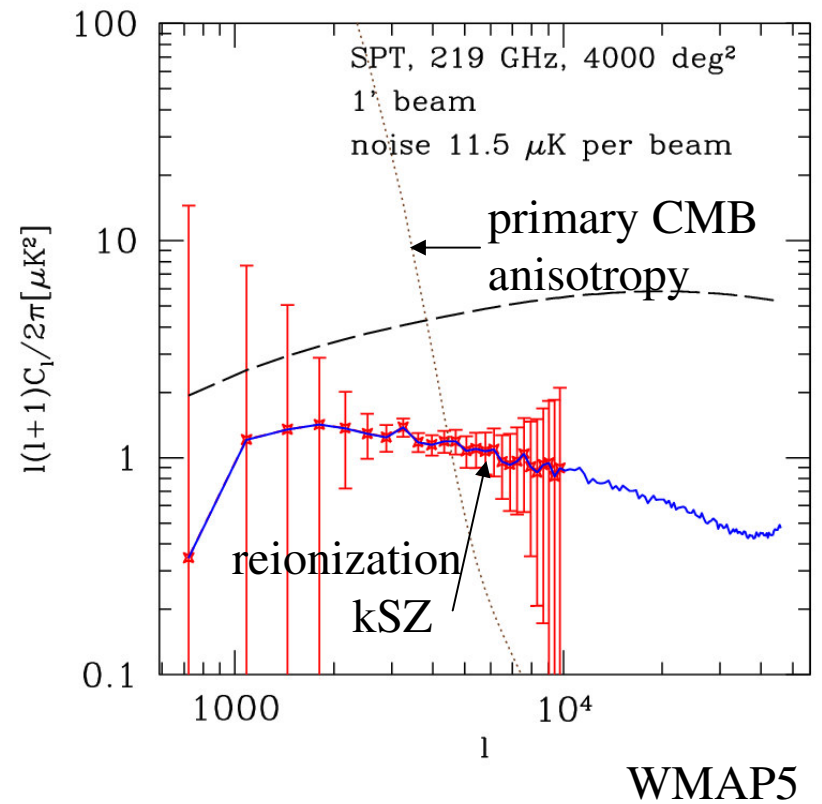
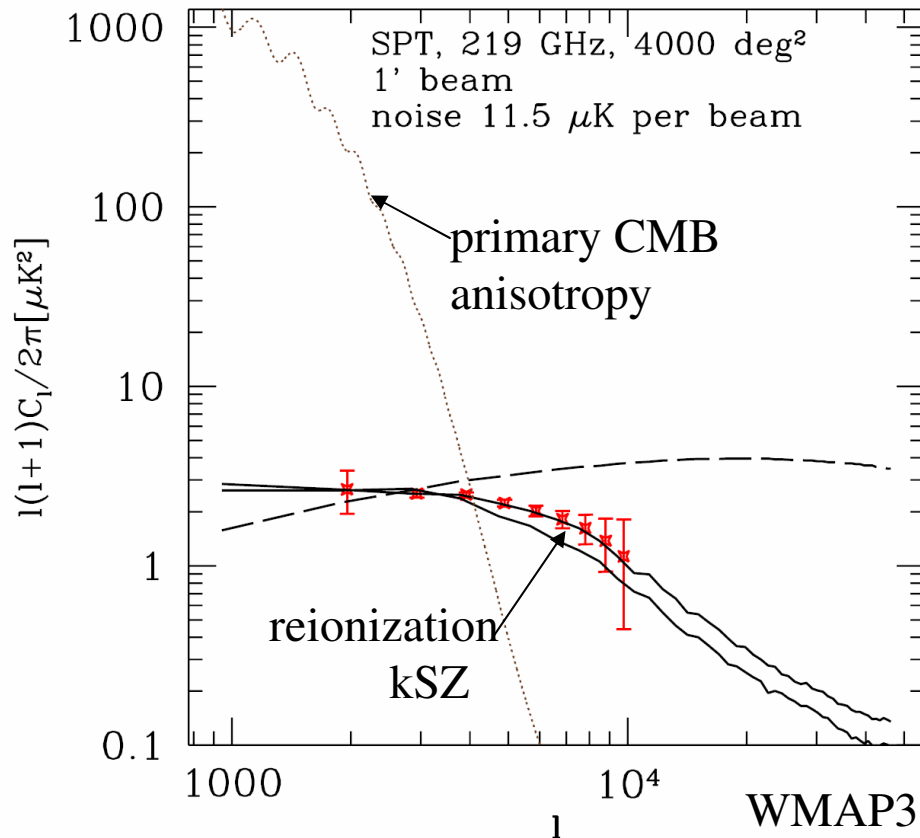
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- Source halos $> 10^8$ solar masses
- Self-regulated reionization



Observability of the kSZ from reionization: sky power spectrum

Effect of Self-Regulation

- Source halos $> 10^9$ solar masses
- Source halos $> 10^8$ solar masses
- Self-regulated reionization



Low Frequency Array (LOFAR)

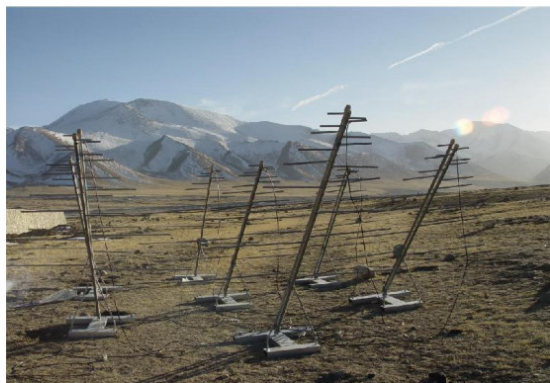


Murchison Widefield Array (MWA)



Primeval Structure Telescope (PAST/21CMA)

Prototype Tests, Ulaanbaatar, Xin Jiang, China



Giant Meterwave Radio Telescope (GMRT)



Square Kilometer Array (SKA)



21-cm Radiation Background

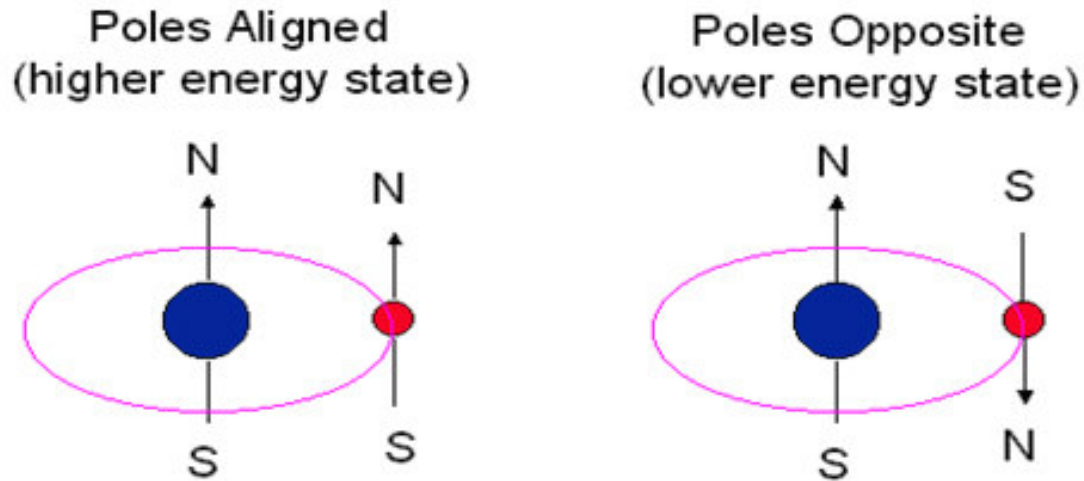
- Foreground emission or absorption by H atoms at redshift z seen against CMB at redshifted wavelength $21(1+z)$ cm.

