

# Cosmic IR Background & COBE, Spitzer and WISE

by

Ned Wright (UCLA)

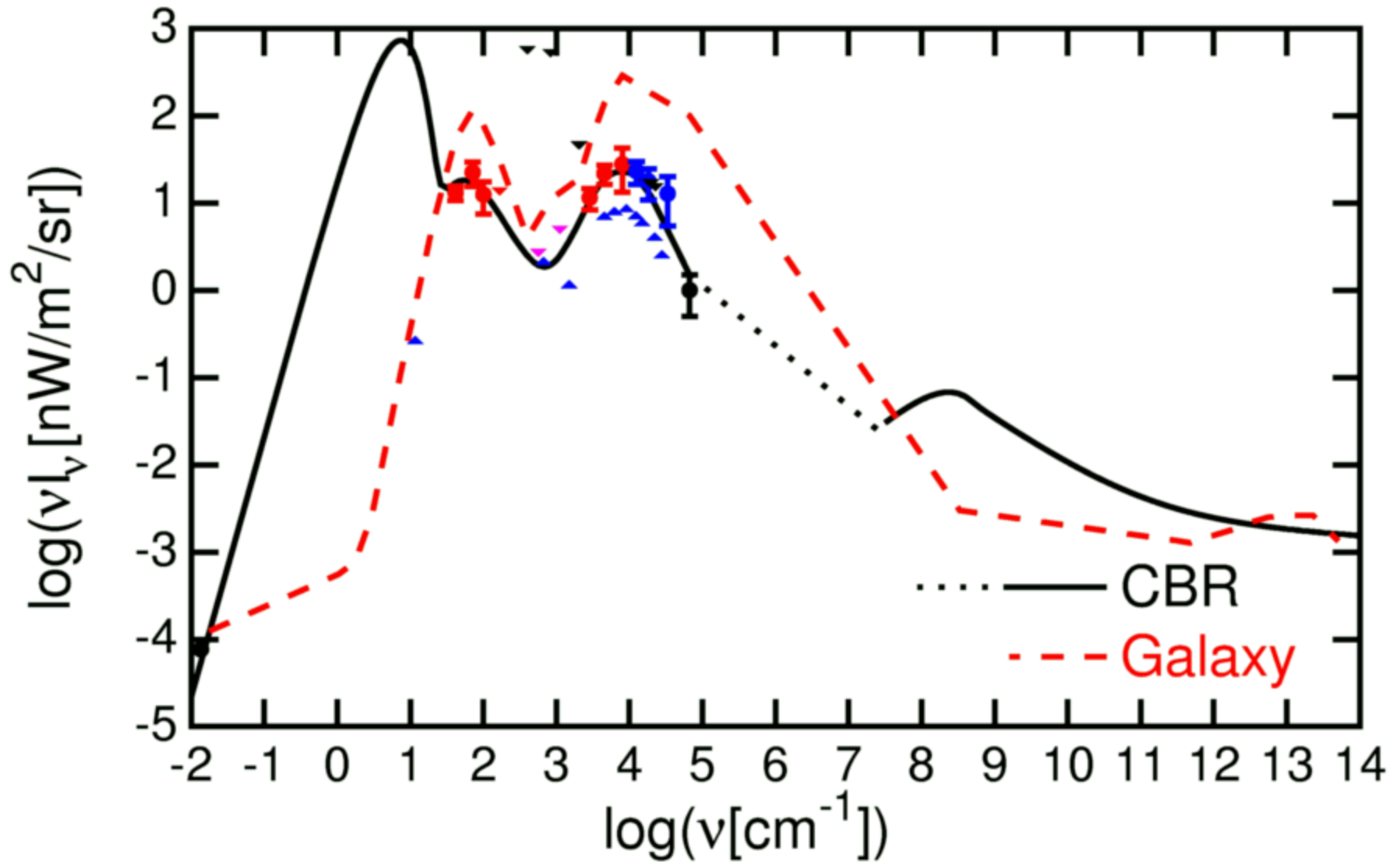
<http://www.astro.ucla.edu/~wright/intro.html>

See:

- <http://www.astro.ucla.edu/~wright/cosmolog.htm>
- <http://www.astro.ucla.edu/~wright/DIRBE>
- <http://www.astro.ucla.edu/~wright/CIBR>
- <http://spitzer.caltech.edu>
- <http://wise.astro.ucla.edu>

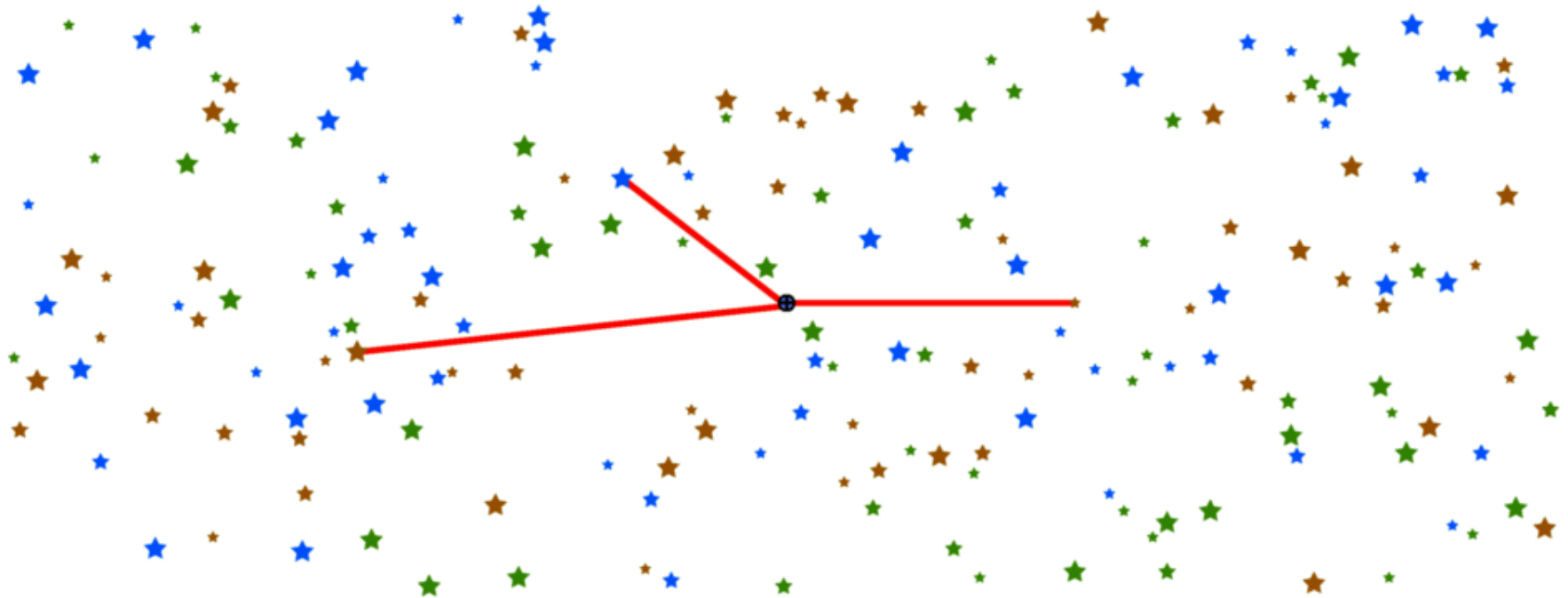
- $I_\nu$  - specific intensity, energy per (area  $\times$  time  $\times$  frequency  $\times$  solid angle).  
 $10^{-20} \text{ W/m}^2/\text{Hz}/\text{sr} = 1 \text{ MegaJansky}/\text{sr} [\text{MJy}/\text{sr}]$ .
- $I = \int I_\nu d\nu$  - bolometric intensity, energy per (area  $\times$  time  $\times$  solid angle).  
 $1 \text{ nW}/\text{m}^2/\text{sr}$ .
- $\nu I_\nu$  - intensity per octave, energy per (area  $\times$  time  $\times$  logarithmic bandwidth  $\times$  solid angle).  
 $1 \text{ nW}/\text{m}^2/\text{sr}$ .
- $J_\nu = (4\pi)^{-1} \int I_\nu d\Omega$ : mean intensity – the angle averaged intensity.
- $F_\nu = \int I_\nu \cos \theta d\Omega$ : monochromatic flux, energy per (area  $\times$  time  $\times$  frequency).  
 $10^{-20} \text{ W}/\text{m}^2/\text{Hz} = 1 \text{ MegaJansky} [\text{MJy}]$ .

# Wide window on the CBR



# Backgrounds

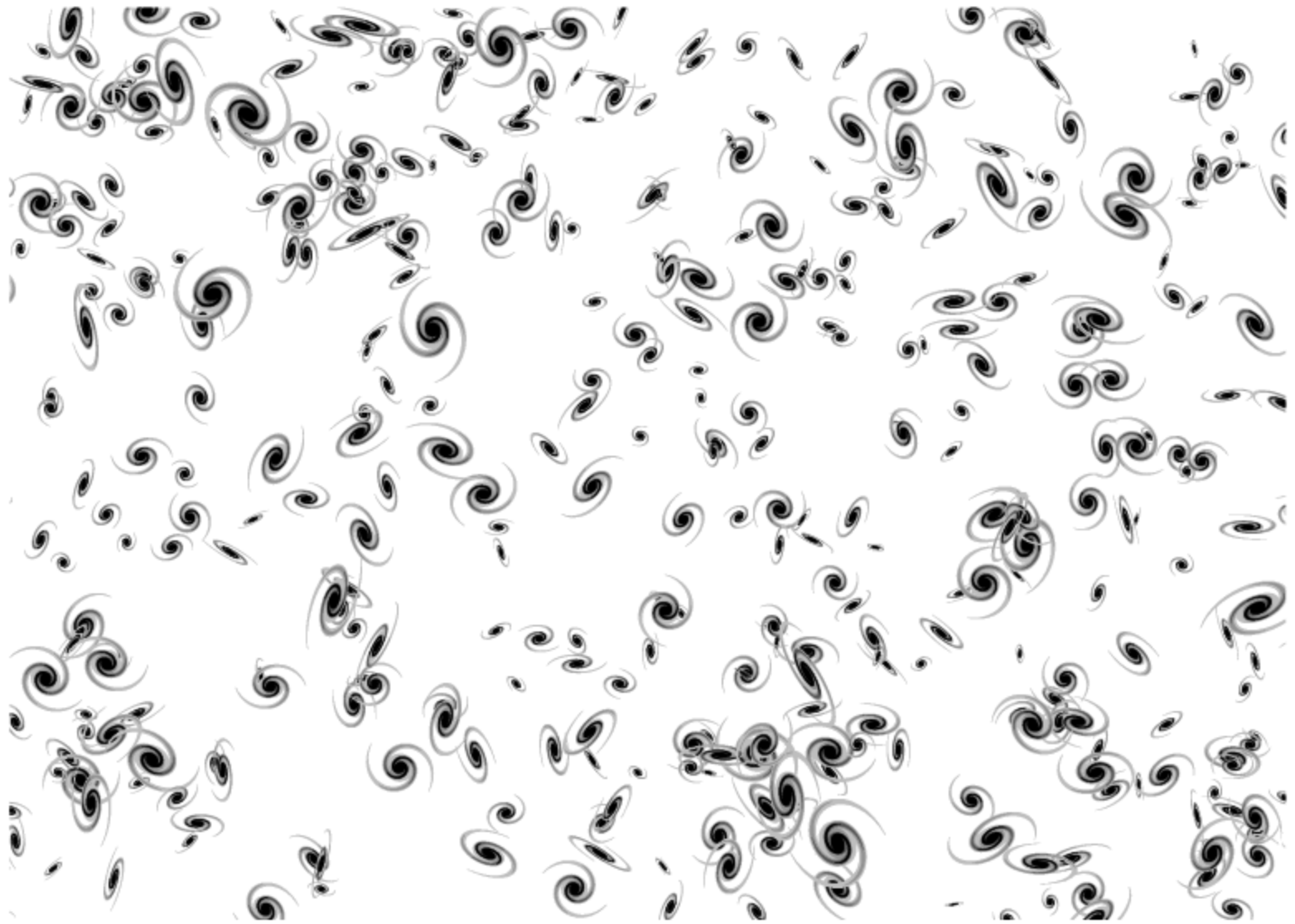
- Microwave – the CMB is 10,000 times brighter than the galactic foreground & the spectrum is very close to a blackbody
- Far Infrared – the FIRB is 10 times fainter than the galaxy with a spectrum similar to the galaxy
- Near IR and Optical – also 10 times fainter than galaxy
- X-ray – the XRB is 10 times brighter than the galaxy

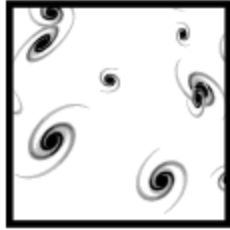


In a homogeneous unchanging Universe every line of sight will end on a star. So why is the night sky not as bright as the surface of a star? The Cosmic Infrared Background is what remains after this Olbers' paradox is resolved.

