#### [Toward] Simulating Cosmological Reionization with Enzo

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#### Personal Introduction

- Assistant professor in Mathematics at SMU.
- Computational/applied math background; expertise in nonlinear and linear solvers for multi-scale problems in computational physics.
- This work focuses on collaborations with UCSD's Laboratory for Computational Astrophysics (M. Norman et al.).
- Goal is to develop high-accuracy, scalable solvers for coupled simulations of radiation, hydrodynamics and ionization.
- Efforts focused on developing solver modules for *Enzo*, an open-source code for cosmological hydro., dark-matter, gravity and chemical ionization (advertised by Kim Tran earlier).
- Target applications are mainly in cosmological reionization, though goal is to develop numerical methods that span to other regimes as well.









## Enzo-RT Goals

- I. Extend Enzo to include radiation transport & ionization, to enable:
  - studies of self-regulated star formation,
  - predictions on the epoch of cosmic reionization, and
  - predictions on the observed properties of early galaxies.
- II. Accurately model stiff cosmological RT and chemical ionization processes.
- III. Integrate new solver in Enzo code so that coupling respects hydrodynamics approach (shocks), but tightly couples physics.
- IV. Enable very-large-scale simulations ( $\mathcal{O}(10^6)$  processors).









#### Coupled Matter-Radiation System

We consider the coupled cosmological PDE system,

$$\nabla^2 \phi = \frac{4\pi G}{a} \left( \rho_b + \rho_{dm} - \rho_0 \right),$$
  

$$\partial_t \rho_b + \frac{1}{a} \mathbf{v}_b \cdot \nabla \rho_b = -\frac{1}{a} \rho_b \nabla \cdot \mathbf{v}_b,$$
  

$$\partial_t \mathbf{v}_b + \frac{1}{a} \left( \mathbf{v}_b \cdot \nabla \right) \mathbf{v}_b = -\frac{\dot{a}}{a} \mathbf{v}_b - \frac{1}{a\rho_b} \nabla p - \frac{1}{a} \nabla \phi,$$
  

$$\partial_t e + \frac{1}{a} \mathbf{v}_b \cdot \nabla e = -\frac{2\dot{a}}{a} e - \frac{1}{a\rho_b} \nabla \cdot (p \mathbf{v}_b) - \frac{1}{a} \mathbf{v}_b \cdot \nabla \phi + G - \Lambda,$$
  

$$\partial_t \mathbf{n}_i + \frac{1}{a} \nabla \cdot (\mathbf{n}_i \mathbf{v}_b) = -3 \frac{\dot{a}}{a} \mathbf{n}_i - \mathbf{n}_i \Gamma_i^{ph} + \alpha_{i,j}^{rec} \mathbf{n}_e \mathbf{n}_j,$$
  

$$\partial_t E_\nu + \frac{1}{a} \nabla \cdot (E_\nu \mathbf{v}_b) = \nabla \cdot (D \nabla E_\nu) + \nu \frac{\dot{a}}{a} \partial_\nu E_\nu + 4\pi \eta_\nu - ck_\nu E_\nu.$$

Here,  $G(E, \mathbf{n}_i)$  and  $\Lambda(E, \mathbf{n}_i)$  are the heating and cooling rates. The  $E_{\nu}$  equation approximates the radiative flux as a function of the energy density gradient,

$$\mathbf{F}_{\nu} = -D\,\nabla E_{\nu},$$

where  $D(E_{\nu}, \nabla E_{\nu}) \in \mathbb{R}^{3 \times 3}$  is the *flux limiter*.

[Bryan et al., Comp. Phys. Comm., 1995; R. et al., J. Comput. Phys., 2009]









## Enzo Operator-Split Numerics

We solve this coupled system in an operator-split framework, solving one component of the system at a time within a time step:

- (i) Project dark matter particles onto mesh to generate  $\rho_{dm}$ .
- (ii) Solve for the gravitational potential  $\phi$  via FFT or MG methods.
- (iii) Advect dark matter particles via Particle-Mesh method.
- (iv) Explicitly evolve  $(\rho_b, \mathbf{v}_b, e)$  and advect  $(E_{\nu}, \mathbf{n}_i)$  with a high-order PPM method.
- (v) Implicitly evolve a stiff reaction-di ffusion PDE system that updates  $(E_{\nu}, e, \mathbf{n}_i)$ .