Emission $\leftrightarrow T_{\text{spin}} > T_{\text{CMB}}$

Absorption $\leftrightarrow T_{\text{spin}} < T_{\text{CMB}}$

Transparent $\leftrightarrow T_{\text{spin}} = T_{\text{CMB}}$

21-cm Level Population of Atomic Hydrogen



$$\frac{n_2}{n_1} = 3 \exp\left(-\frac{h\nu_0}{kT_{spin}}\right)$$

3 Ways to Change the 21-cm Level Population

- An H atom can:
 - Absorb a 21-cm photon from the CMB
(CMB Pumping)
 - Collide with another atom (or an electron or ion)
(Collisional Pumping)
 - **Absorb a UV photon at 1215 Angstrom to make Lyman- α transition of H atom, then decay to one of 21-cm levels (“Wouthuysen-Field Effect”)**
(Lyman- α Pumping)

Stages of 21-cm Background

- Dark Ages
 - $z \geq 150$, $T_{\text{spin}} = T_{\text{CMB}} \rightarrow$ nothing
 - $20 \leq z \leq 150$, $T_{\text{spin}} < T_{\text{CMB}} \rightarrow$ absorption
 - $z \leq 20$, $T_{\text{spin}} > T_{\text{CMB}}$ in minihalos \rightarrow emission
- **Epoch of Reionization ($6 \leq z \leq 20$)**
 - $T_{\text{spin}} > T_{\text{CMB}}$ in minihalos \rightarrow emission
 - **After sources turn on, Lyman- α pumping \rightarrow**
 - Without heating, $T_{\text{spin}} < T_{\text{CMB}} \rightarrow$ IGM in absorption
 - **With heating, $T_{\text{spin}} > T_{\text{CMB}} \rightarrow$ IGM in emission**

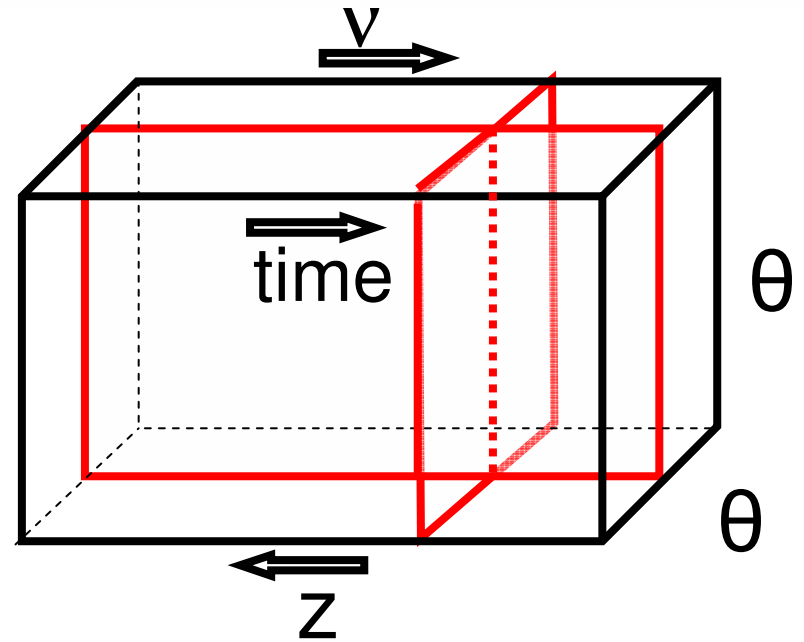
The Redshifted 21cm Signal From the EoR

- The measured radio signal is the **differential brightness temperature**

- $$\delta T_b = T_b - T_{\text{CMB}}: \quad \delta T_b \approx 25 x_{\text{HI}}(1 + \delta) \left(\frac{1 + z}{10} \right)^{1/2} \left[1 - \frac{T_{\text{CMB}}(z)}{T_s} \right] \left[\frac{H(z)/(1 + z)}{dv_{\parallel}/dr_{\parallel}} \right] \text{ mK}$$

(for WMAP3 cosmological parameters).

- Depends on:
 - x_{HI} : neutral fraction
 - δ : overdensity
 - T_s : spin temperature
- For $T_s \gg T_{\text{CMB}}$, the dependence on T_s drops out
- The signal is a spectral *line*: carries **spatial, temporal, and velocity information**.



The image cube: images stacked in frequency space

$$\delta T_b \approx 25 x_{\text{HI}}(1 + \delta) \left(\frac{1+z}{10} \right)^{1/2} \left[1 - \frac{T_{\text{CMB}}(z)}{T_s} \right] \left[\frac{H(z)/(1+z)}{dv_{\parallel}/dr_{\parallel}} \right] \text{ mK}$$

- For $T_s \gg T_{\text{CMB}}$, the dependence on T_s drops out and brightness temperature fluctuations then depend only on L.O.S. velocity gradient & H I density fluctuations,

$$\delta_{\rho_{\text{HI}}} = (1 + \delta_{\rho})(1 + \delta_{x_{\text{HI}}}) - 1 = \delta_{\rho} + \delta_{x_{\text{HI}}} + \delta_{\rho}\delta_{x_{\text{HI}}}$$

- For *linear* perturbations in the matter density, the *peculiar* velocity gradient and density perturbation are related in Fourier space as follows:

$$(aH)^{-1} \frac{dv_{\parallel}}{dr_{\parallel}}(\mathbf{k}) = -\mu^2 \delta_{\rho}(\mathbf{k})$$

$$\mu = \hat{\mathbf{k}} \cdot \hat{\mathbf{r}} = \text{cosine of angle between the wavevector } \mathbf{k} \text{ and the line of sight}$$

- To lowest order, we can then write in Fourier space,

$$\delta T_b(\mathbf{k}) \approx (\delta T_b)_{\text{average}} (\delta_{\rho_{\text{HI}}} + \mu^2 \delta_{\rho})$$

- the 21cm 3D Power Spectrum can then be decomposed as follows:

$$P_{\Delta T}(\mathbf{k}) = P_{\mu^0}(k) + P_{\mu^2}(k)\mu^2 + P_{\mu^4}(k)\mu^4$$

where

$$P_{\mu^0} = \overline{\delta T_b}^2 P_{\delta_\zeta, \delta_\zeta}$$

$$P_{\mu^2} = 2 \overline{\delta T_b}^2 P_{\delta_\zeta, \delta_\rho} \quad (\text{where } \delta_\zeta \equiv \delta_{\rho_{\text{HI}}})$$

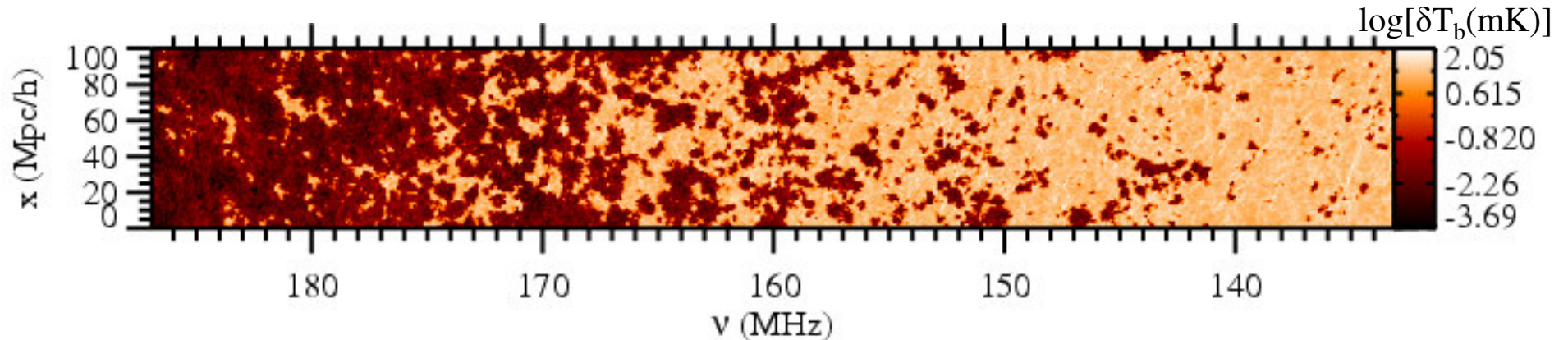
$$P_{\mu^4} = \overline{\delta T_b}^2 P_{\delta_\rho, \delta_\rho}$$

- In that case, there is a separation in powers of the angle cosine which makes it possible to use 21cm survey data to solve for the power spectrum of matter density fluctuations ==> can solve for cosmological parameters
- When reionization patchiness must also be taken explicitly into account, too, then one can still do this, but requires knowledge of the patchiness and cross-correlations from simulations, to find the best fit for reionization terms, too.