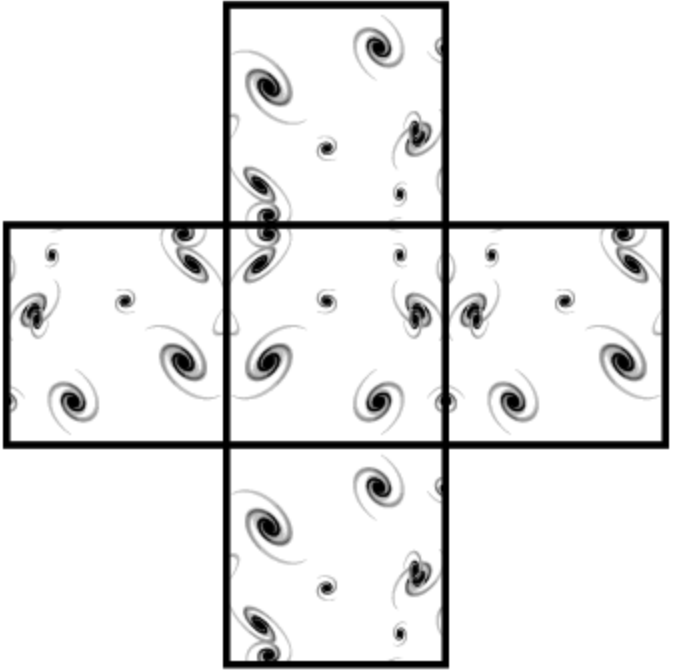
# Sources of the CIRB

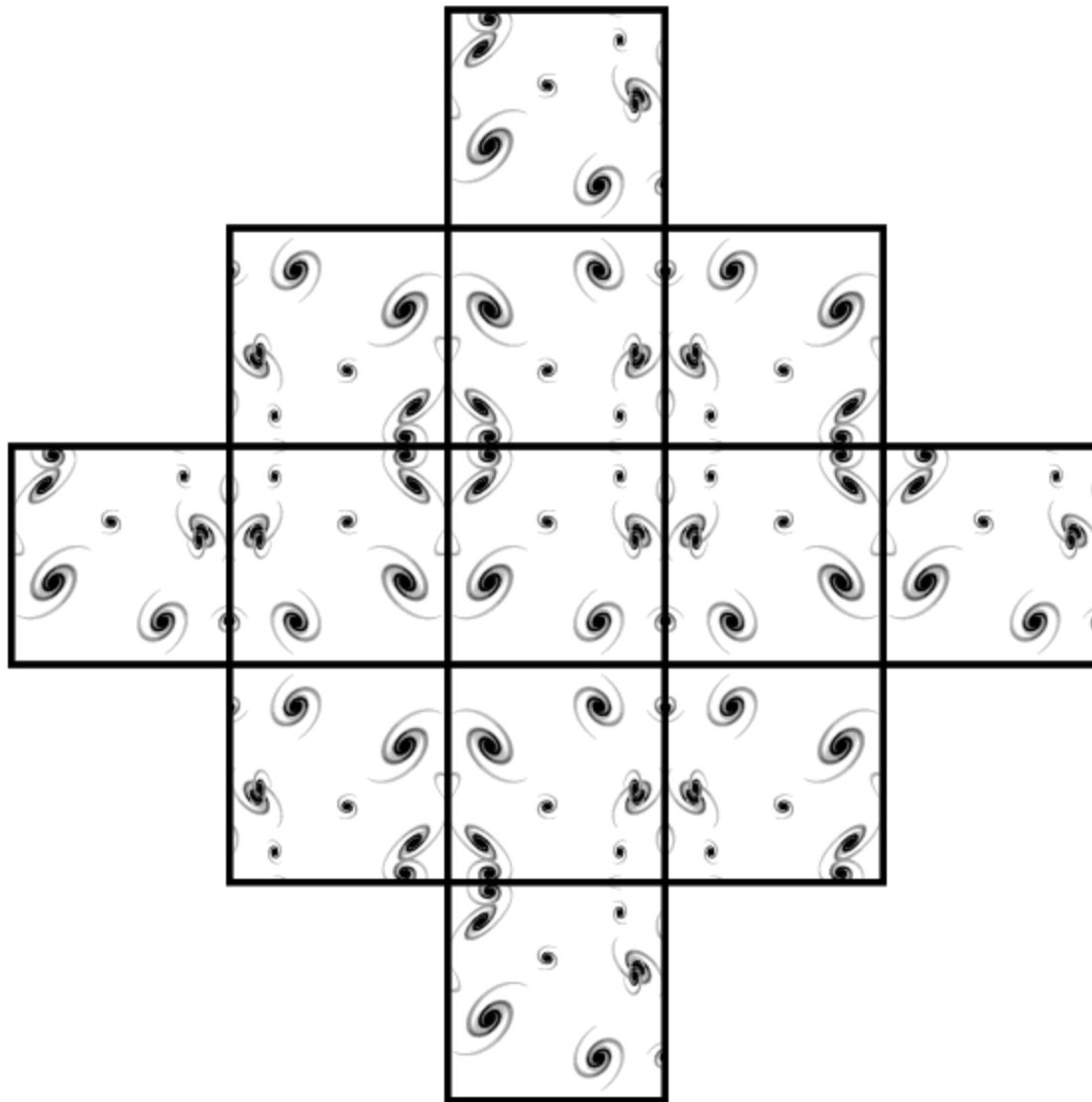


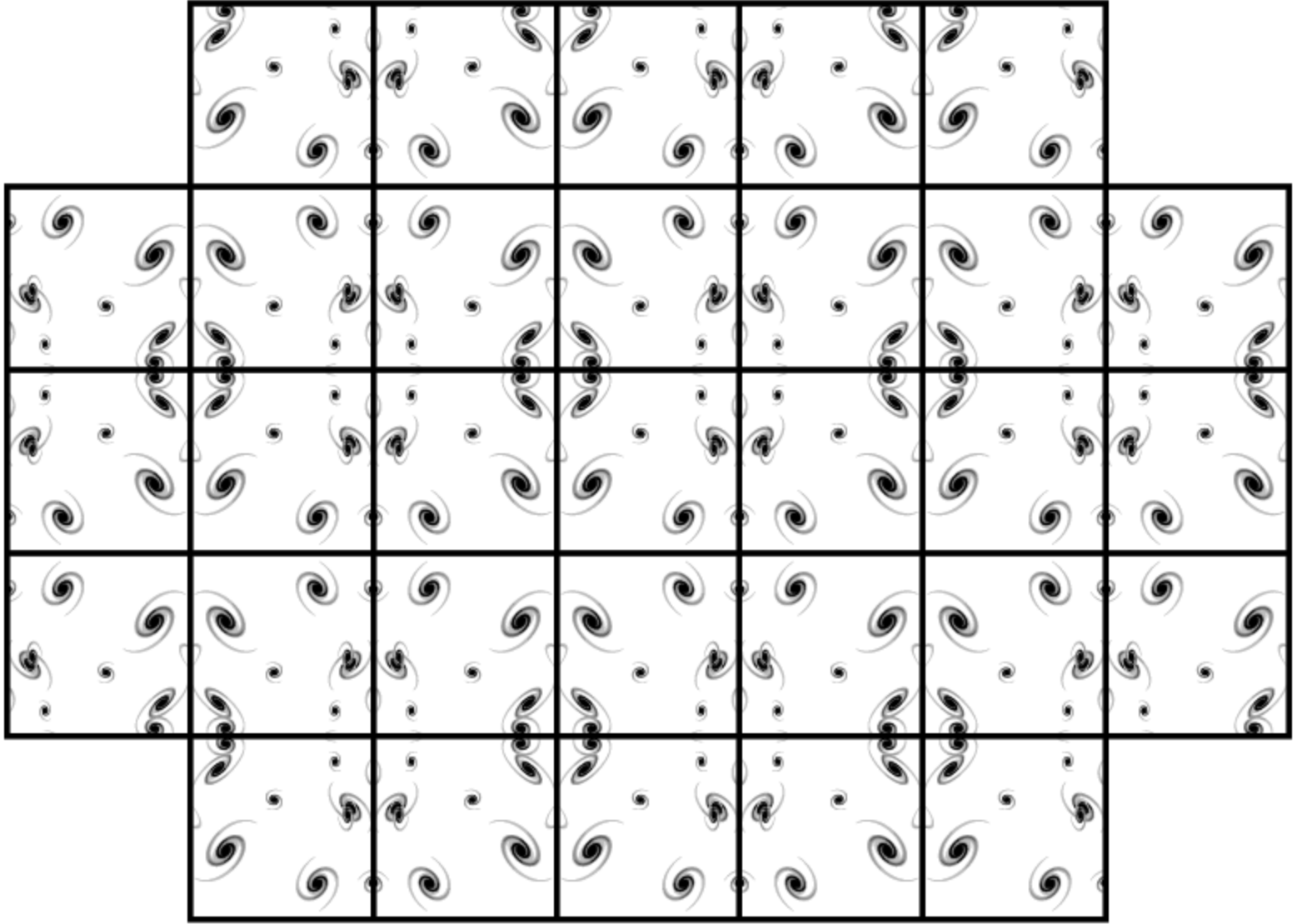


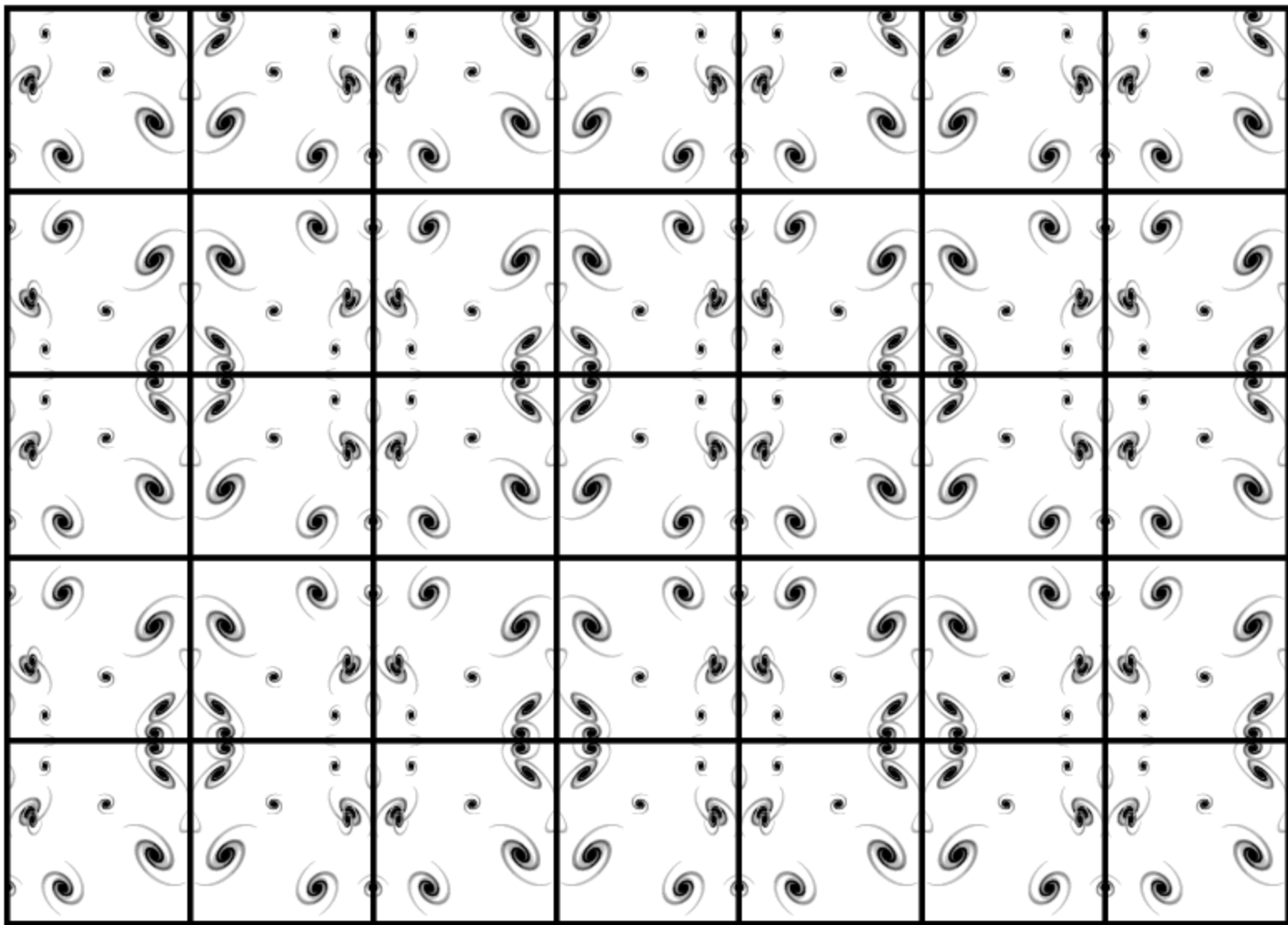






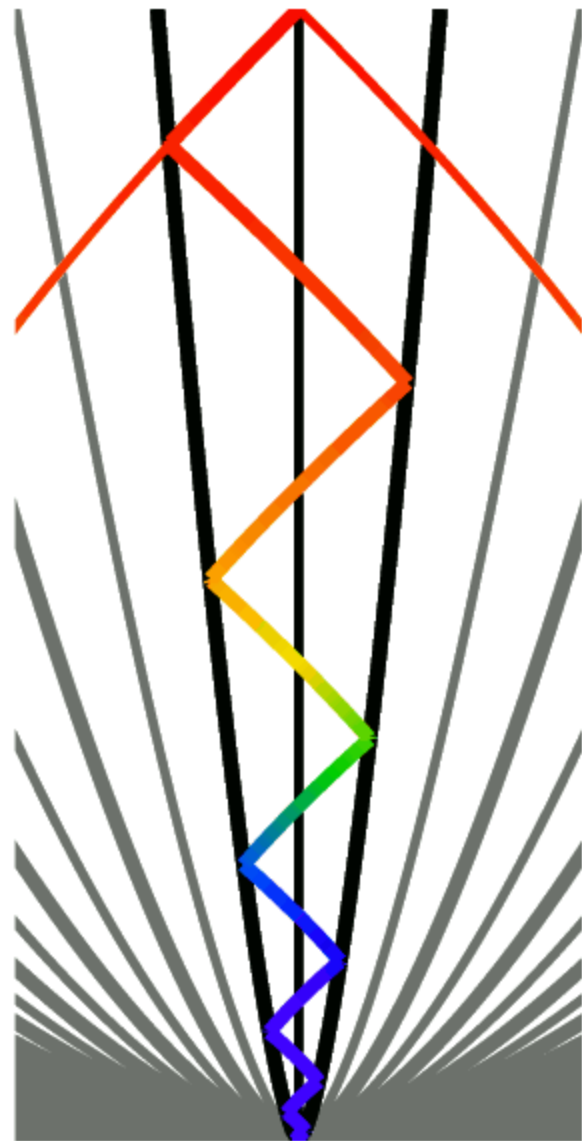






## Expanding Mirrored Box

- The Universe is *homogeneous* & *isotropic*.
- Reflections in a mirrored box look like the Universe.
- Expanding mirrored box gives a redshift with  $\lambda \propto a(t)$ .



Consider a mirrored box, where the mirrors move with the galaxies in the Hubble flow, and define the *comoving luminosity density*:

$$\ell(t) = \frac{\sum_{\text{in box}} L_{gal}(t)}{V_{box}(t_0)}$$

For a static, unchanging Universe the energy density in the box now is

$$u(t_0) = \frac{4\pi J}{c} = \int_{-\infty}^{t_0} \ell(t_0) dt \rightarrow \infty \quad [\text{Olber's Paradox}]$$

In an expanding Universe with scale factor  $a(t)$ , and  $1 + z = a(t_0)/a(t)$ , then:

$$u(t_0) = \frac{4\pi J}{c} = \int_0^{t_0} \frac{\ell(t)}{1+z} dt$$

The total energy produced,  $\int \ell(t)dt$ , is more than the CIRB energy density because it does not have the  $(1+z)$  factor in the denominator. For the baryon density given by Big Bang Nucleosynthesis,  $\Omega_B h^2 = 0.02$ , if 1% of the baryons are converted from hydrogen to helium releasing 0.7% of their mass into energy, then

$$\int \ell(t)dt = 0.02 \times 0.01 \times 0.007 \times 10539.4 \frac{\text{eV}}{\text{cm}^3} = \frac{1}{68} \frac{\text{eV}}{\text{cm}^3}$$

$c/[4\pi]$  times this energy density is 56 nW/m<sup>2</sup>/sr.

The relationship between time and redshift is

$$t = \frac{2}{3H_o(1+z)^{1.5}} \quad \text{so} \quad \frac{dt}{dz} = \frac{-1}{H_o(1+z)^{2.5}}$$

for the Einstein-de Sitter Universe with  $\Omega = 1$ , and the Hubble constant is  $H_o = (2/3)/t_o$ .

Thus if  $\ell(t)$  is CONSTANT, the energy density of the CIRB would be

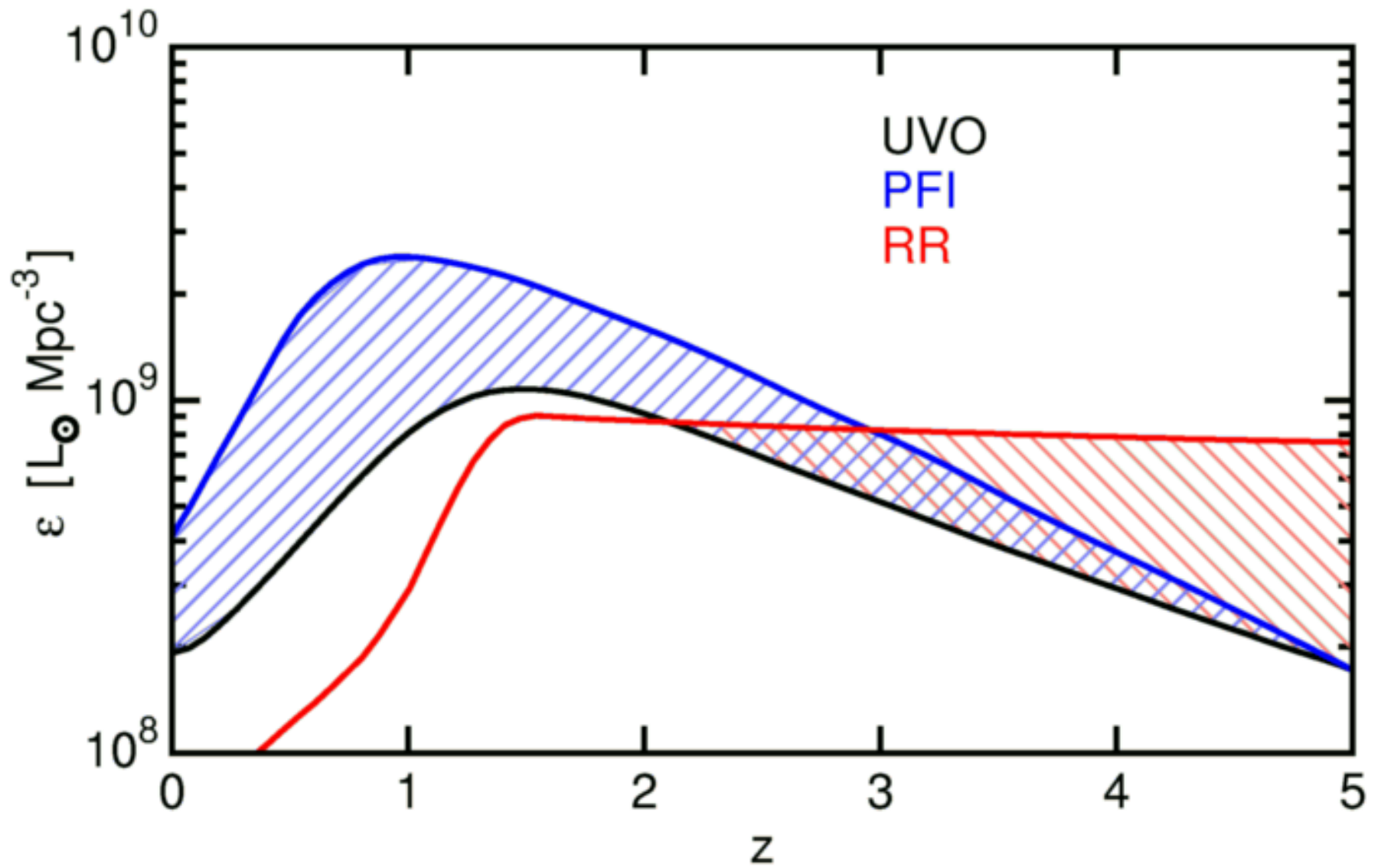
$$u = \frac{1}{H_o} \int \frac{\ell(t_o)}{(1+z)^{3.5}} dz = \frac{3}{5} \int \ell(t) dt = 0.6\ell(t_o)t_o$$

If  $\ell(t)$  is proportional to  $(1+z)$ , then the energy density of the CIRB would be

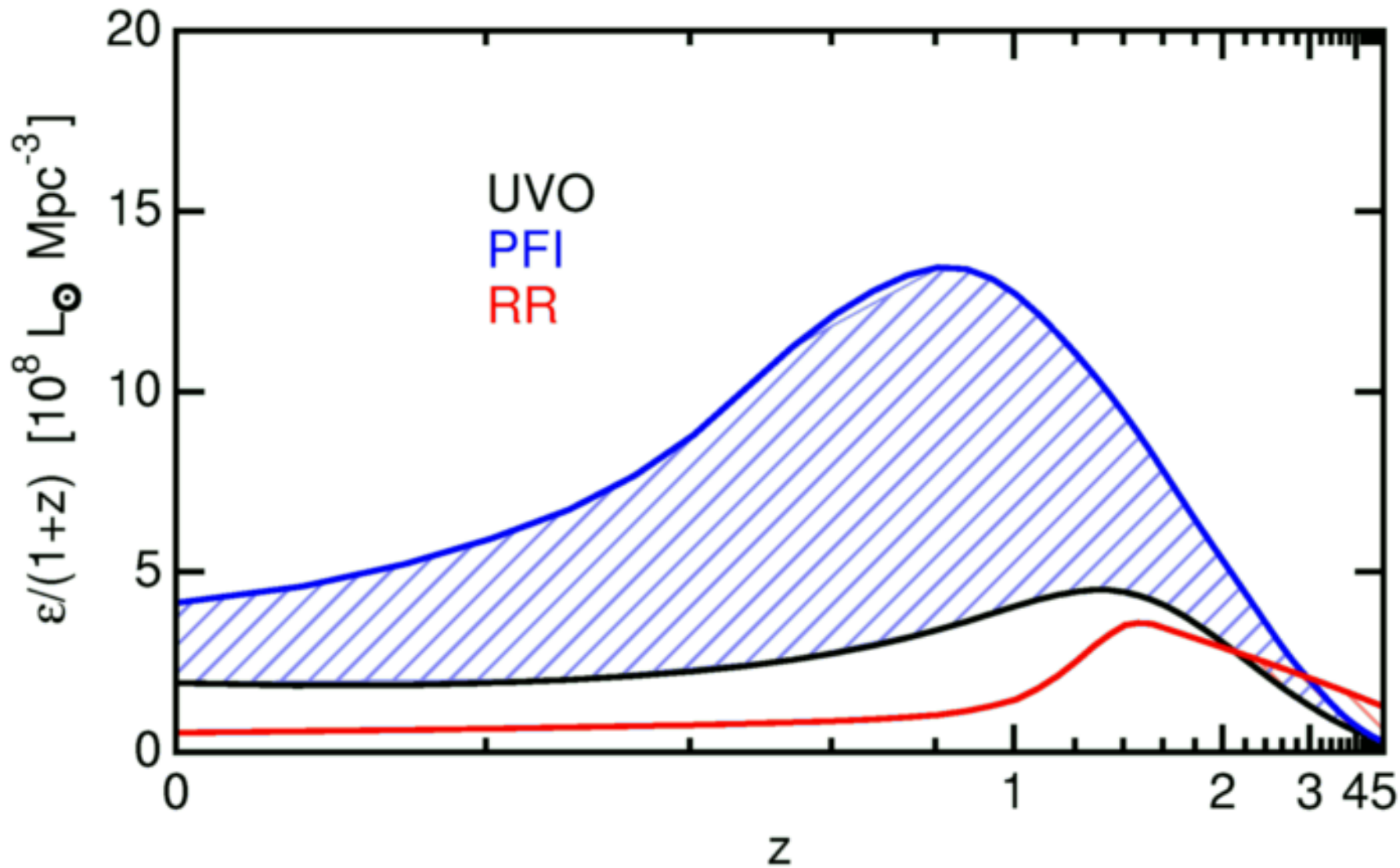
$$u = \frac{1}{H_o} \int \frac{(1+z)\ell(t_o)}{(1+z)^{3.5}} dz = \frac{1}{3} \int \ell(t) dt = \ell(t_o)t_o$$