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#### Gray FLD Radiation Transfer Equation

Prior to investigating the multi-frequency case, we begin using a single, frequency-integrated ("grey") radiation energy,

$$E(x,t) = \int_{\nu_0}^{\infty} E_{\nu}(x,t,\nu) \, d\nu = \tilde{E}(x,t) \int_{\nu_0}^{\infty} \chi_E(\nu) \, d\nu,$$

where  $\chi_E : I\!\!R \to I\!\!R$  is an assumed radiation energy density spectrum.

With this approximation, the integrated radiation energy equation becomes

$$\partial_t E + \frac{1}{a} \nabla \cdot (E \mathbf{v}_b) = \nabla \cdot (D \nabla E) - \frac{\dot{a}}{a} E + 4\pi \eta - ckE,$$

where the new coupling terms are integrated over frequency:

$$\eta = \int_{\nu_0}^{\infty} \eta_{\nu} \, d\nu$$
$$k = \left( \int_{\nu_0}^{\infty} k_{\nu} E_{\nu} \, d\nu \right) / \left( \int_{\nu_0}^{\infty} E_{\nu} \, d\nu \right).$$









#### Rad-Hydro Verification Results



Smearing of a radiation front in a vacuum due to FLD approximation (convergence wrt  $\Delta x$ ).

Thermal equilibration: energy conserv. error,

$$\left(\int \left|\Delta E_{\mathsf{tot}}\right| \mathrm{d}x\right) / \left(\int E_{\mathsf{tot}}(0) \mathrm{d}x\right),$$

L→H and H→L equilibration, under various nonlinear and linear tolerances ( $\varepsilon$ ,  $\delta$ ). [Turner & Stone, *ApJ*, 2001]



Radiating shock front convergence wrt  $\Delta x$ . [Lowrie & Edwards, *ShWav*, 2008].

ε	δ	$L \rightarrow H$	$H \rightarrow L$	
1e-7	1e-9	7.0e-12	1.4e-12	
1e-10	1e-9	3.1e-13	3.5e-14	
1e-4	1e-9	3.6e-11	1.8e-11	
1e-7	1e-12	7.0e-12	1.4e-12	
1e-7	1e-3	7.0e-12	1.4e-12	







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## Isothermal Radiative Ionization Test

Tests of isothermal ionization of a static, initially-neutral hydrogen region:

- L = 6.6 kpc.
- Monochromatic emission:  $\dot{N}_{\gamma} = 5 \cdot 10^{48}$  photons s<sup>-1</sup>, at  $h\nu = 13.6$  eV.
- Fixed gas temperature,  $T = 10^4$  K.
- Case-B recombination rate,  $\alpha_B = 2.59 \cdot 10^{-13} \text{ cm}^3 \text{ s}^{-1}$ .
- Constant number density  $n = 10^{-3}$  cm<sup>-3</sup>.

Analytical solution given by

$$r_{I} = r_{S} \left[ 1 - e^{-t\alpha_{B}n_{H}} \right], \text{ where}$$

$$r_{S} = \left[ \frac{3\dot{N}_{\gamma}}{4\pi\alpha_{B}n_{H}^{2}} \right]^{1/3}.$$

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[Iliev et al., MNRAS, 2006]









## Isothermal Ionization Results



NW: I-front position history.

NE: Coarse radiation field  $(N_x = 16)$ .



SE: Fine radiation field  $(N_x = 128)$ .













# Isothermal Cosmological Ionization Test

Repeat of previous test, but in a cosmologically expanding universe.