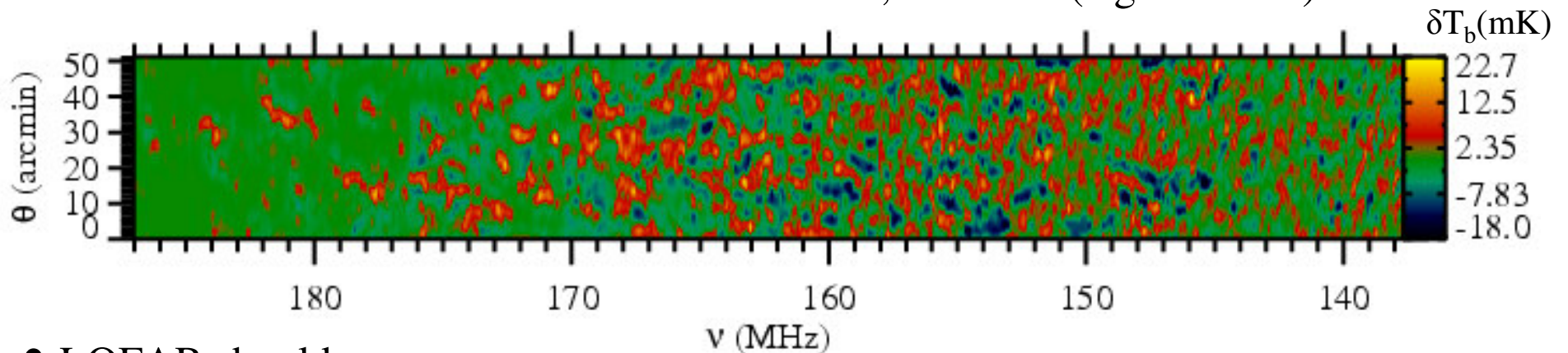
Reionization topology revealed by fluctuations in 21-cm brightness temperature, δT_b , along the line of sight

Iliev, Mellema, Pen, Bond, & Shapiro (2008), MNRAS, 384, 863 (astro-ph/0702099)

- mapping the sky along the LOS: high-resolution cuts in position-redshift space



- beam- and bandwidth-smoothed : 3 arcmin, 0.2 MHz (e.g. LOFAR)



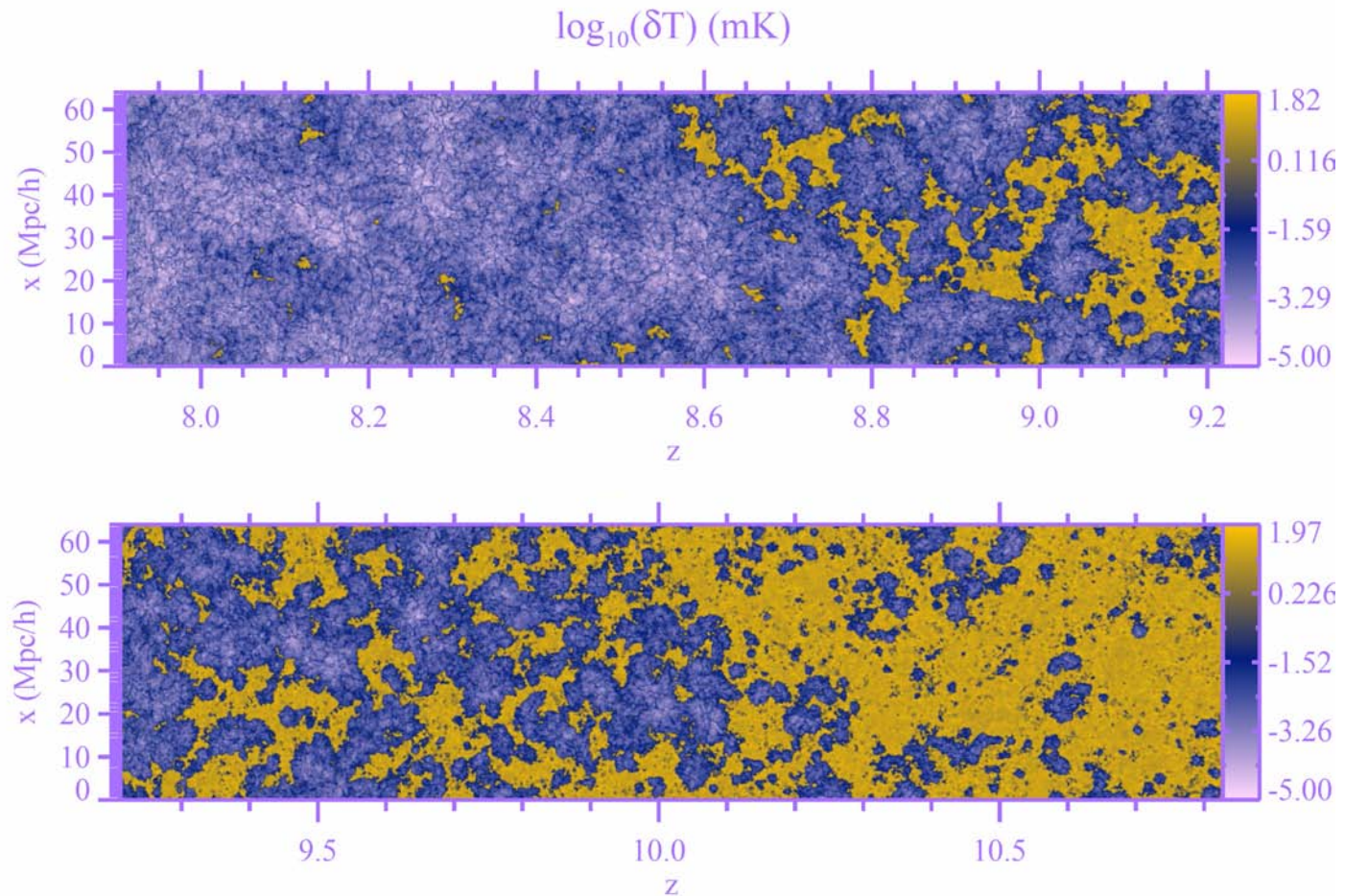
- LOFAR should see
large ionized bubbles!

Case: $f_\gamma = 250$, subgrid clumping factor $C(z)$, WMAP3

Reionization topology revealed by fluctuations in 21-cm brightness temperature, δT_b , along the line of sight

(Shapiro, Iliev, Mellema, Pen, and Merz 2008, AIP 1035, 68; astro-ph/0806.3091)