# Luminosity density vs. redshift



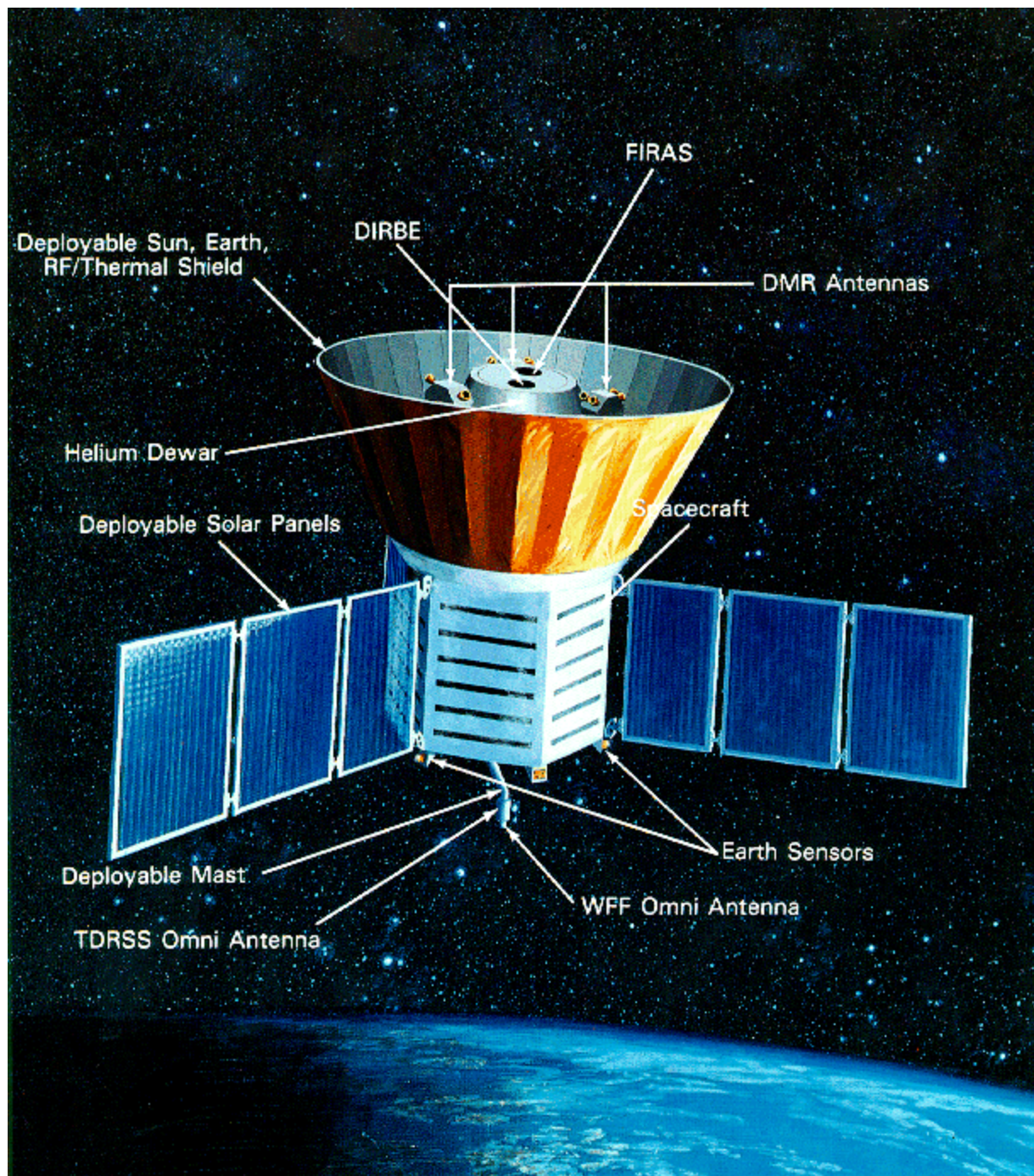
# $L/(1+z)$ vs. time



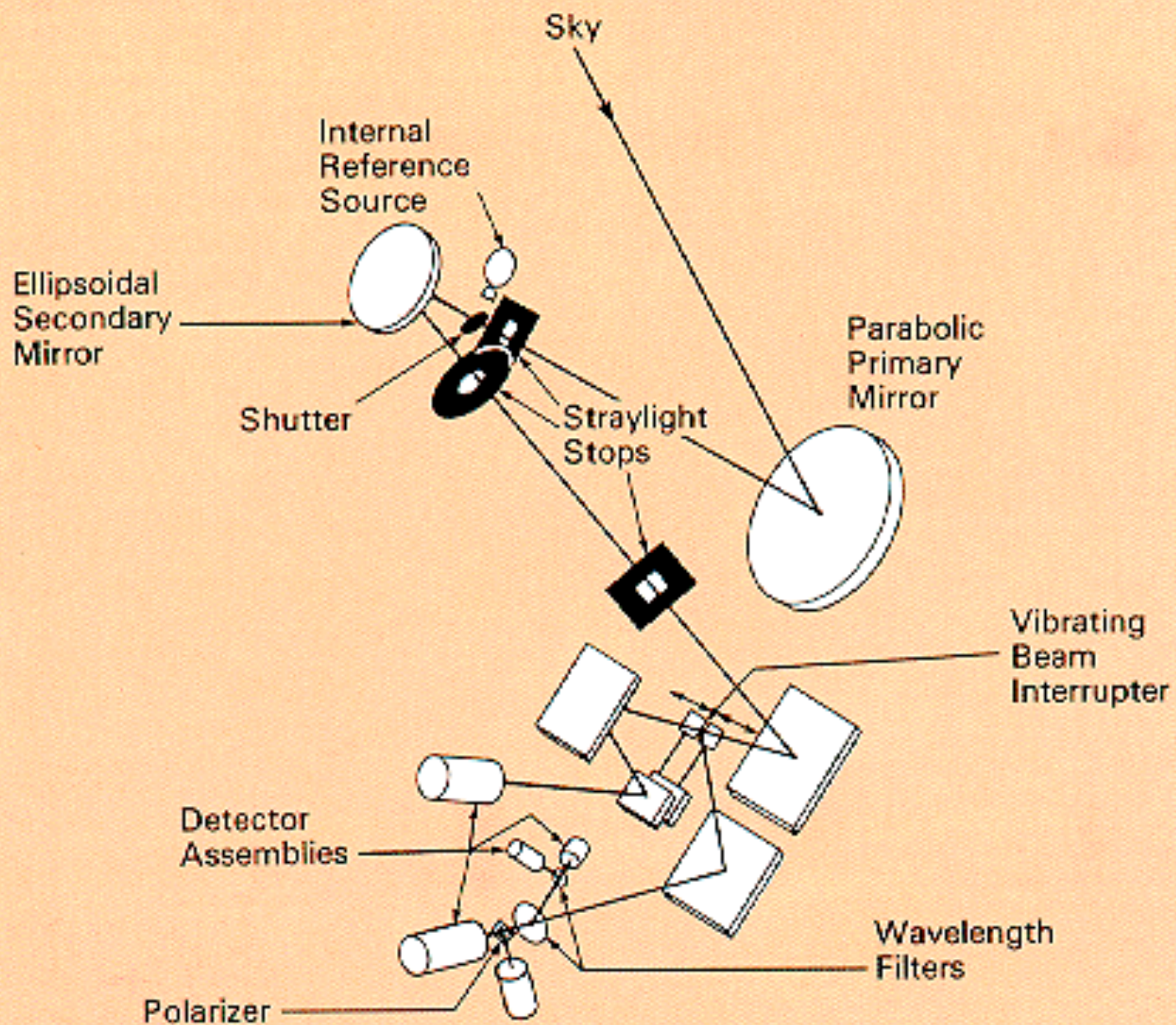
# Fluctuation Analysis

$$\begin{aligned}\text{var}(F) &= \int \left( \frac{L}{4\pi D_L(z)^2} \right)^2 (1+z)^3 n_{CM}(z) D_A(z)^2 c dt d\Omega \\ &= (4\pi)^{-2} \int \frac{n_{CM}(z) L^2}{(1+z) D_L(z)^2} c dt d\Omega\end{aligned}$$

- Diverges like  $1/z_{\min}$  or  $(F_{\text{cut}})^{0.5}$
- Always dominated by the brightest unmasked sources: galaxy edges, dwarf galaxies,  $z=1$  sources
- Contribution from high  $z$  highly suppressed
- Good for about 2X lower flux than catalog limit





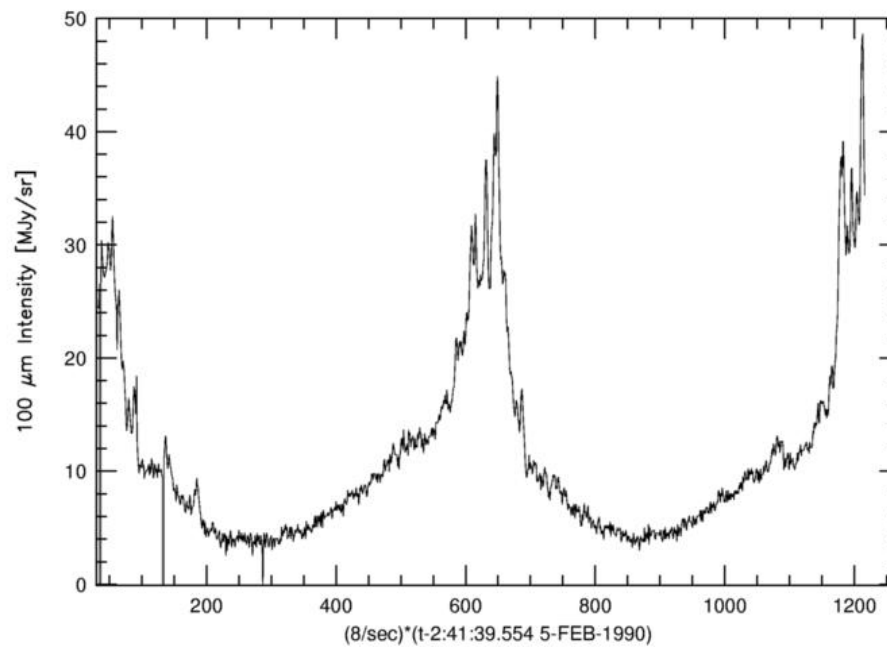
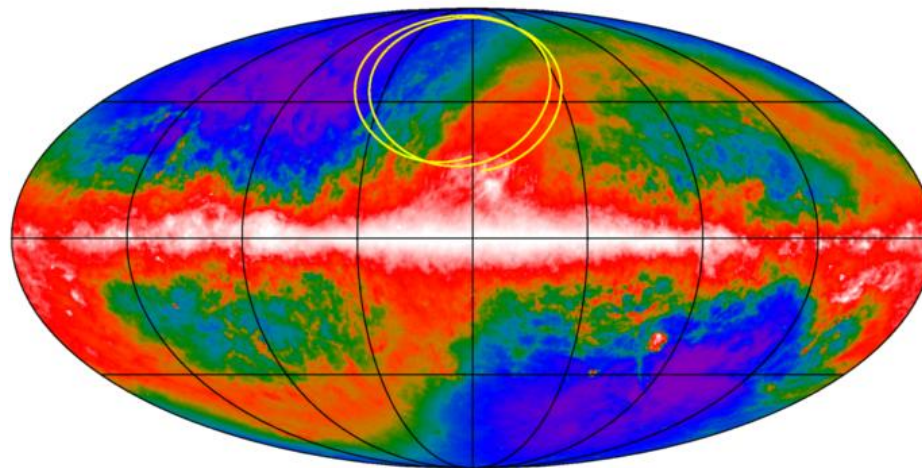