Four tests:

<i>q</i> <sub>0</sub>	$z_i$	$L_i$ [kpc]	$ ho_{b,i}~[{ m g~cm^{-3}}]$	H <sub>0</sub>	$\Omega_m$	$\Omega_{\Lambda}$	$\Omega_b$
0.5	4	80	1.18e-28	0.5	1.0	0	0.2
0.05	4	60	2.35e-28	1.0	0.1	0	0.1
0.5	10	36	1.18e-28	0.5	1.0	0	0.2
0.05	10	27	2.35e-28	1.0	0.1	0	0.1

Analytical solution given by

$$r_I(t) = r_{S,i} \left( \lambda e^{-\tau(t)} \int_1^{a(t)} e^{\tau(b)} [1 - 2q_0 + 2q_0(1 + z_i)/b]^{-1/2} db \right)^{1/3},$$

where

$$\tau(a) = \lambda \left[ F(a) - F(1) \right] \left[ 6q_0^2 (1+z_i)^2 \right]^{-1},$$
  

$$F(a) = \left[ 2 - 4q_0 - 2q_0 \frac{1+z_i}{a} \right] \left[ 1 - 2q_0 + 2q_0 \frac{1+z_i}{a} \right]^{1/2},$$
  

$$\lambda = \alpha_B n_{H,i} / H_0 / (1+z_i).$$

[Shapiro & Giroux, ApJ, 1987]









## Cosmological Ionization Results













#### Hydrodynamic Radiative Ionization Test

Re-do first test in a dynamic medium (L = 15 kpc); combines radiation, hydrodynamics and hydrogen ionization:

- Emission source has a  $T = 10^5$  K blackbody spectrum.
- Initial temperature set to  $T = 10^2$  K.
- Front transitions from R-type to D-type as it reaches Strömgren radius.
- Eventually stalls at  $r_f = r_S \left(\frac{2T_i}{T_e}\right)^{2/3}$ , where  $T_i$  and  $T_e$  are the temperatures behind and ahead of the i-front.

R-type and D-type propgation have different analytical solutions,

$$r_I^R = r_S \left[ 1 - e^{-t\alpha_B(T_i)n_H} \right]^{1/3}, \qquad r_I^D = r_S \left[ 1 + \frac{7c_s t}{4r_S} \right]^{4/7},$$

with true solution somewhere between the two.

[Whalen & Norman, ApJS, 2006; Iliev et al., MNRAS, 2009]









## Hydrodynamic Radiative Ionization Results











#### The Need for Multi-Frequency

However, in comparison with other RT codes, we get peculiar results:

Right: n profiles at 250 Myr (cyan is Enzo) [1].

Below: Enzo n profiles results at 200, 500 Myr.

Lower right: Zeus-MP results (solid is MF, dashed is monochromatic).















# Summary of Current Results

We've achieved a second-order accurate (space and time) coupled solver for cosmological radiation, hydrodynamics, self gravity and chemical ionization.

- Captures shock fronts due to trusted PPM hydrodynamics approach.
- Accurately solves couplings between radiation, ionization and gas energy, due to implicit formulation and coupled solvers.
- Highly scalable, with current tests using > 4000 processors, thanks to reliance on optimal multigrid methods.

However, the grey radiation approximation has its shortcomings:

- Single radiation field allows full absorption by hydrogen, even though higher-frequency radiation should pass through.
- Leads to increased absorption in optically-thick regions, with no pre-heating ahead of I-fronts.









# Continuing Work

- Enhance stability of chemistry computations at large  $\Delta t$ 
  - Fastest dynamics occur in ionization, ostensibly allowing complete ionization of a volume in a single step  $[\mathcal{O}(1) \rightarrow \mathcal{O}(10^{-5})]$ .
  - Time inaccuracy can over-shoot to give negative densities.
- Extend radiation approximation to multi-frequency case,  $E \rightarrow E_{\nu}$ 
  - New physical couplings and interpolation of  $\nu$ -space, based offof [Ricotti, Gnedin & Shull, ApJ, 2002]
  - Large-scale solvers for coupled reaction-diffusion systems
- Extend implicit solver software to AMR grids
  - Regular-grid geometric multigrid  $\rightarrow$  AMG or FAC methods