- mapping the sky along the LOS: high-resolution cuts in position-redshift space

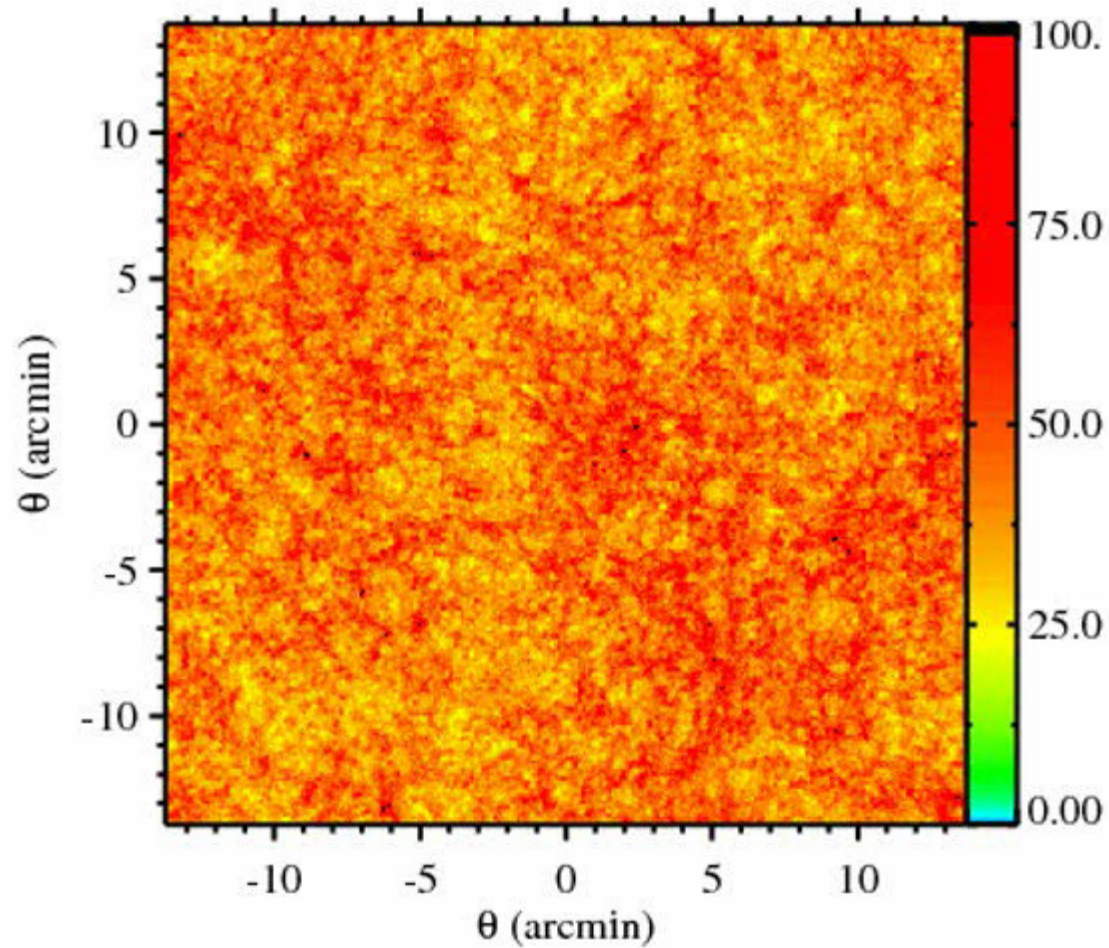


- LOFAR should see
large ionized bubbles!

90 Mpc box, WMAP3+

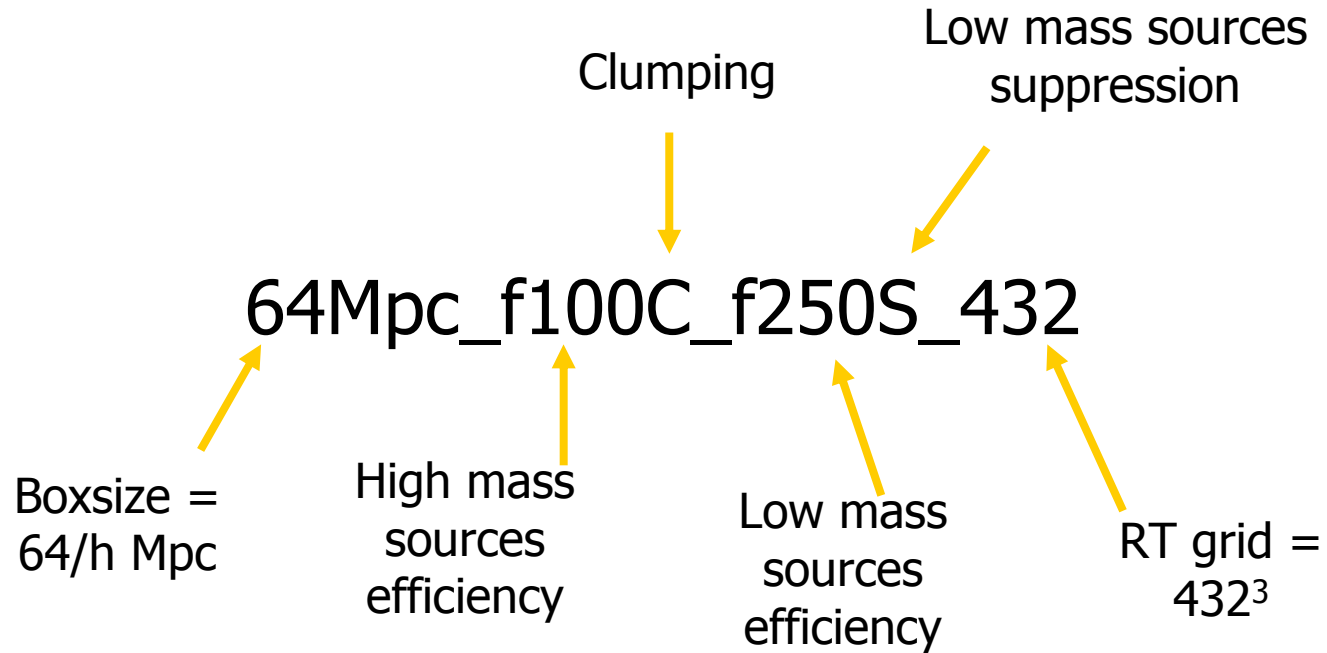
Sky Maps of 21cm Background Brightness Temperature Fluctuations During Epoch of Reionization : Travel through Time

δT (mK) at $z=23.52$ (57.902 MHz)



Notation

- Our simulations are characterized by

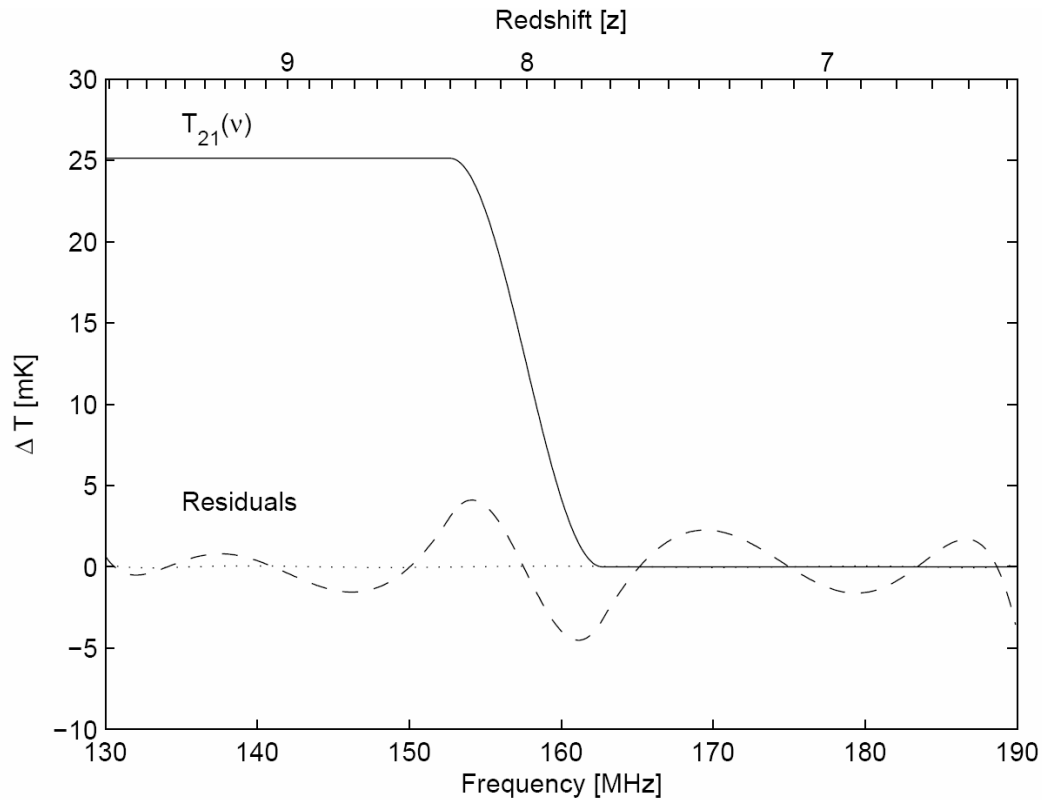


Statistical Measurements of the 21cm Background During the EoR

- The sensitivity of the upcoming EoR experiments will be too low to image 21cm from reionization pixel by pixel: Statistical measurements needed.
 - **First goal:** to reliably detect signatures from reionization (and separate them from foreground and instrumental effects).
 - **Second goal:** to interpret them in terms of astrophysics (source population and properties).
- Luckily, the 21cm line signal is rich in properties:
 1. Global signals: mean signal, fluctuations.
 2. Angular properties: power spectra
 3. Frequency properties: correlation length, Kaiser effect
 4. Non-Gaussianity.