# DIRBE Beam Size

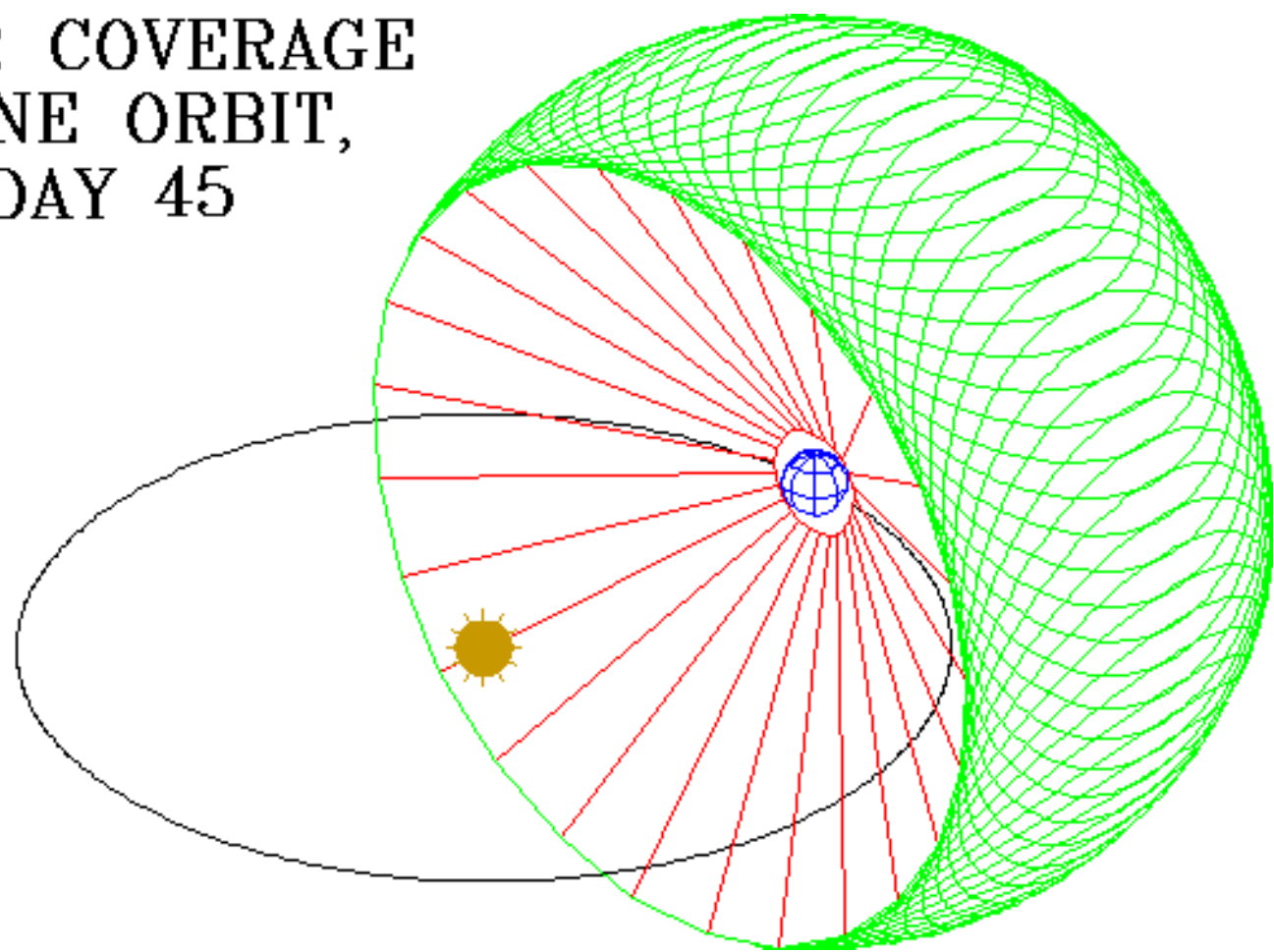


$0.7^\circ$

# Band 8 with 5-Feb-90 Scan

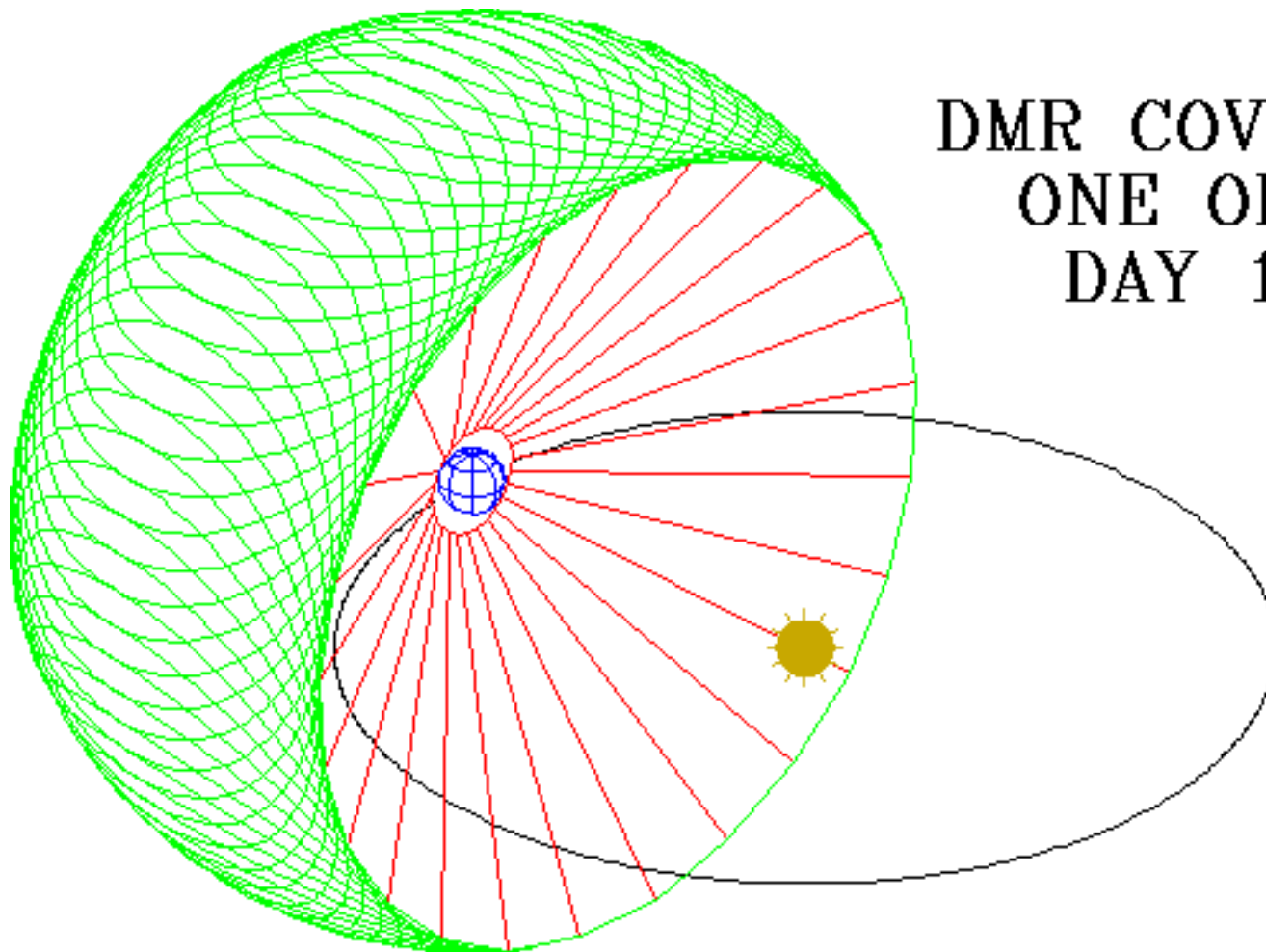


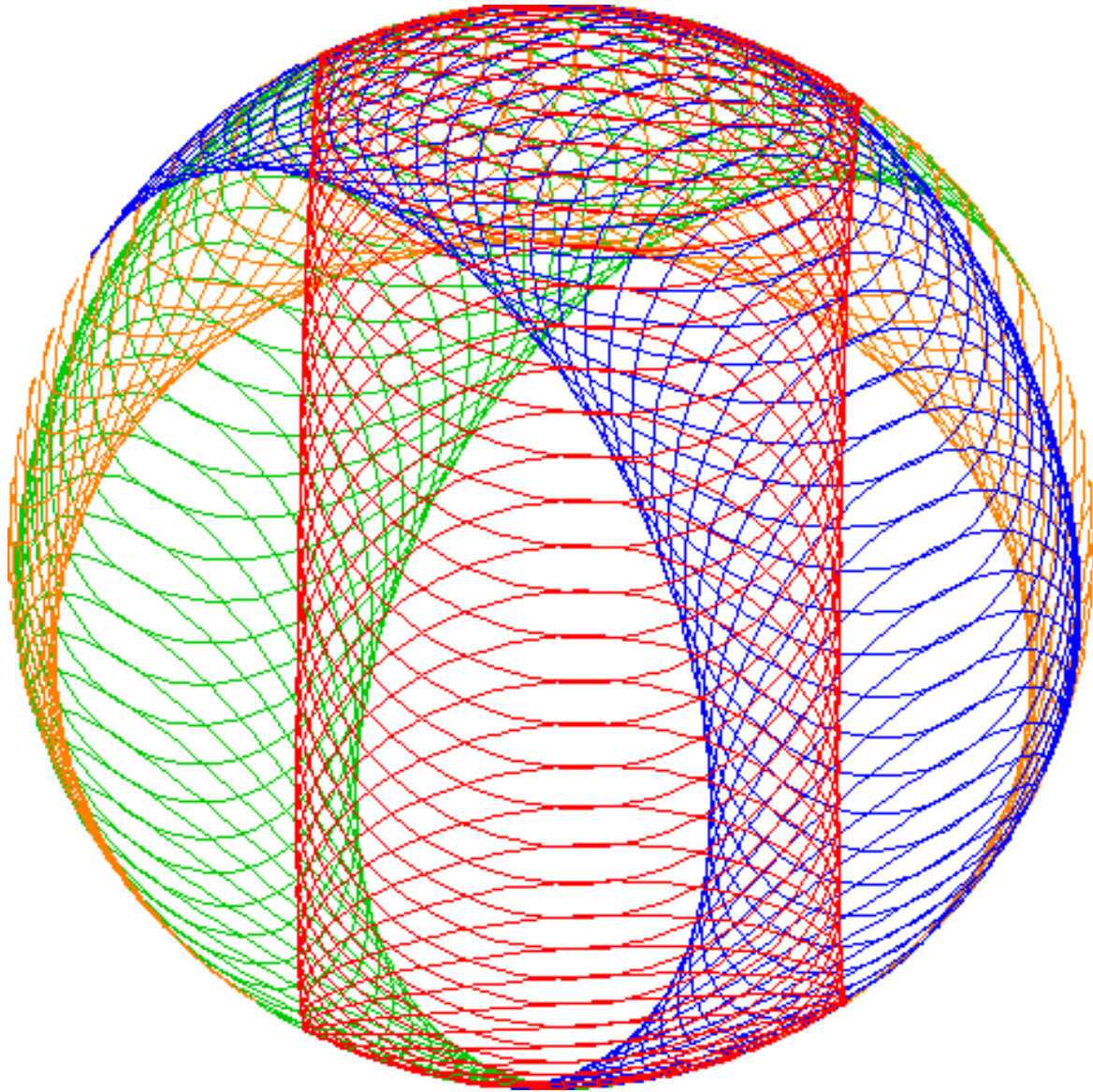
DMR COVERAGE  
ONE ORBIT,  
DAY 45

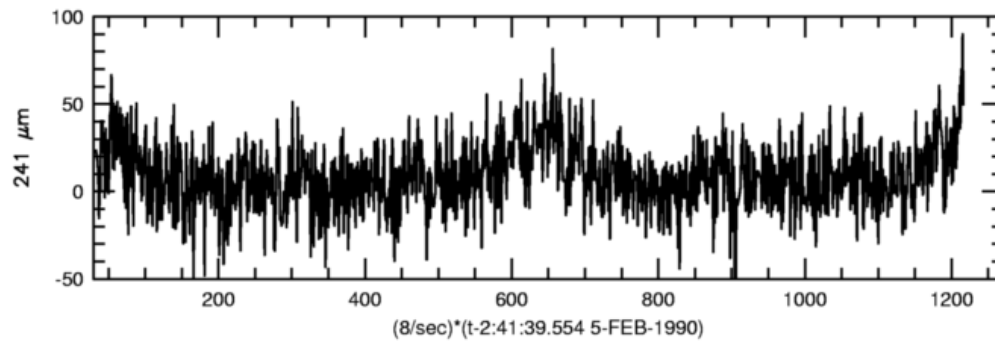
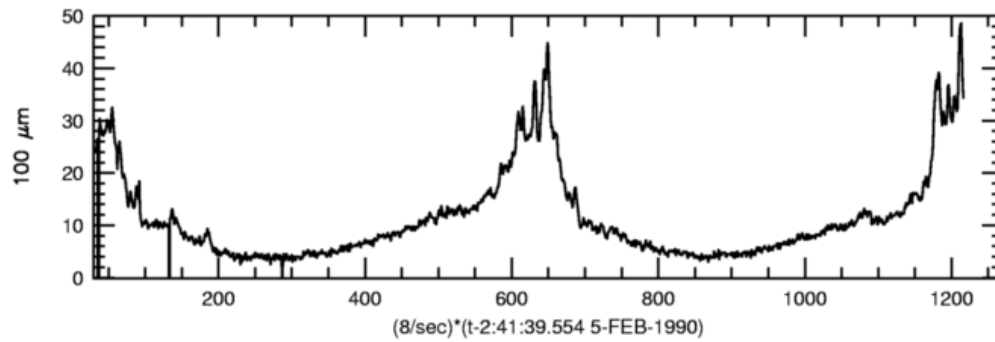
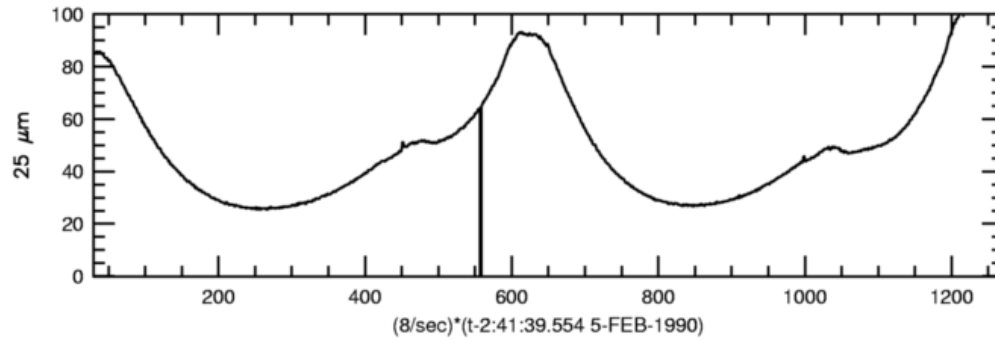
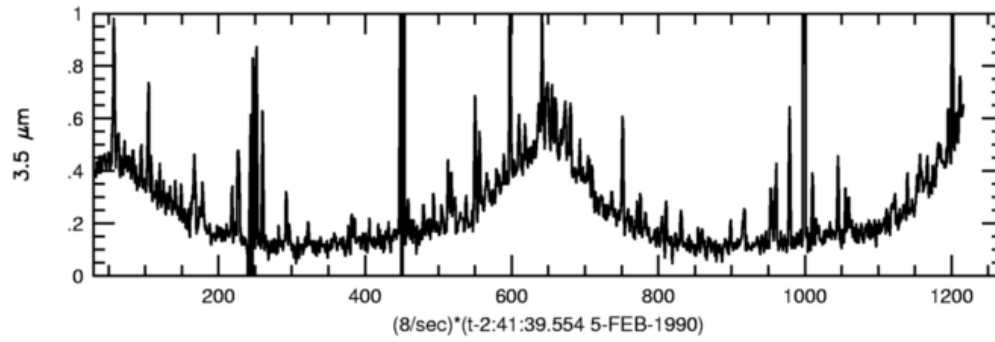




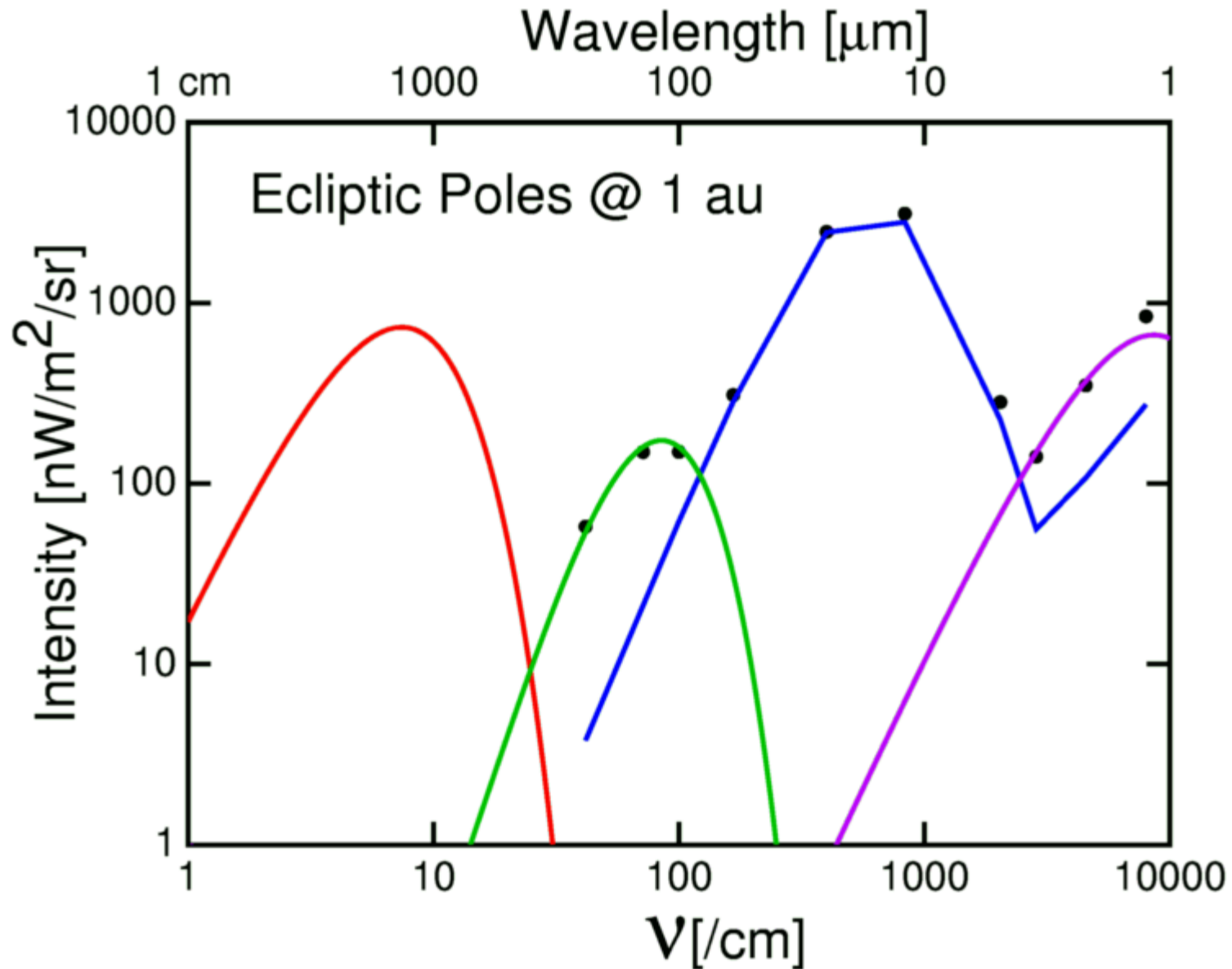
**DMR COVERAGE  
ONE ORBIT,  
DAY 135**



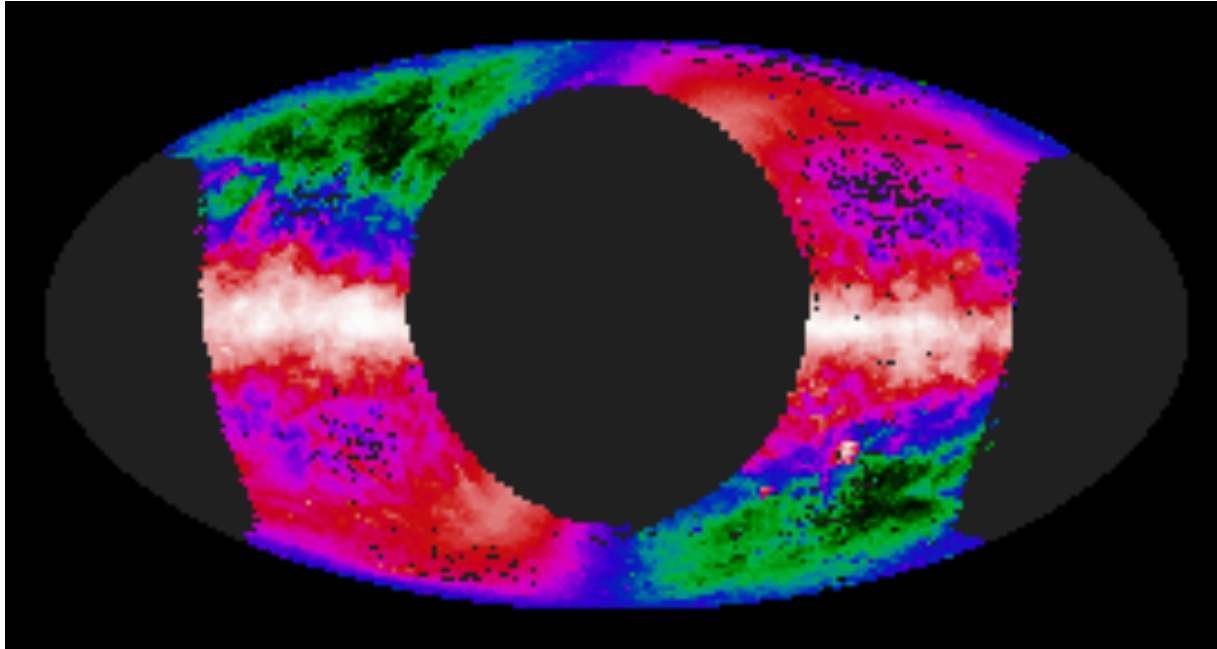




# Bump Chart: Where is the CIRB?

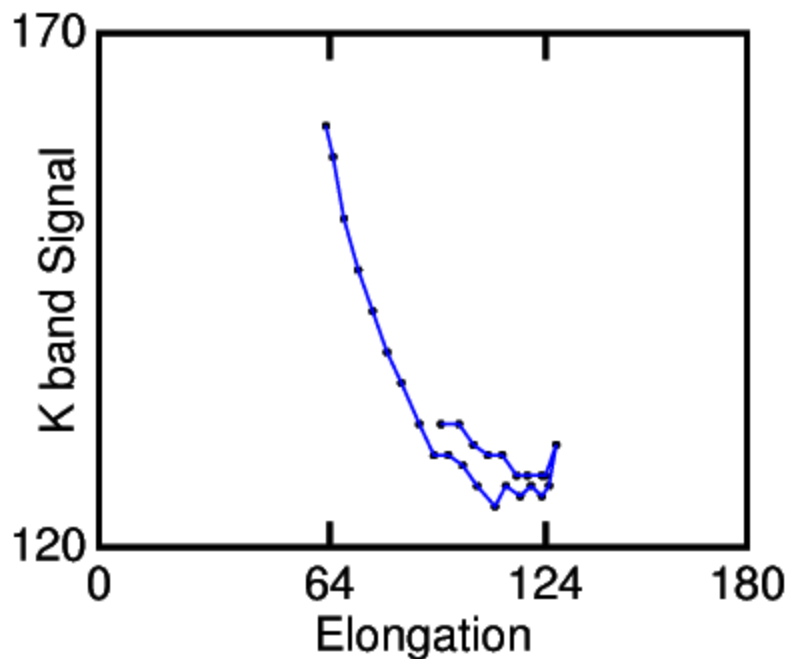
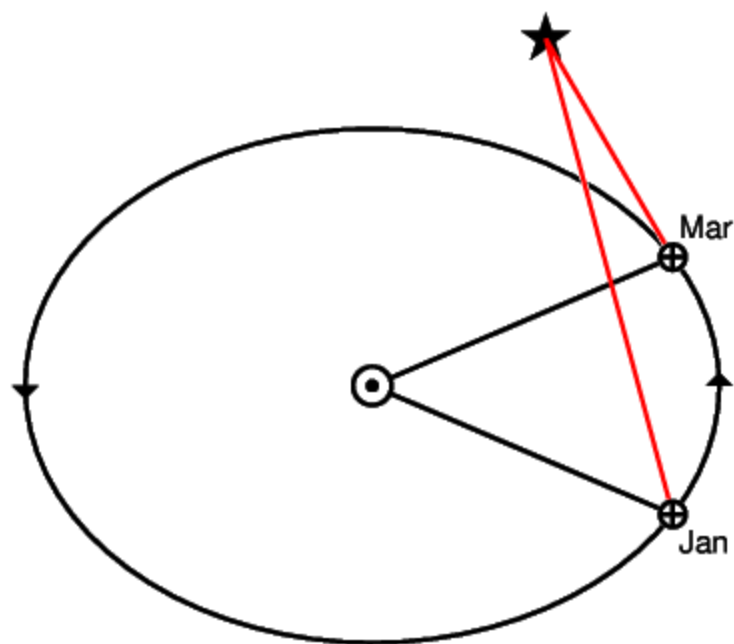


# DIRBE 100 $\mu\text{m}$ Weekly Maps



Note the triangles of zodiacal emission along the ecliptic on either side of the solar exclusion hole.

## Extrapolation to Outside Solar system?



- Observe same spot on the sky through different amounts of interplanetary dust.
- Fit a model to the change in intensity *vs.* elongation (or time).

## Interplanetary Dust Cloud Models

The observed intensity depends on a line-of-sight integral of an emissivity that depends on the dust density,  $n(x, y, z, t, \text{size, shape, composition})$ , a function of at least 7 variables. The data are given by  $I(\nu, l, b, t)$ , a function of only 4 variables.

Simplification is needed. For a time independent solution, the density depends only on the integrals of the motion:  $E$ ,  $L$ , and  $L_z$  which map into  $a$ ,  $e$  and  $i$ . The eccentricity has little effect on the line-of-sight integral, so assume the density is function only of  $r$  and  $i$ .

Let  $n(r, i) = R_1(r)I_1(i) + R_2(r)I_2(i) + \dots$

## Radial Density Profile

The Poynting-Robertson effect gives a torque of

$$T = -\sigma \frac{r v_{orb}}{c^2} \frac{L}{4\pi r^2}$$

and since the angular momentum is  $L = m\sqrt{GMr}$  this gives an inward drift rate of

$$\frac{dr}{dt} = \frac{T}{dL/dr} = -(2\pi)^{-1} \frac{\sigma}{m} \frac{L}{c^2} r^{-1}$$

A steady state requires that

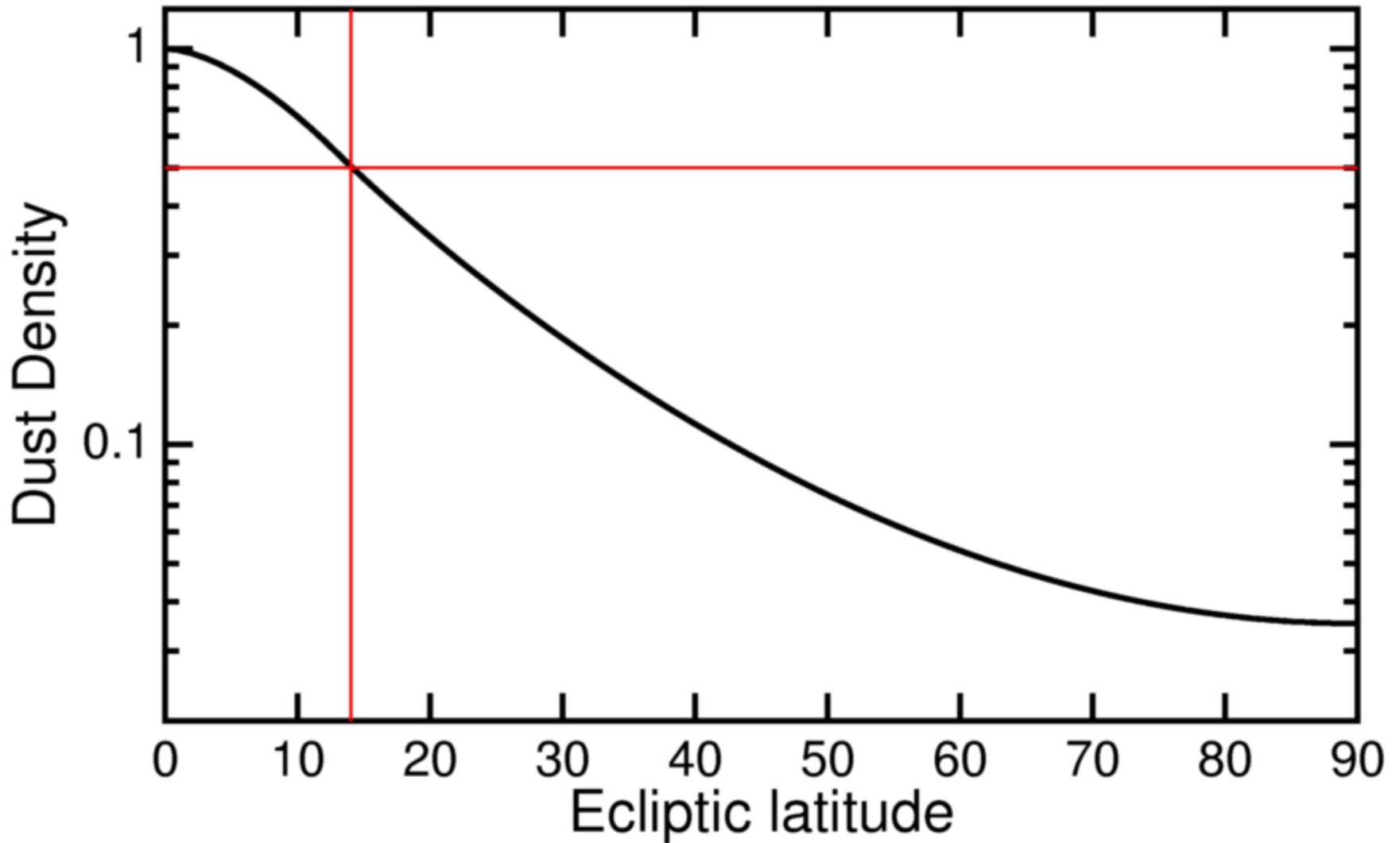
$$\frac{\partial}{\partial r} \left( n(r) r^2 \frac{dr}{dt} \right) = 0$$

and thus that  $n(r) \propto r^{-1}$ .

FIT gives  $r^{-1.23}$

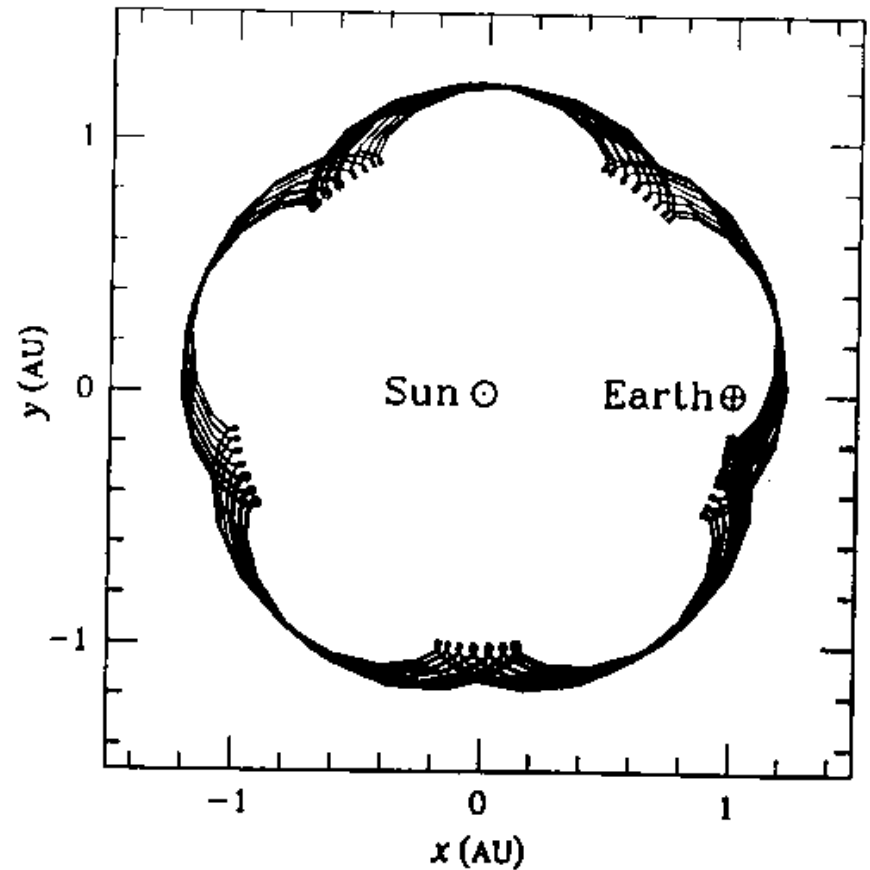
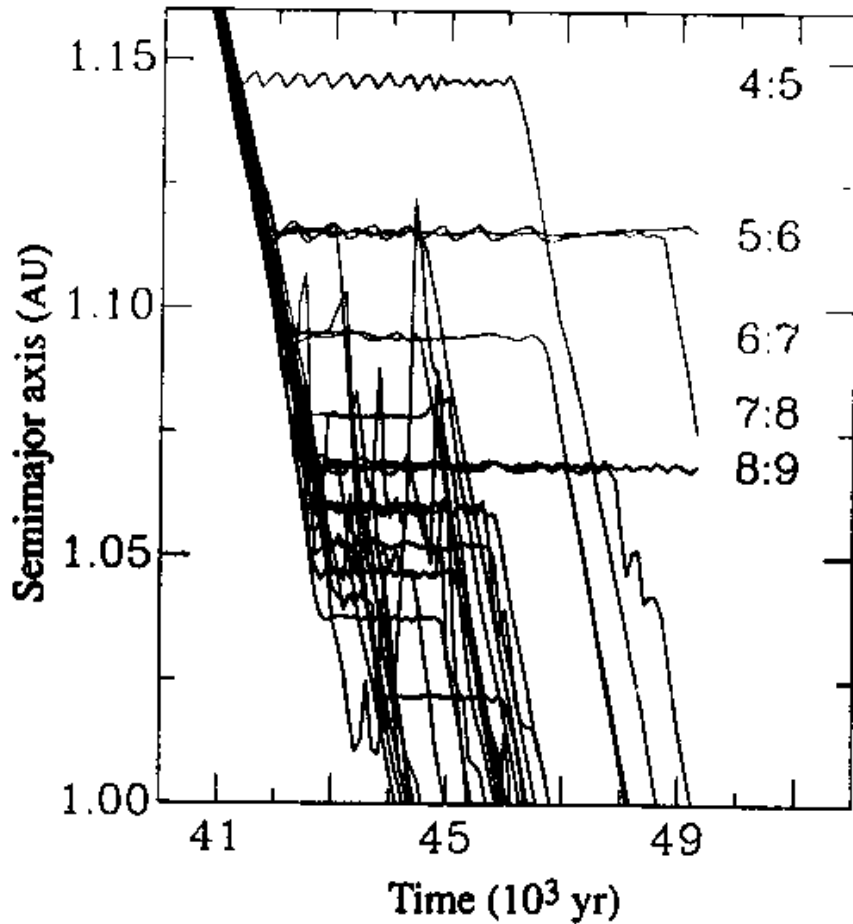