EDGES (Mileura Station, Western Australia)

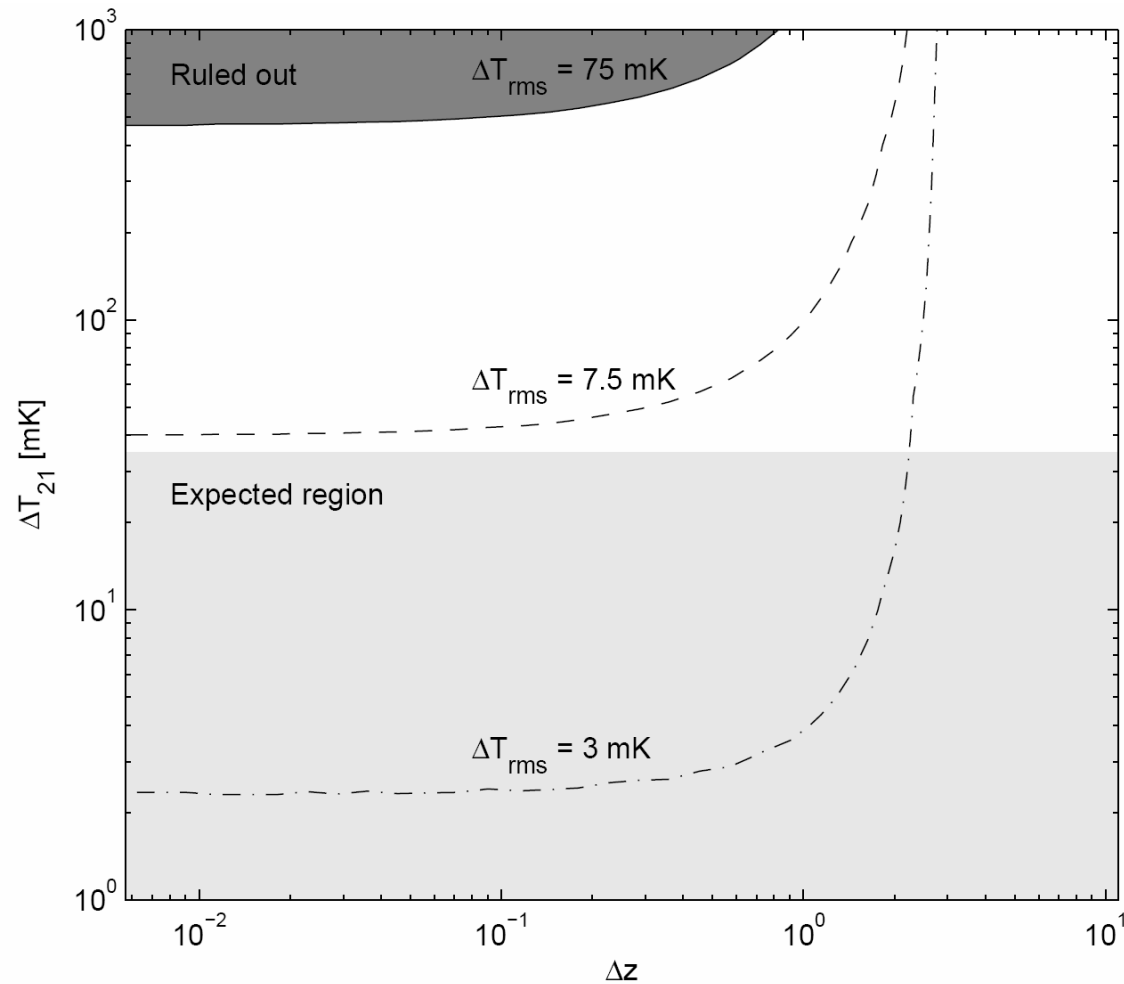
Bowman, Rogers, and Hewitt (2008)



- Search for a “jump” in the global mean brightness temperature when reionization ends

EDGES (Mileura Station, Western Australia)

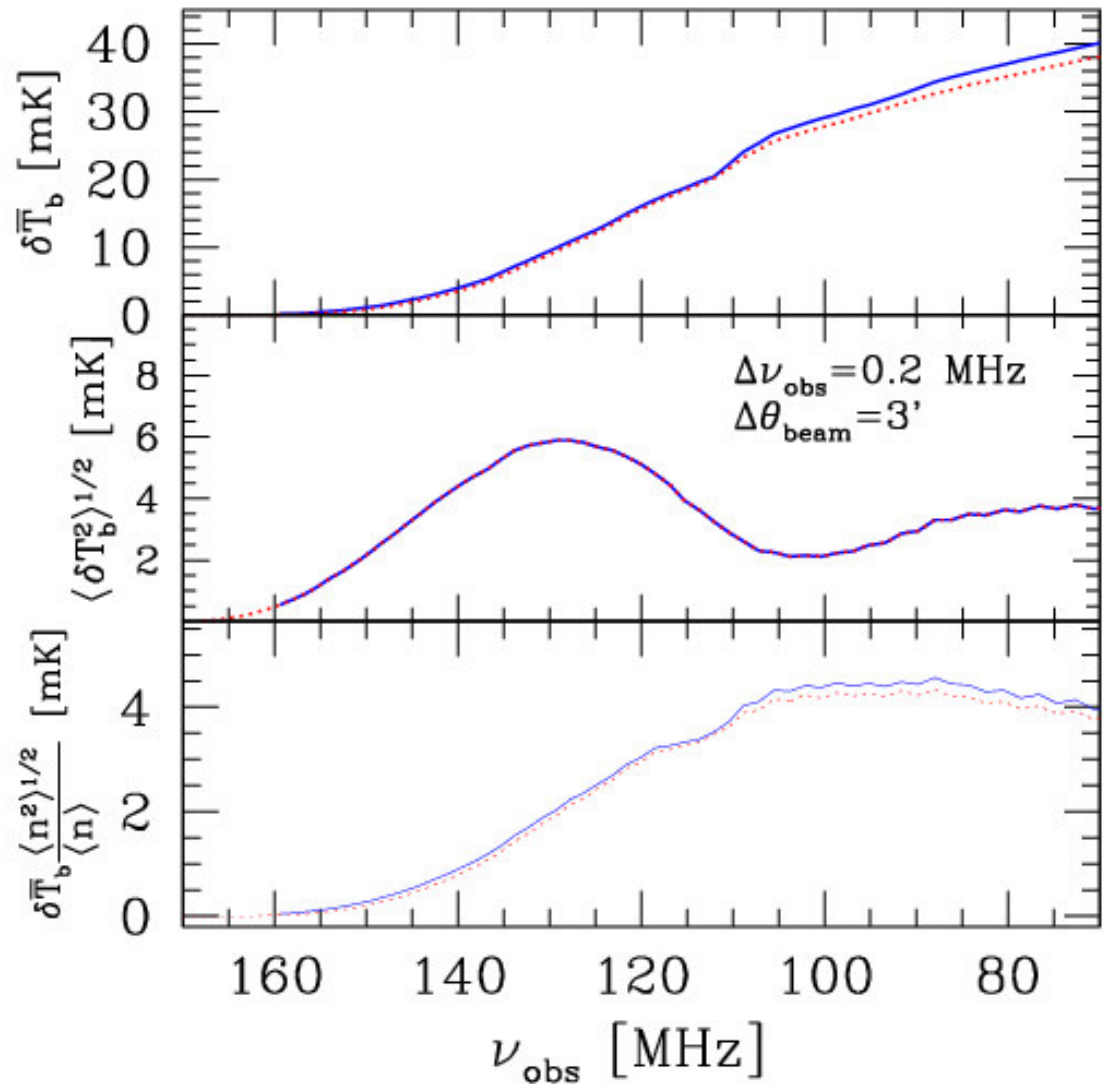
Bowman, Rogers, and Hewitt (2008)



- First results only rule out extreme cases.