# Dust Density vs inclination

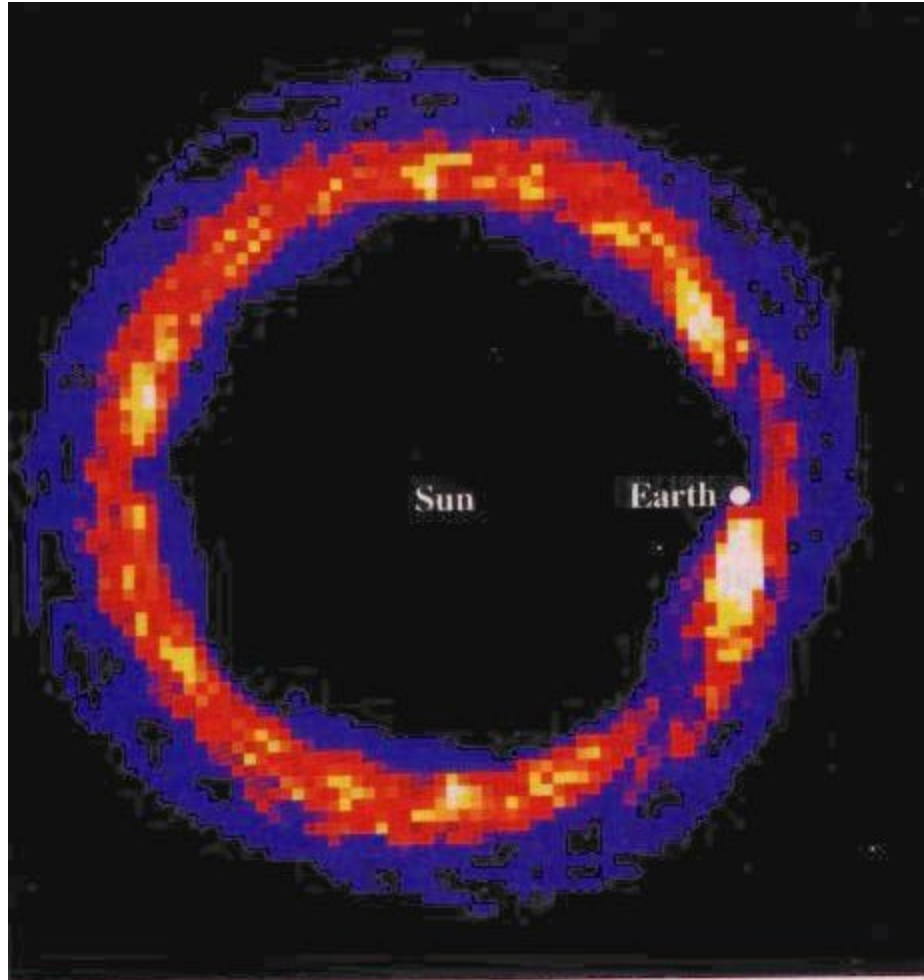


# A Circumsolar Ring of Asteroidal Dust in Resonant Lock with the Earth

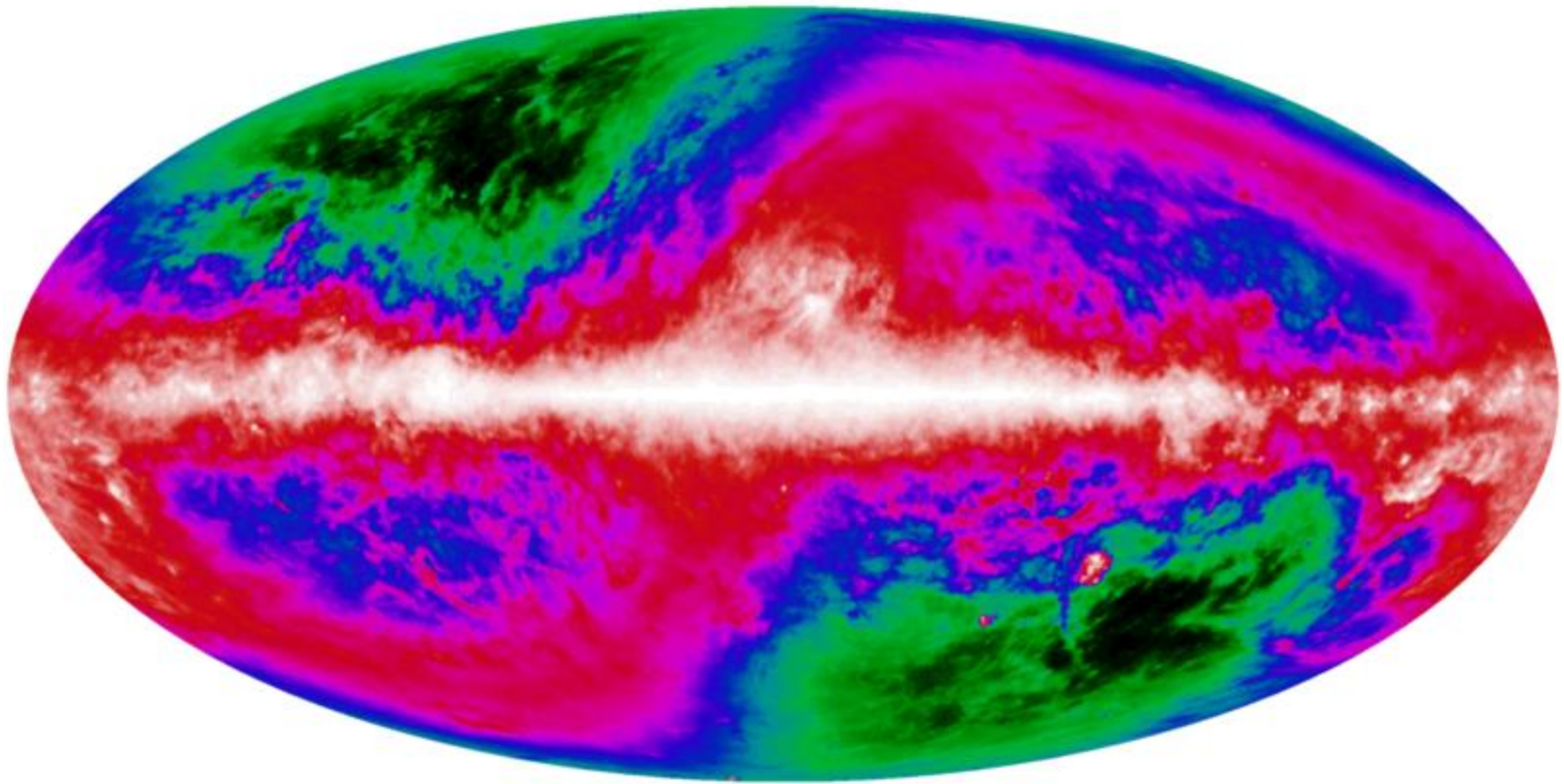
Dermott et al., 1994, Nature, 369, 719



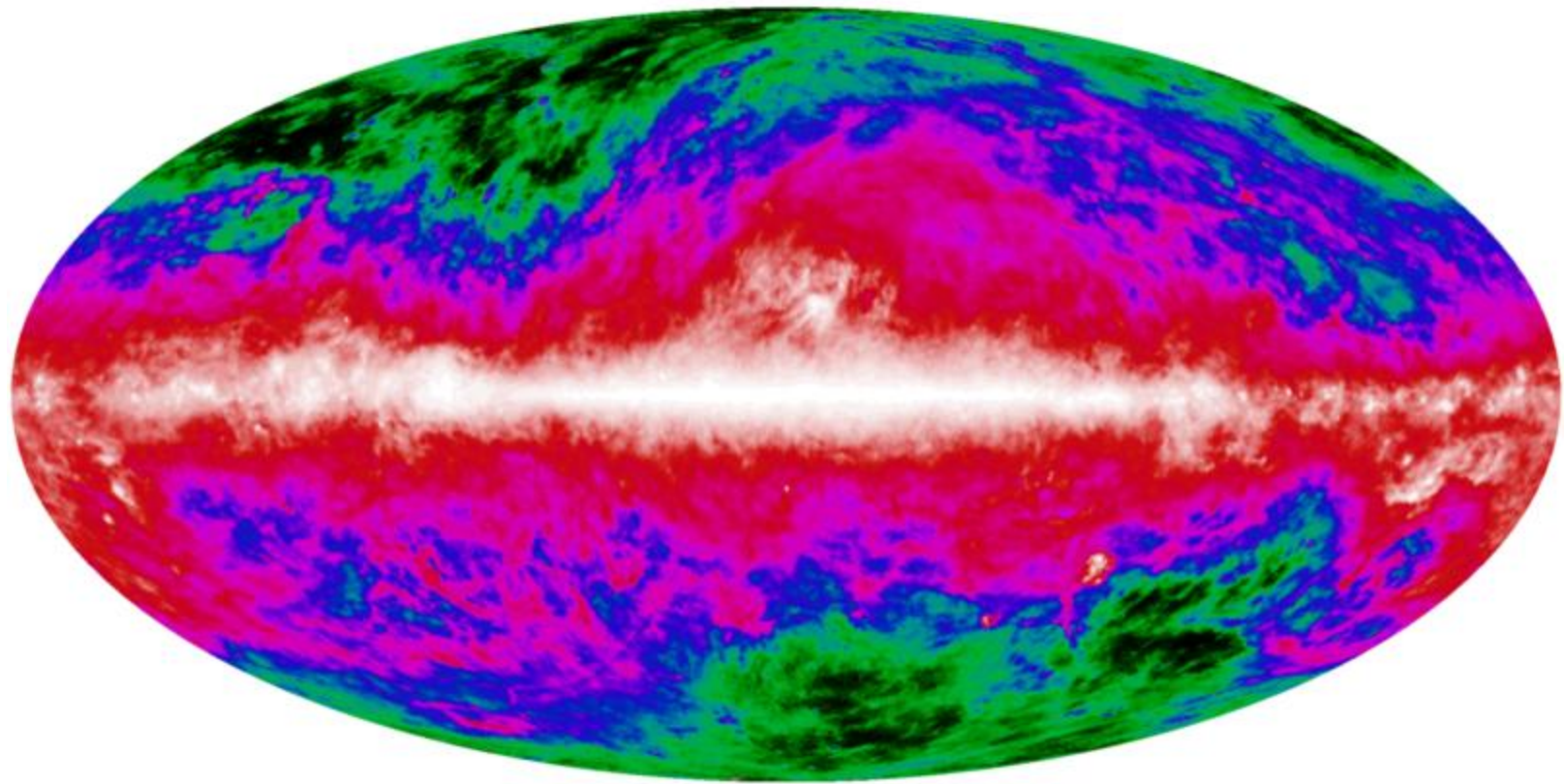
# A Ring and a Trailing Clump



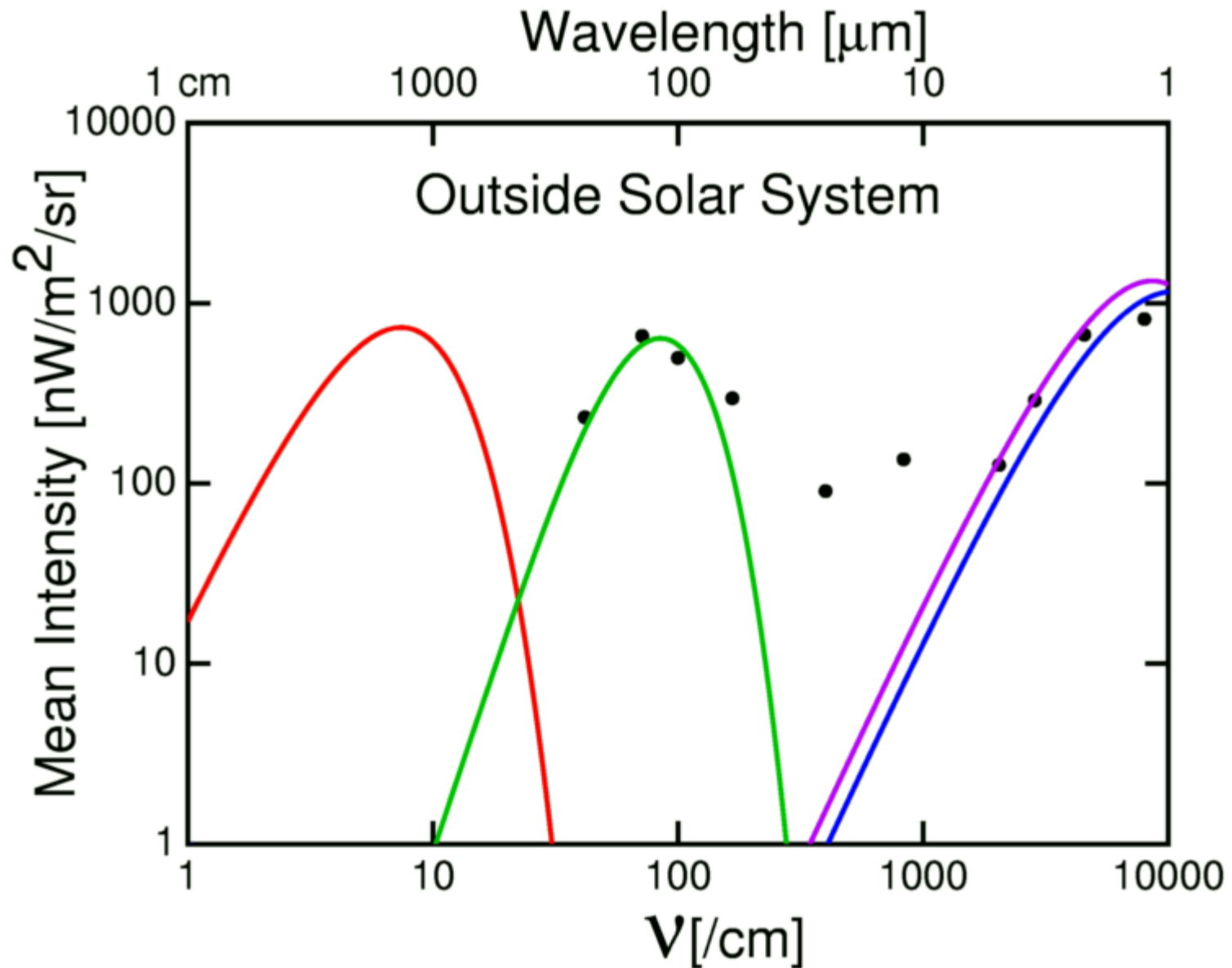
# 100 Micron Total



# Zodi Subtracted 100 Micron

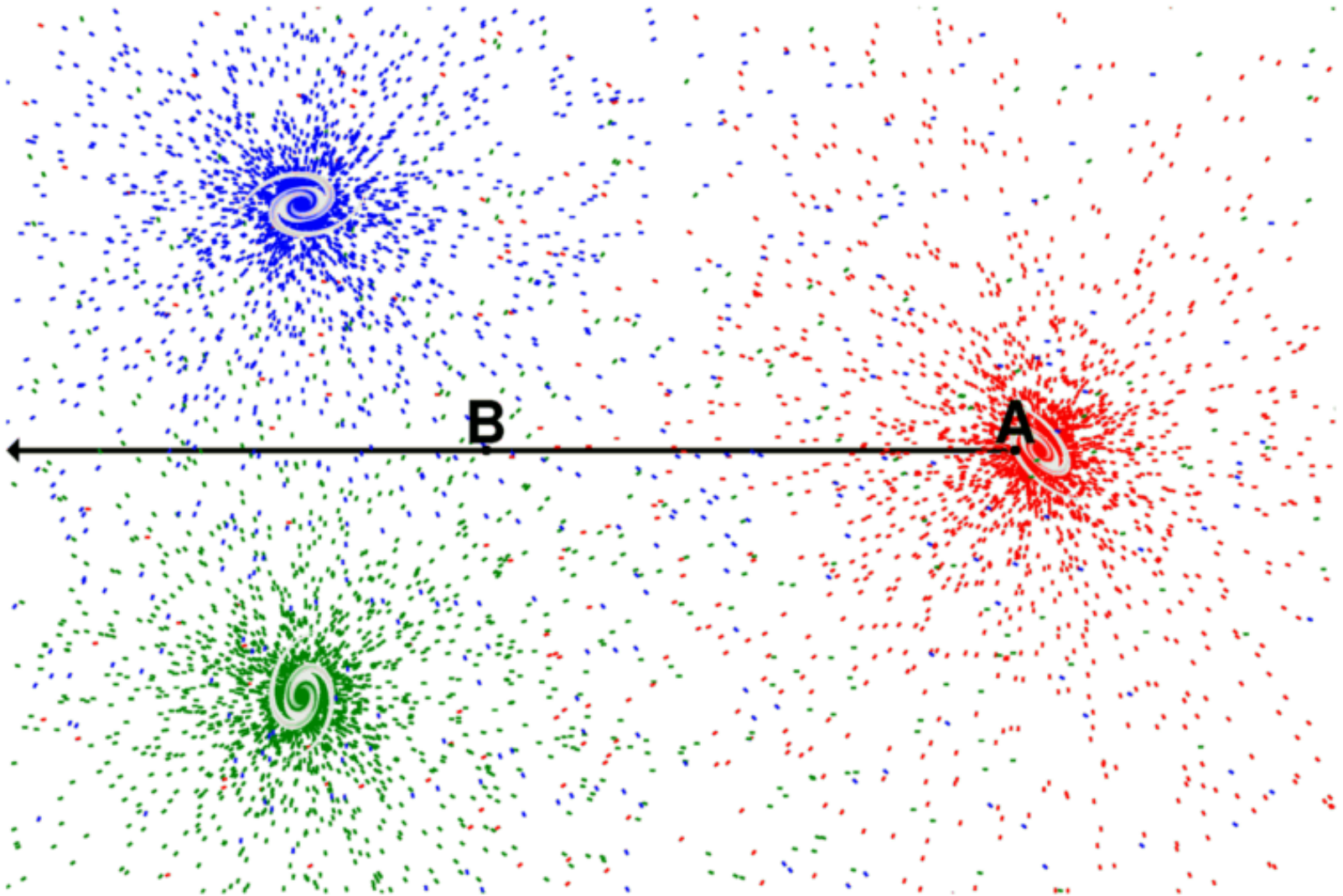


# Still no CIRB Bump:

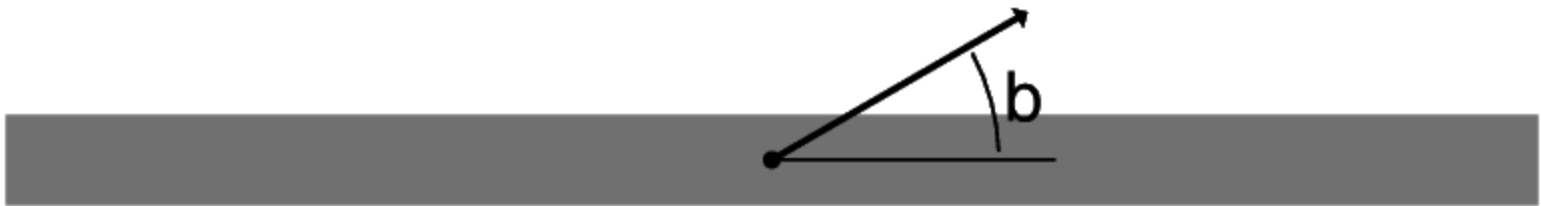




We want  $vJ_v$  at B but sit at A



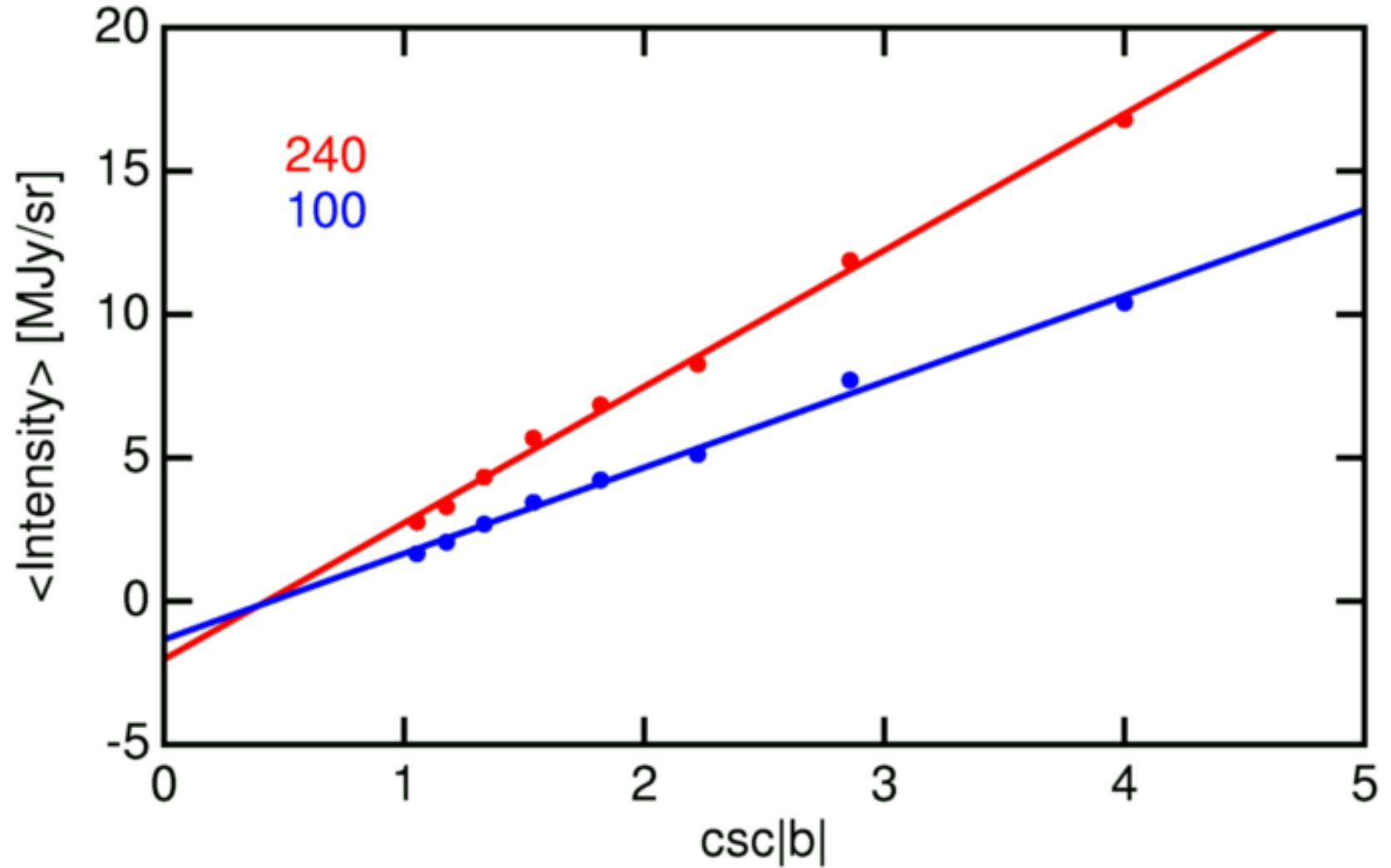
## Extrapolation to NO Galaxy?



- Galaxy is a *very* thin disk
- Average column density  $\propto \csc |b|$

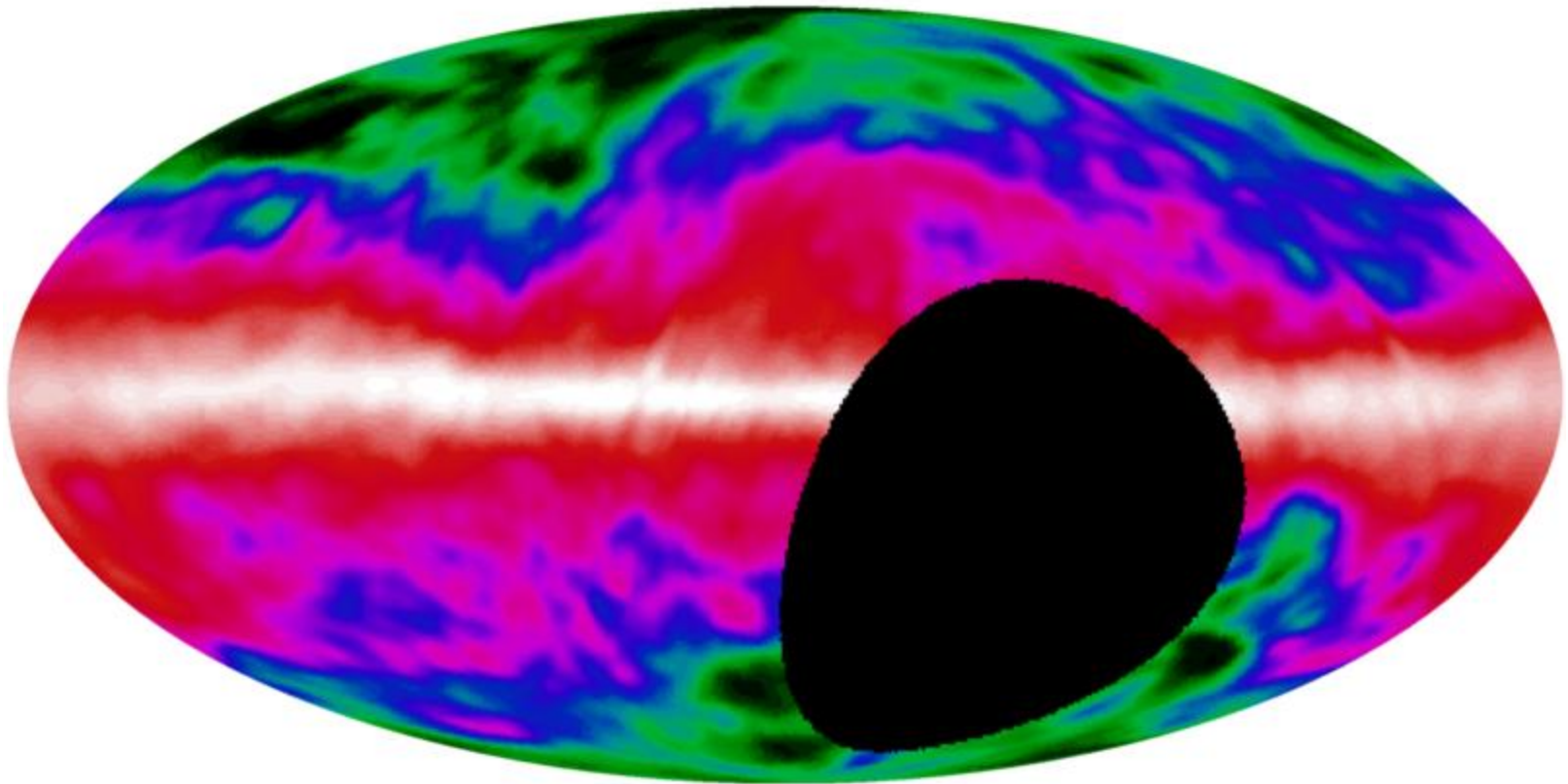


# Extrapolation to $\csc|b|=0$ in Far IR



# Atomic Hydrogen Map

21 cm H I emission

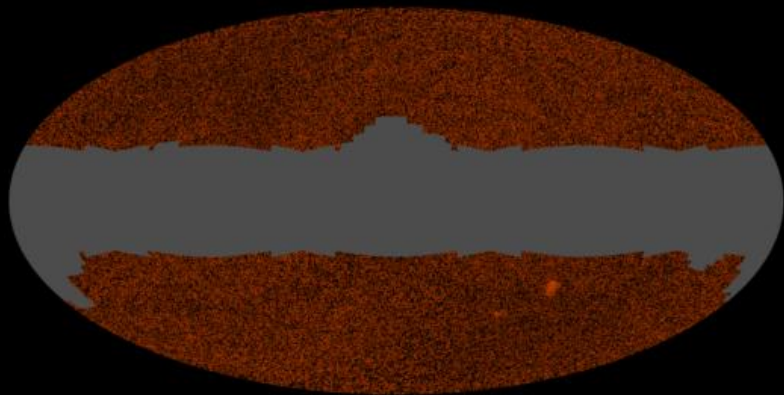
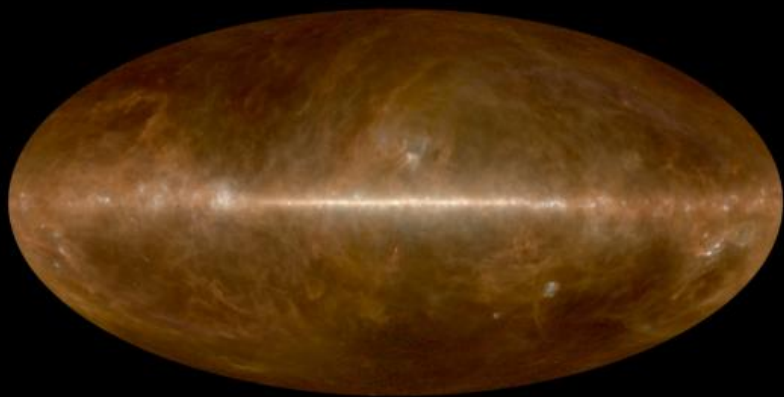


“The COBE Diffuse Infrared Background Experiment Search for the Cosmic Infrared Background. I. Limits and Detections”, Hauser *et al.* (1998, ApJ, 508, 25)

“The COBE Diffuse Infrared Background Experiment Search for the Cosmic Infrared Background. II. Model of the Interplanetary Dust Cloud”, Kelsall *et al.* (1998, ApJ, 508, 44)

“The COBE Diffuse Infrared Background Experiment Search for the Cosmic Infrared Background. III. Separation of Galactic Emission from the Infrared Sky Brightness”, Arendt *et al.* (1998, ApJ, 508, 74)

“The COBE Diffuse Infrared Background Experiment Search for the Cosmic Infrared Background. IV. Cosmological Implications”, Dwek *et al.* (1998, ApJ, 508, 106)



## DIRBE Team IRB Results

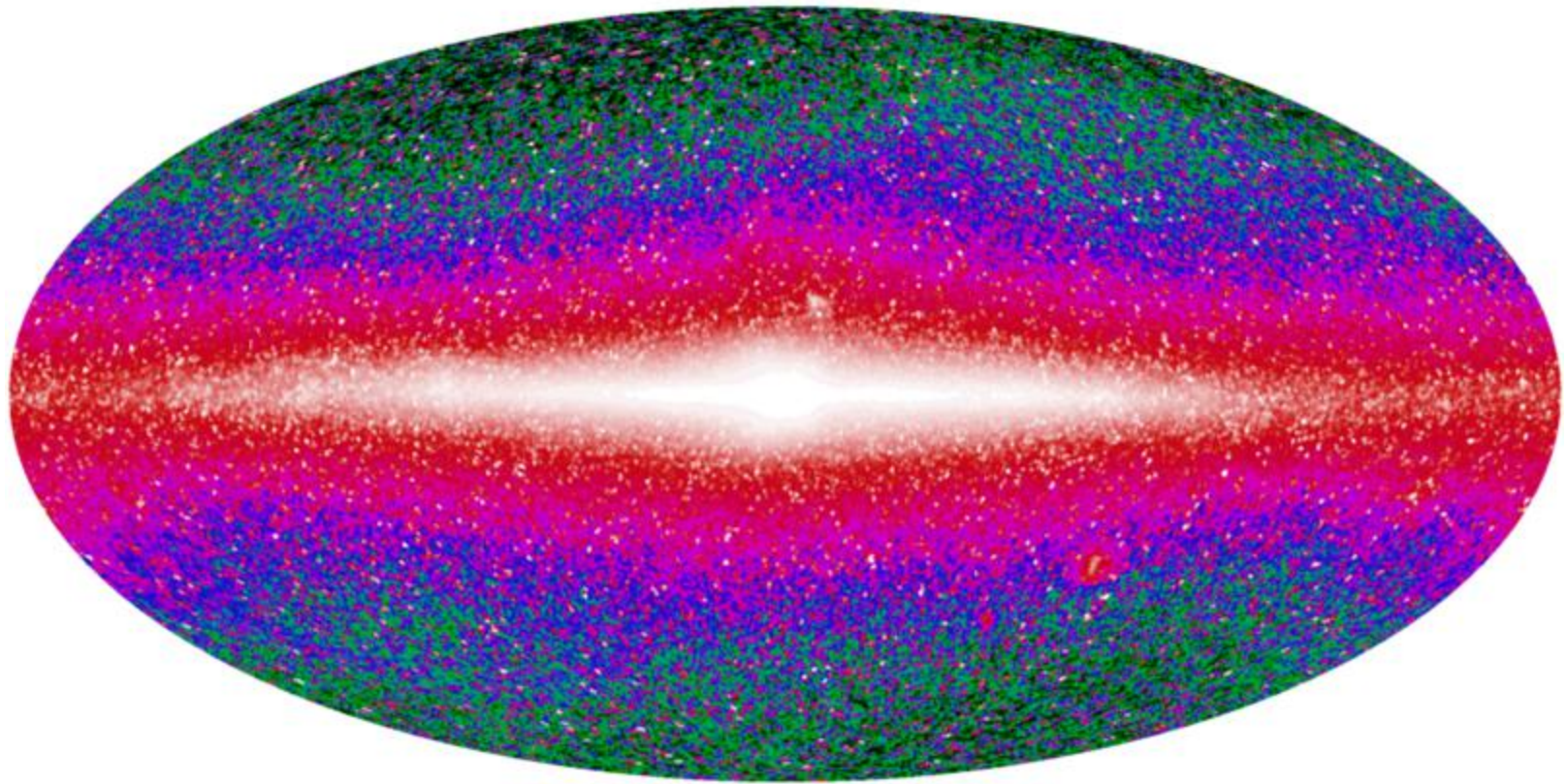
Hauser, . . . , Wright (1998, ApJ, 508, 25)

$\lambda$ [ $\mu\text{m}$ ]	$\nu I_\nu$ [ $\text{nW m}^{-2} \text{sr}^{-1}$ ]	$I_\nu$ [ $\text{MJy sr}^{-1}$ ]
1.25	< 75.	< 0.031
2.2	< 39.	< 0.029
3.5	< 23.	< 0.027
4.9	< 41.	< 0.067
12.	< 468.	< 1.87
25.	< 504.	< 4.2
60.	< 75.	< 1.5
100.	< 38.	< 1.27
140.	$25.0 \pm 6.9$	$1.17 \pm 0.32$
240.	$13.6 \pm 2.5$	$1.09 \pm 0.2$

$$I_\nu \approx (1.3 \pm 0.4) \times 10^{-5} (\tilde{\nu}/100)^{0.64 \pm 0.12} B_\nu(18.5 \pm 1.2 \text{ K})$$

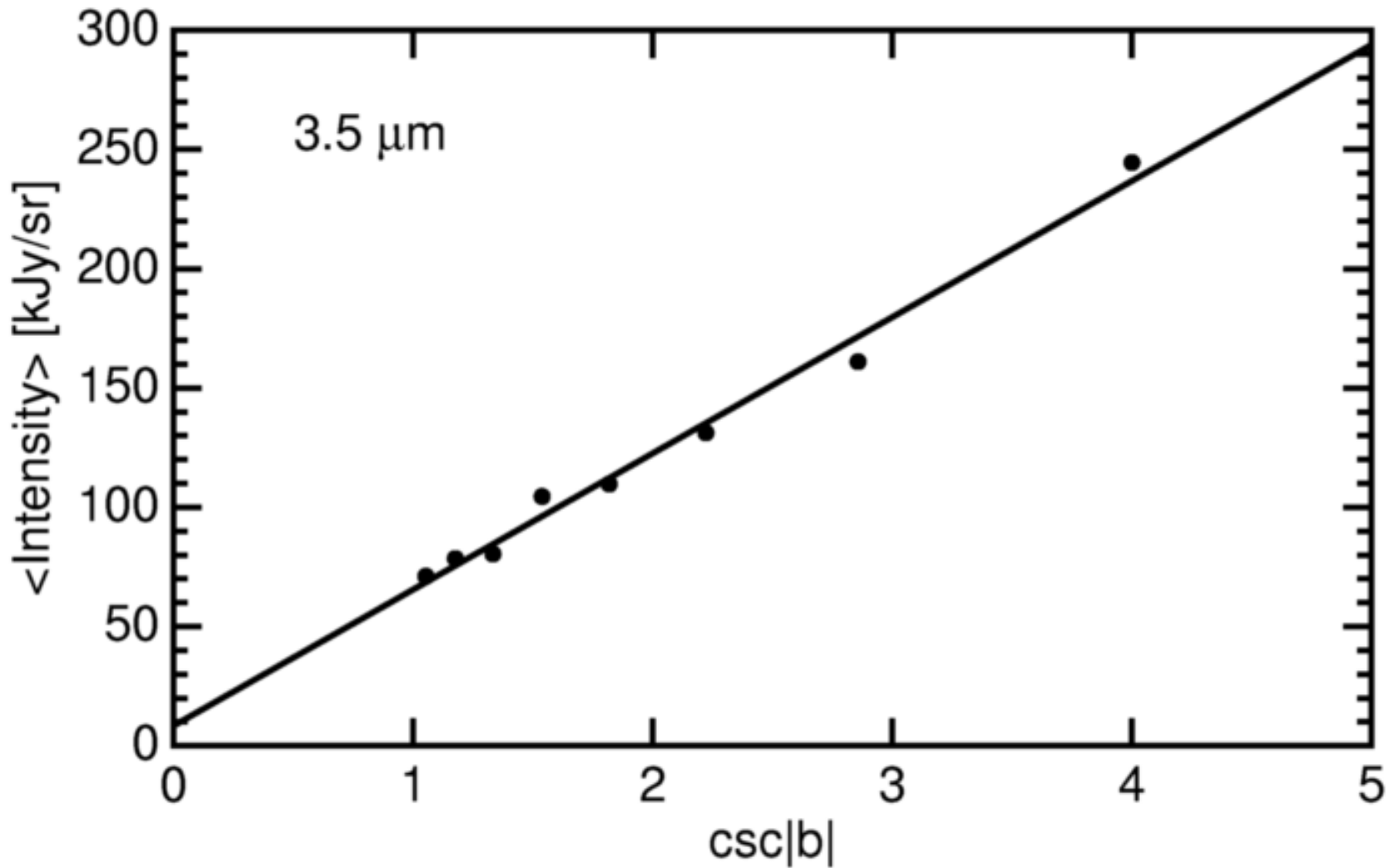
Fixsen *et al.*, 1998, ApJ, 508, 123

# Zodi Subtracted 3.5 Microns





# Extrapolation to $\csc|b|=0$ at $3.5 \mu\text{m}$



UNIVERSITY OF CALIFORNIA  
LOS ANGELES

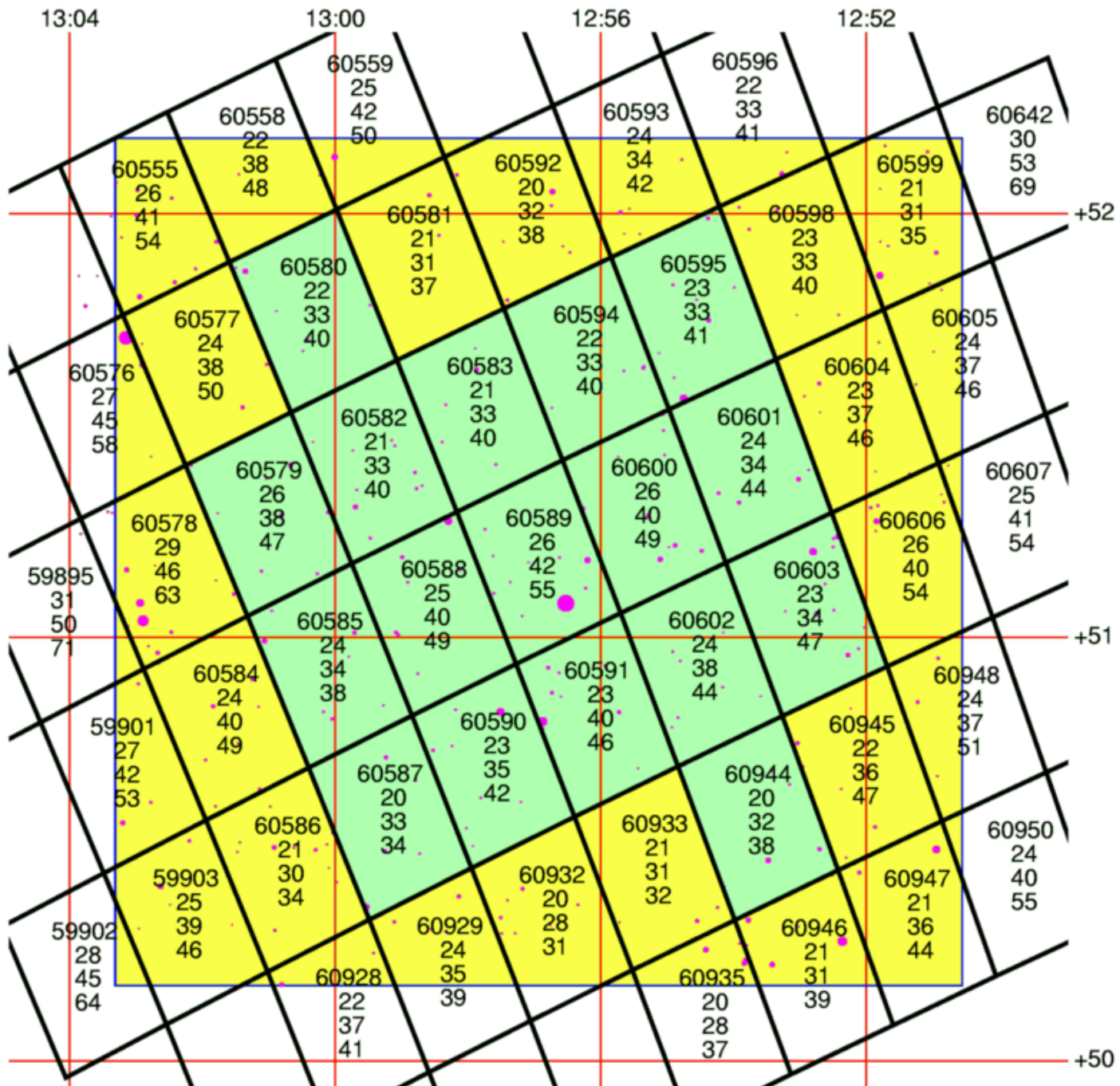
Detecting the Cosmic Infrared Background at  $2.2\mu\text{m}$   
with Ground Based and Space Based Observations

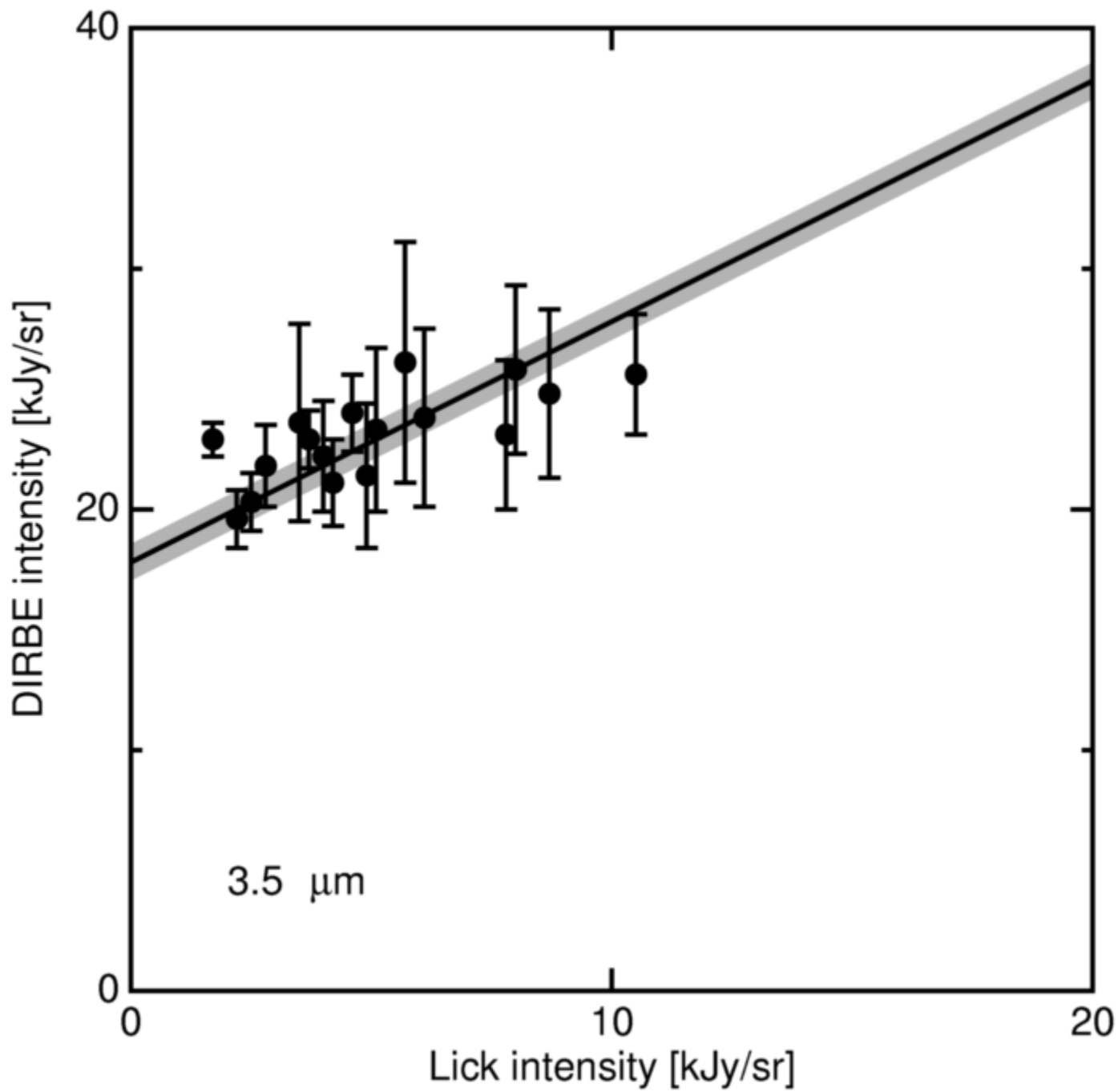
A dissertation submitted in partial satisfaction of the  
requirements for the degree  
Doctor of Philosophy in Physics and Astronomy

by

Varoujan Gorjian  
1998

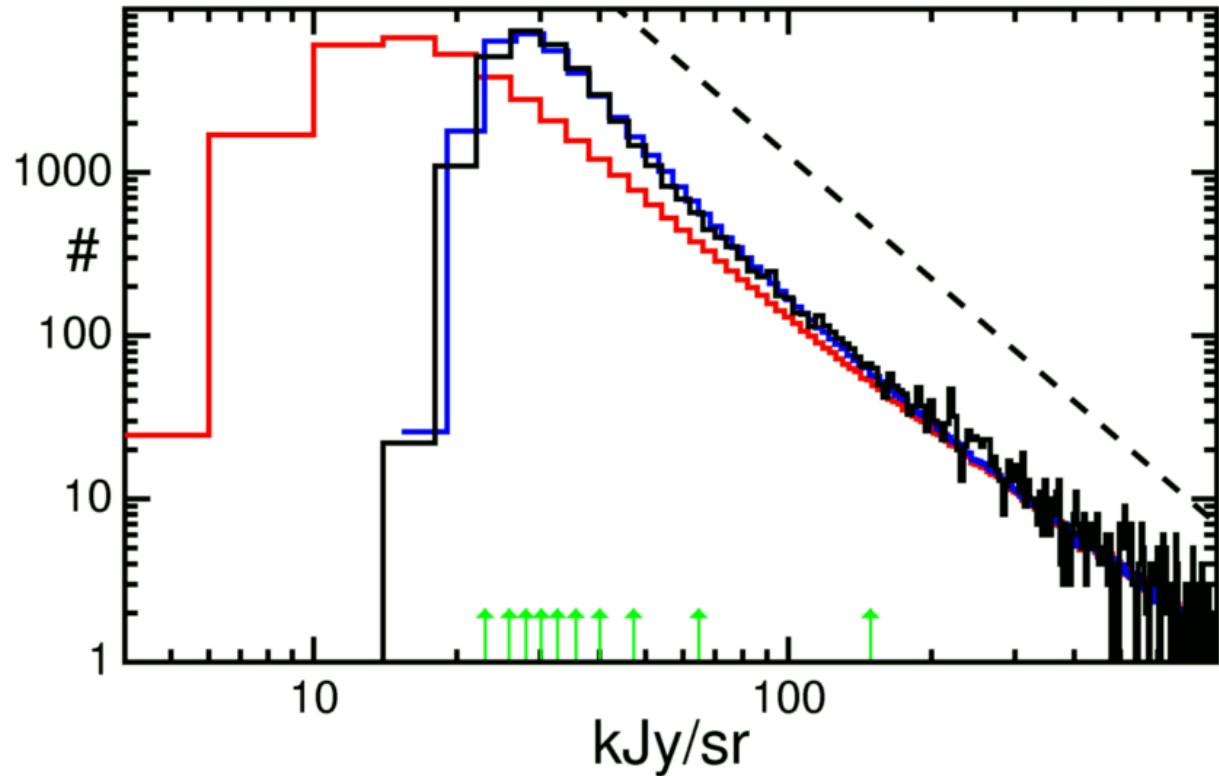






# Wright & Reese, 2000

- Generated many fake star fields,  $|b| > 64^\circ$
- Properly allocated fluxes to DIRBE pixels
- **Resulting model histogram**
- **Histogram shifted by 14.4 kJy/sr**
- Actual zodi-subtracted DIRBE data
- Dashed: Euclidean



$$\text{CIBR} = 23.1 \pm 5.9 \text{ kJy/sr at } 2.2 \mu\text{m}$$
$$14.4 \pm 3.7 \text{ kJy/sr at } 3.5 \mu\text{m}$$