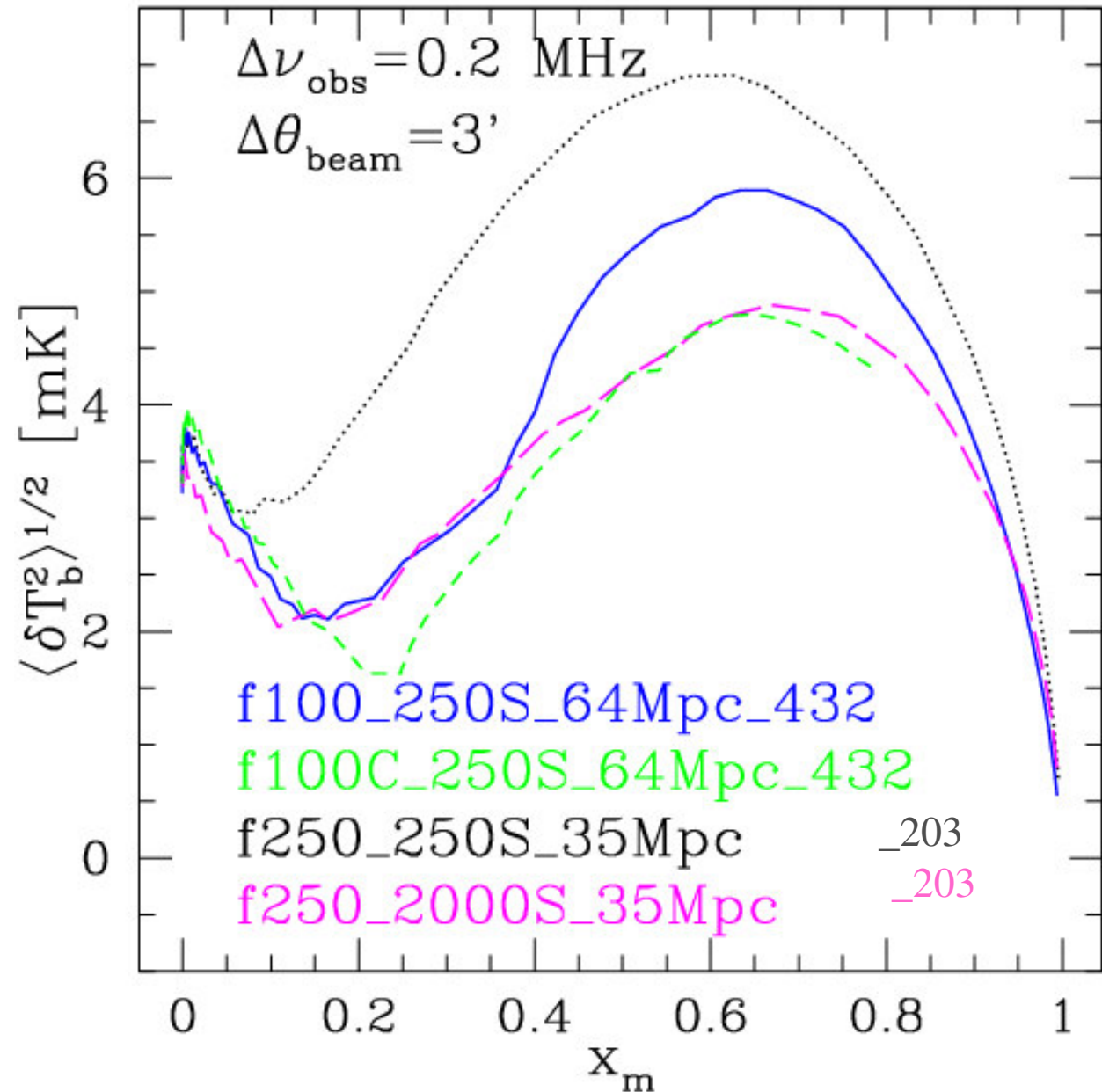
Global Signals

- In principle, a single dish telescope could measure the change of the global signal with frequency: contrary to early expectations, however, simulations do not show a sharp transition.
- The corresponding measurement by an interferometer would be the change of the 21cm (rms) fluctuations.
- Simulations:
64Mpc_f100_f250S_432
and
64Mpc_f100_f250S_216



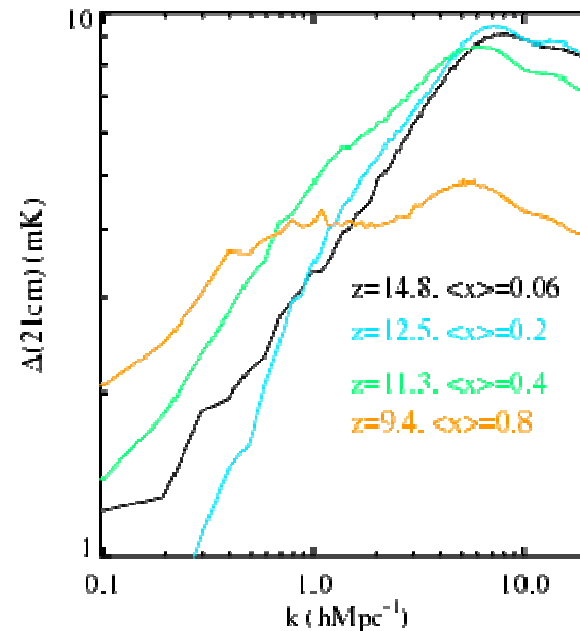
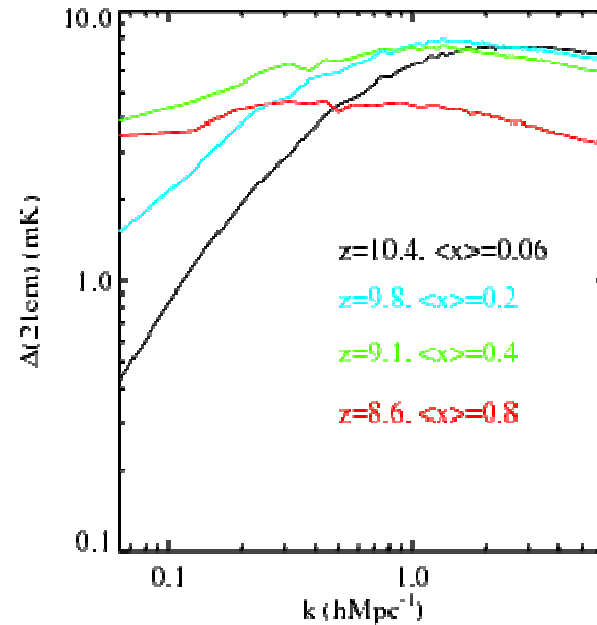
Evolution of 21cm Brightness Temperature Fluctuations

- When plotted against the **mean mass-weighted ionization fraction $x_m(\text{H II})$** , the evolution of fluctuations shows roughly similar behaviour for different (simulation) resolution and source parameters, but the amplitude differs.
- **Peak** around $x_m(\text{H II}) \sim 0.6-0.7$ (shifts to lower values for higher angular resolution).



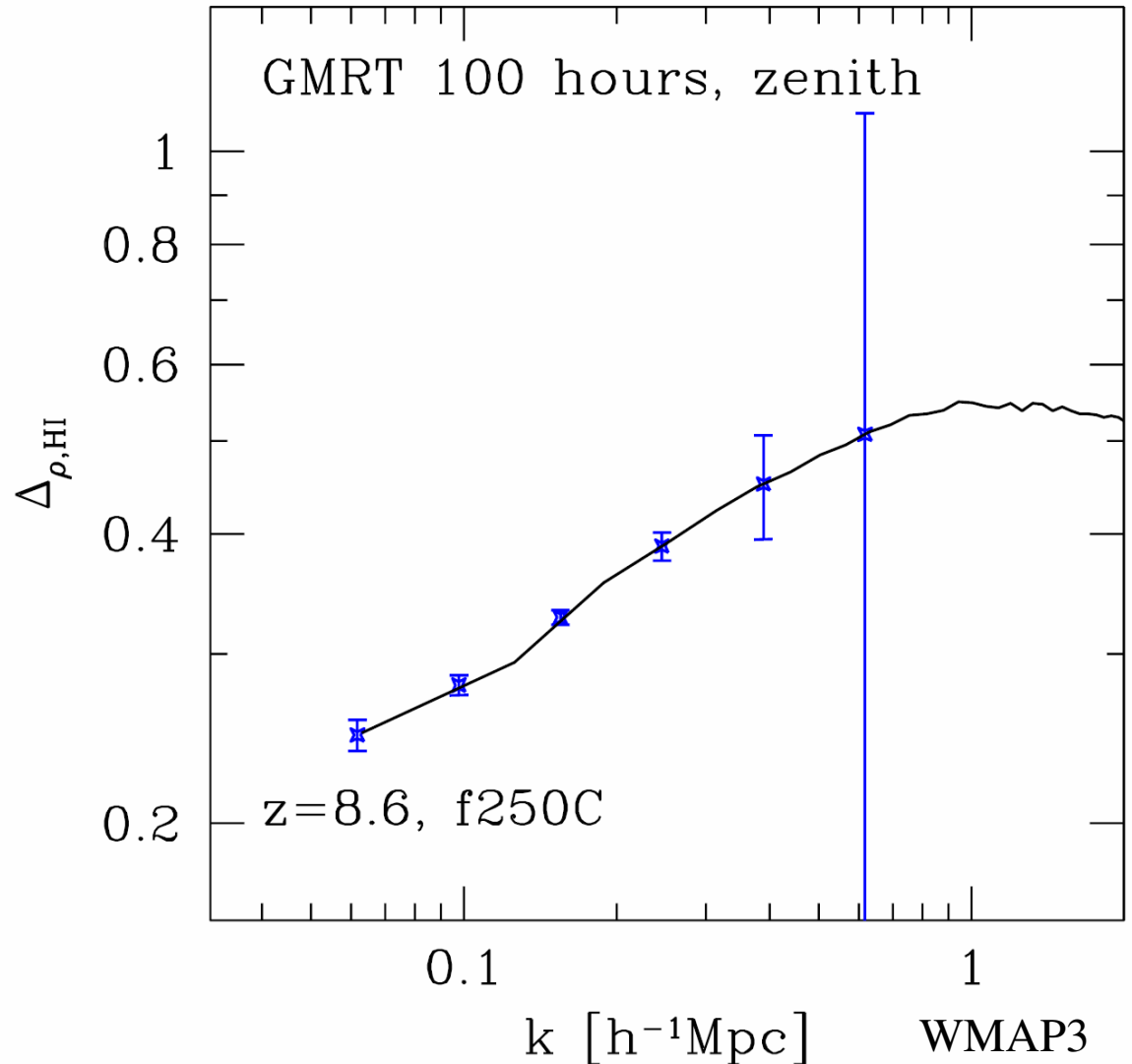
21cm Brightness Temperature Fluctuation Power Spectra

- Information about the **length scales** can be obtained from the power spectra.
- Power shifts to larger scales as reionization progresses, and the power spectrum **flattens**.
- Note that the angular power spectrum is measured directly by an interferometer, the multipole l is equivalent to $\sqrt{u^2+v^2}$ in a visibility map.



Observability of the 21-cm signal : 3D power spectrum of the neutral hydrogen density

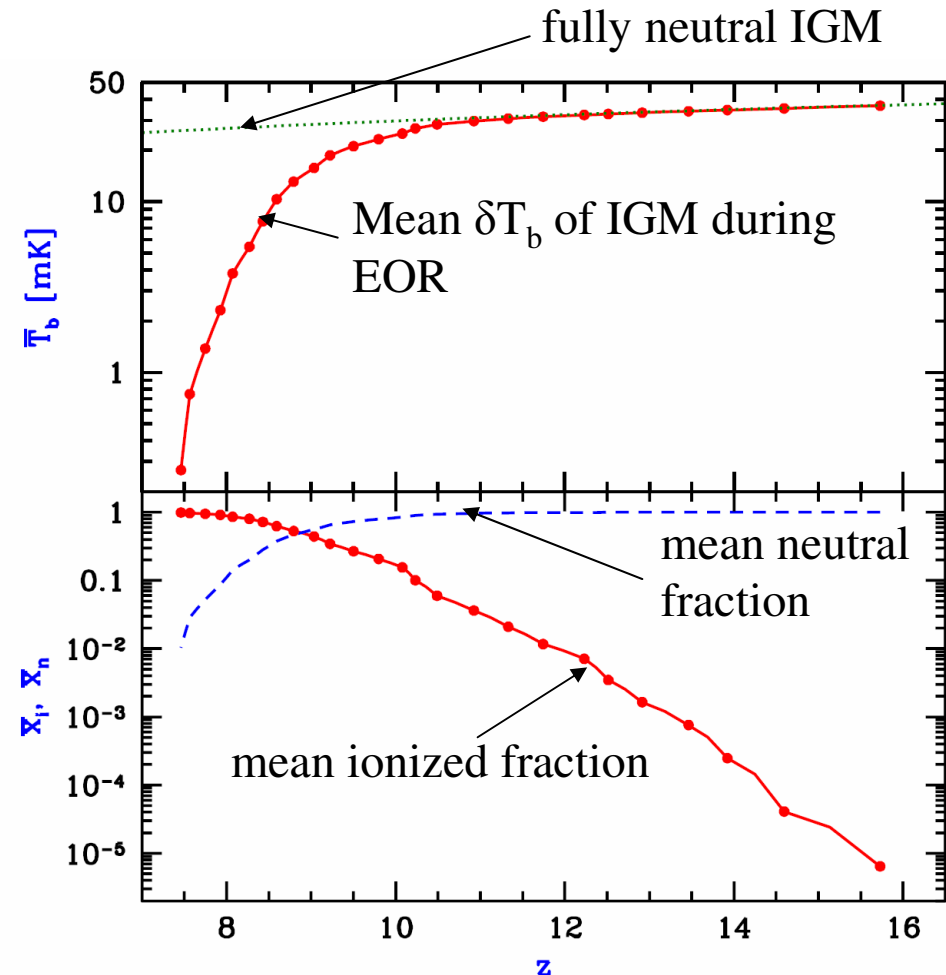
- We show our predicted 21-cm signal at $z = 8.6$ ($x \sim 0.5$) for case with $z_{\text{ov}} = 6.6$, with GMRT sensitivity, for 100 hrs integration, 15 MHz bandwidth and $T_{\text{sys}} = 480$ K, pointing at zenith.



Measuring the History of Cosmic Reionization Using the 21cm PDF

Ichikawa, Barkana, Iliev, Mellema & Shapiro 2009, MNRAS, submitted (astro-ph/0907.2932)

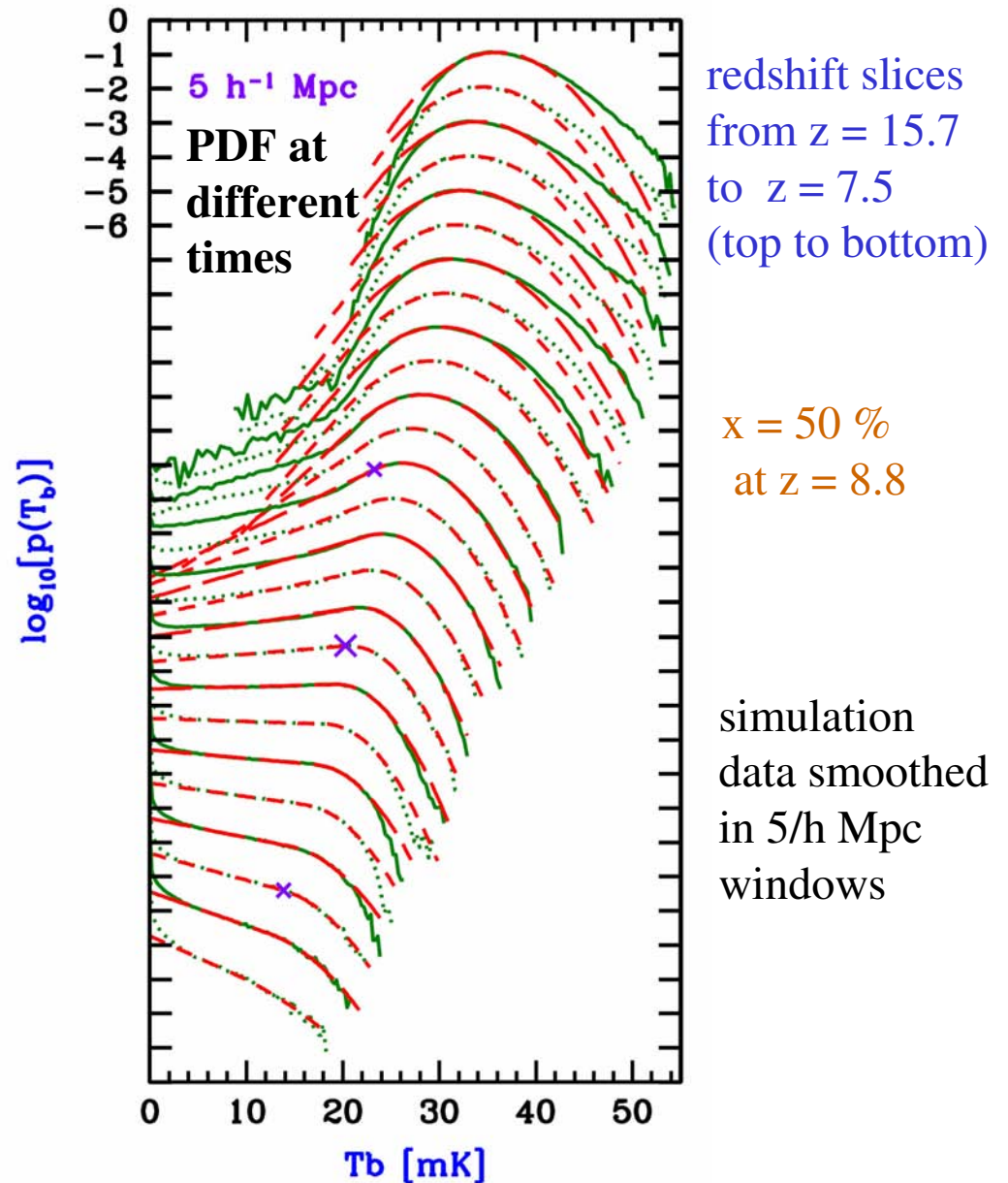
- The **Probability Distribution Function (PDF)** of 21-cm differential brightness temperature is a highly nonGaussian signature of the evolving patchy reionization during the EOR.
- Two peaks in PDF:
 $\delta T_b = 0$ mK (ionized regions)
 ≈ 20 mK (neutral regions)
- We use our reionization simulation of large volume $100/h$ comoving Mpc to predict this 21cm PDF and propose a simple empirical fit
- For the simplest parameterization, upcoming 21cm radio surveys like MWA can recover the reionization history to 1 – 10% accuracy at middle to late stages of EOR
- More realistic fits with more free parameters \rightarrow 2nd generation surveys



Measuring the History of Cosmic Reionization Using the 21cm PDF

Ichikawa, Barkana, Iliev, Mellema & Shapiro 2009. MNRAS. submitted (astro-ph/0907.2932)

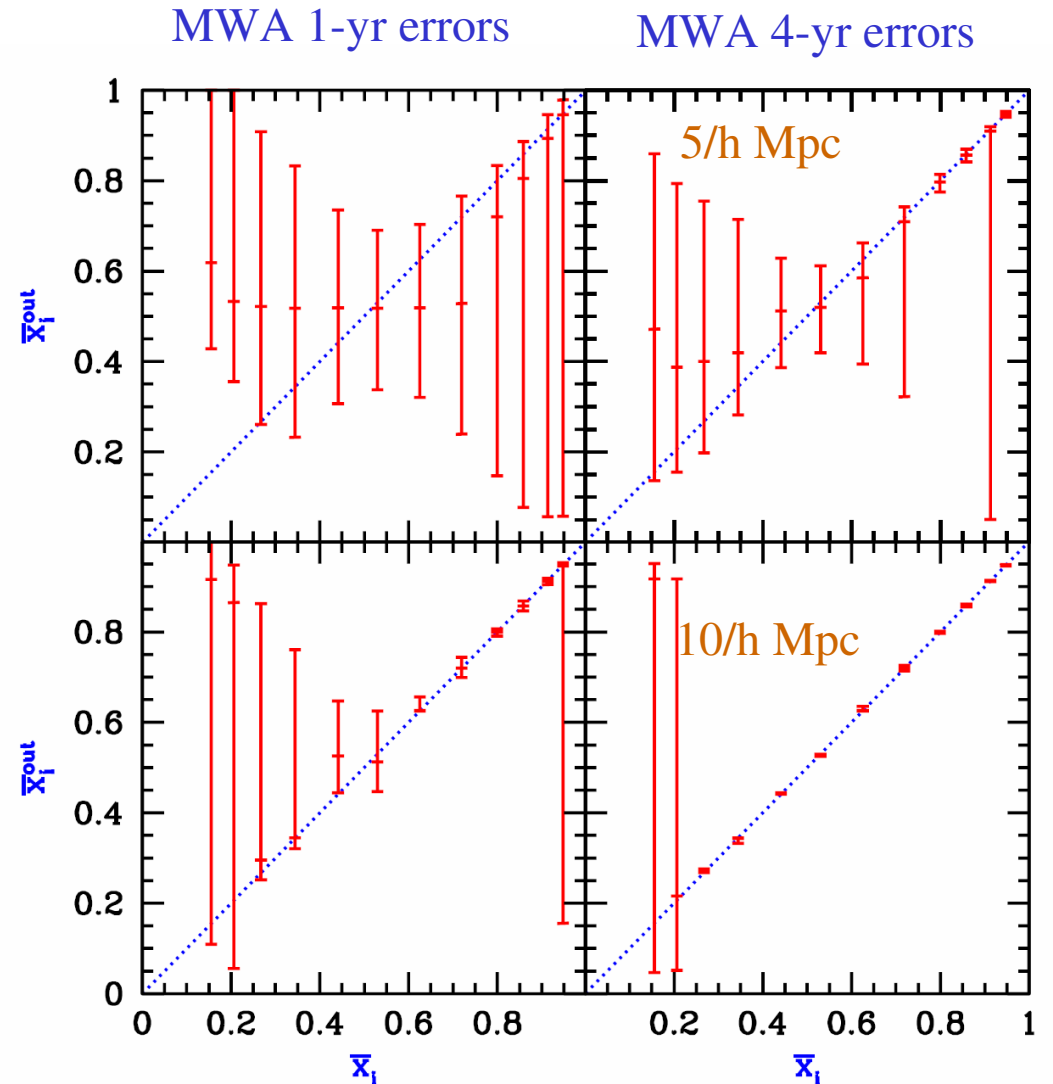
- The **Probability Distribution Function (PDF)** of 21-cm differential brightness temperature is a highly nonGaussian signature of the evolving patchy reionization during the EOR.
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 $\delta T_b = 0$ mK (ionized regions)
 ≈ 20 mK (neutral regions)
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Median values measured by MWA for ionized fraction of universe (and 16 and 84th percentiles)