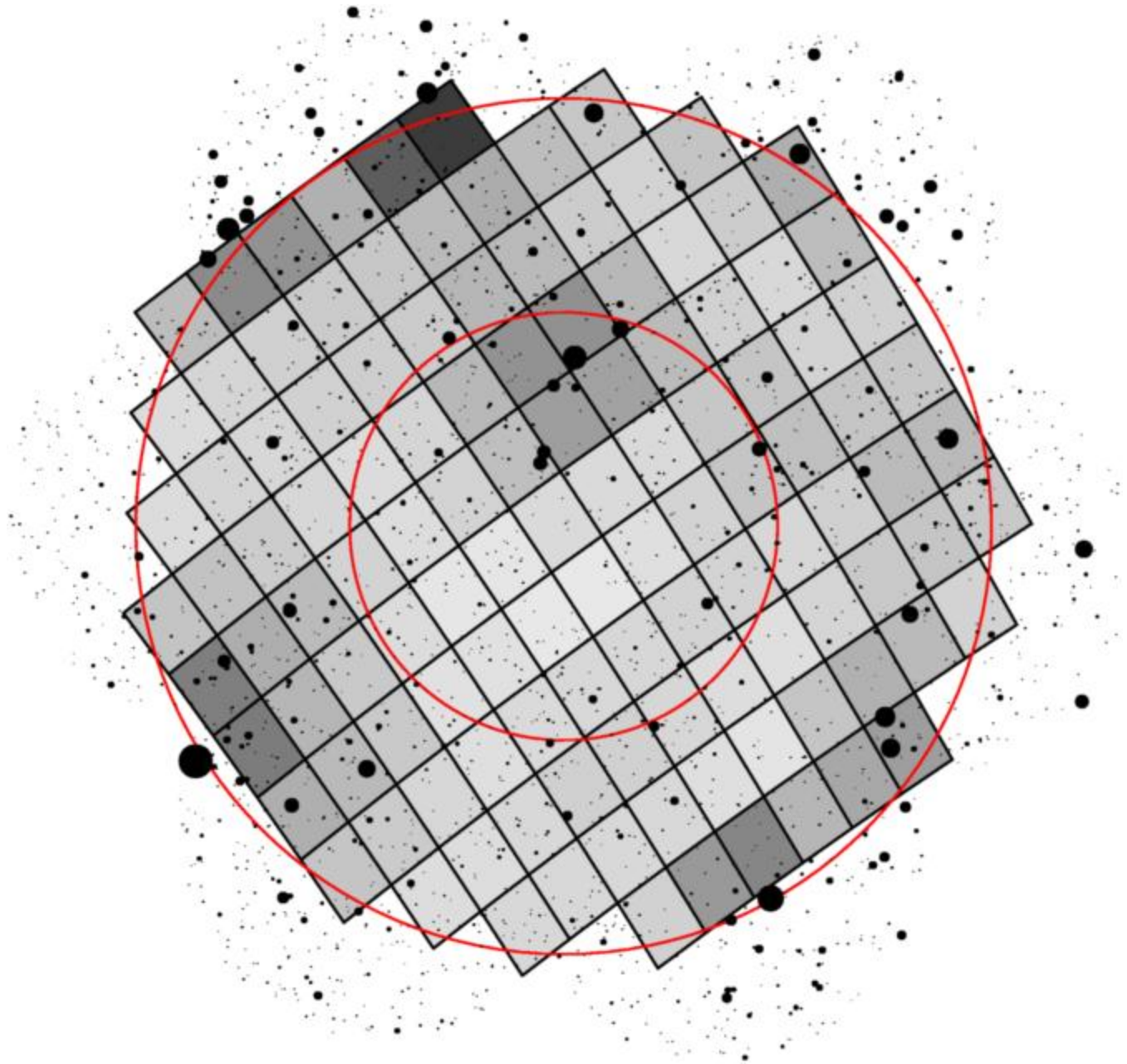
Levenson  
2008 PhD

Intensity, Isotropy and Origin  
of the Cosmic Infrared Background

A dissertation submitted in partial satisfaction  
of the requirements for the degree  
Doctor of Philosophy in Physics

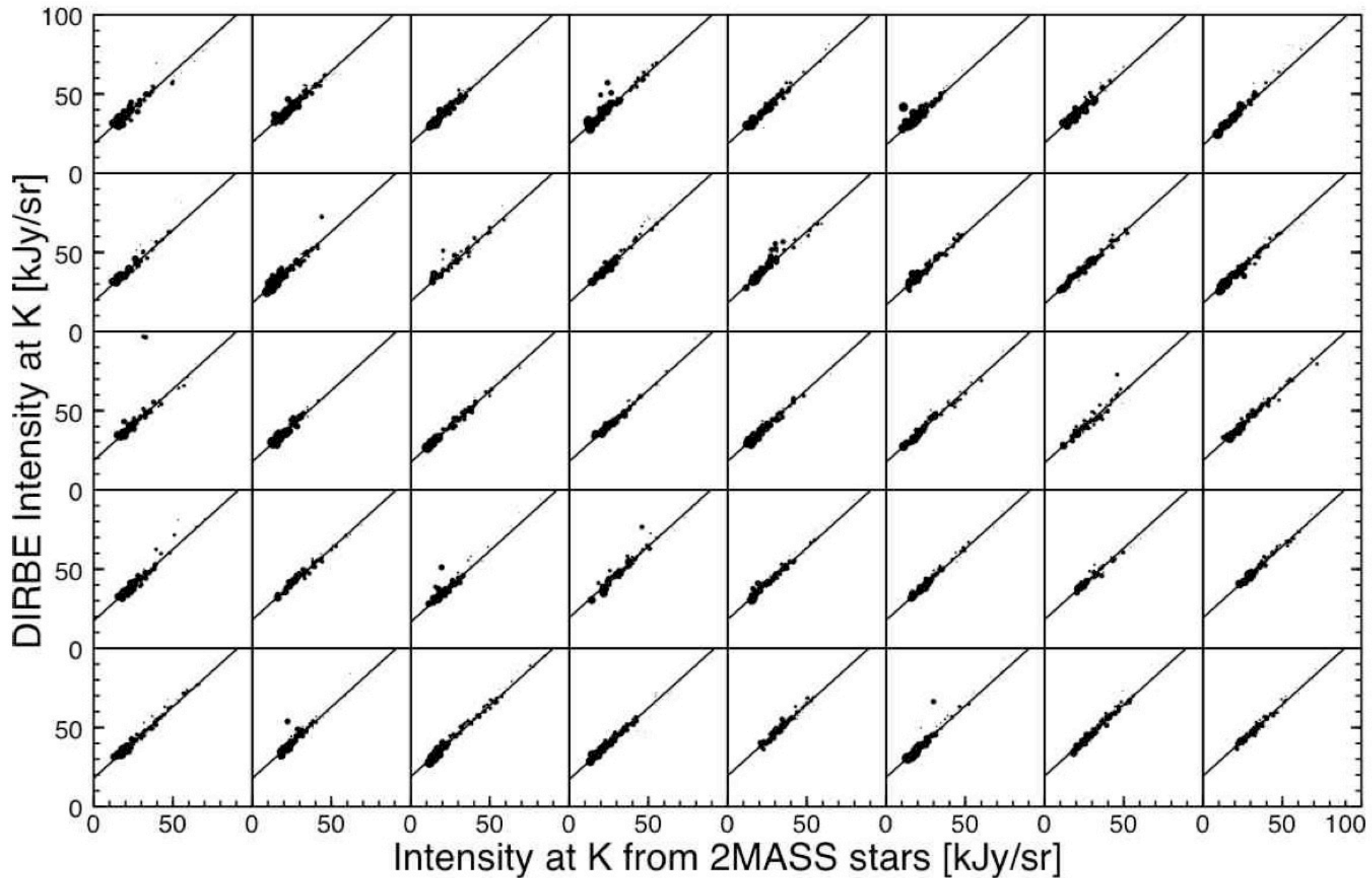
by

Louis Robert Levenson

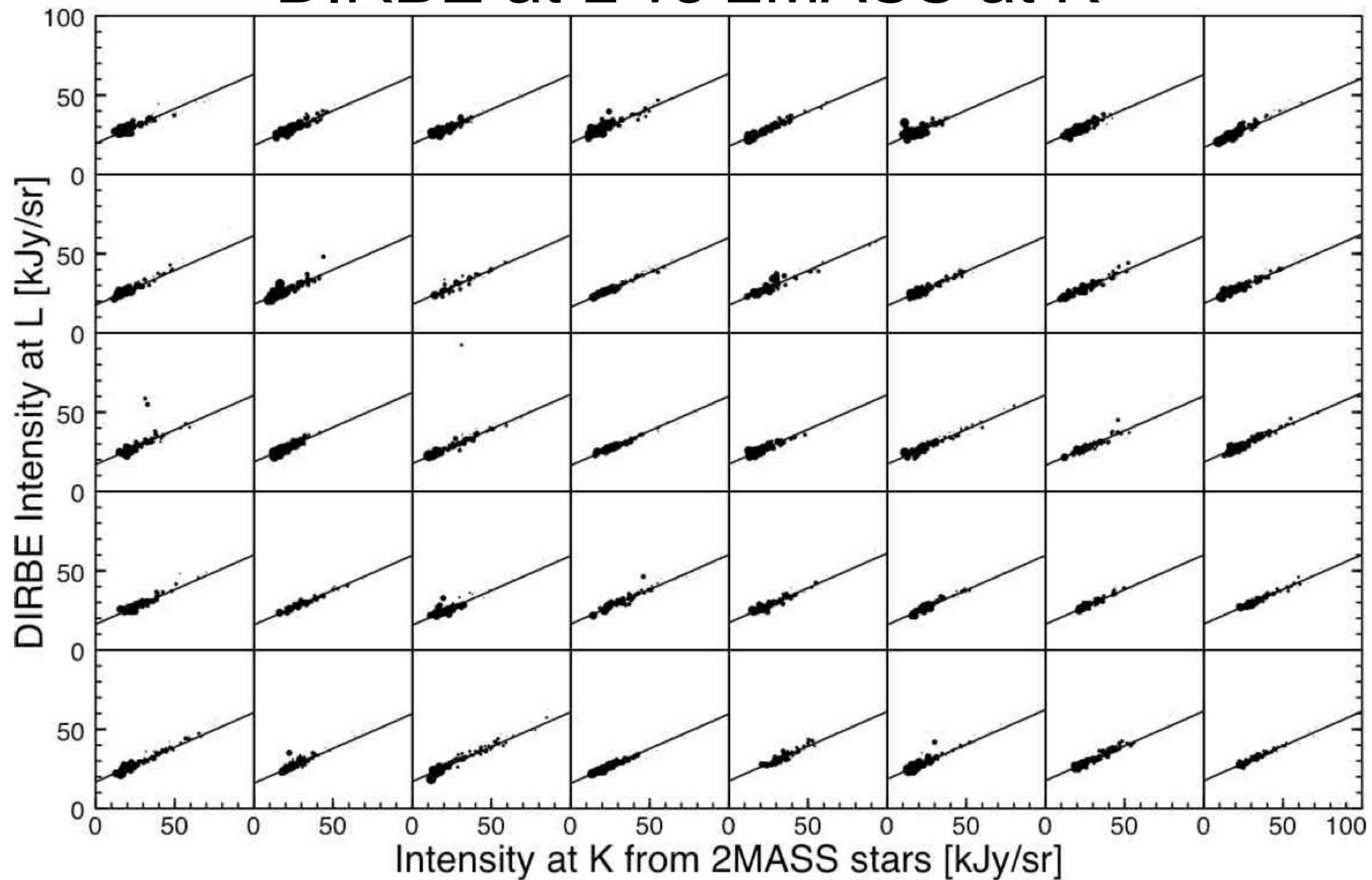




# DIRBE vs. 2MASS Fits at K



# DIRBE at L vs 2MASS at K

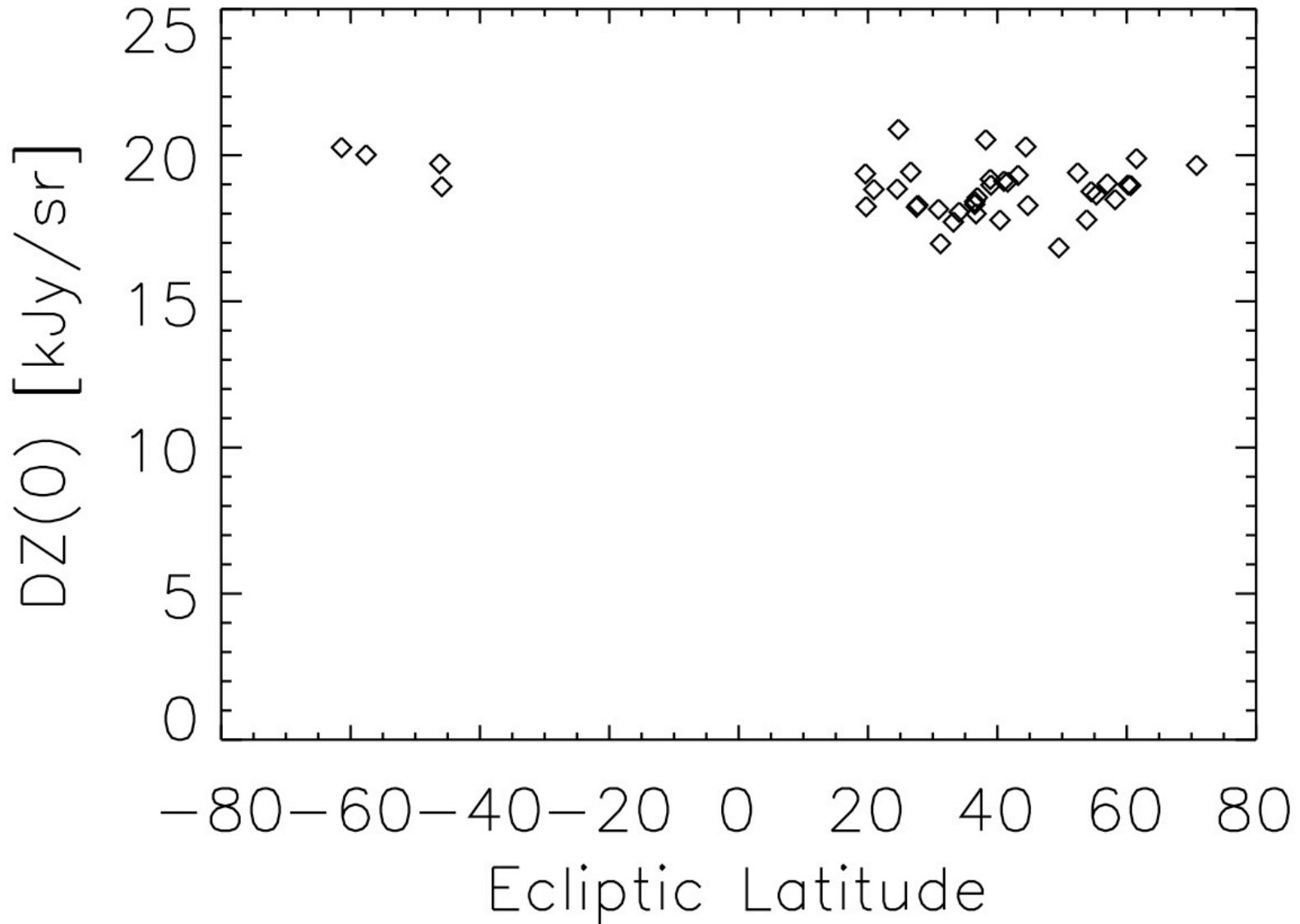


We don't need arc-  
sec resolution to  
subtract galactic stars.  
DIRBE was almost  
good enough.



# DIRBE-2MASS vs ecliptic latitude

## K-Band



# NO ZODI Principles

$$I_{obs}(\lambda, l, b, t) = Z(\lambda, l, b, t; p) + I_c(\lambda, l, b)$$

- Weak No Zodi: No restriction on  $I_c$
- Strong No Zodi:  $I_c(\lambda, l, b)$  independent of  $(l, b)$  for  $|b| > 20^\circ$  and  $\lambda = 25 \mu\text{m}$ .
- Very Strong No Zodi:  $I_c(\lambda, l, b) = 0$  for  $|b| > 20^\circ$  and  $\lambda = 25 \mu\text{m}$ .

# The Incredible Weakness of Nozoding

- For a spherical shell at distance  $R > 1$  AU from the Sun,

$$I \propto \left(1 - R^{-2} \sin^2 \theta\right)^{-0.5}$$

$$\Delta I_{rms}/\langle I \rangle \approx 0.04/R^2 < 1\% \text{ for } R > 2.$$

- For a uniform density dust cloud with a cavity of radius  $R > 1$  AU,

$$I_{bol} \propto \sin^{-1}(R^{-1} \sin \theta)/\sin \theta$$

$$\Delta I_{rms}/\langle I \rangle \approx 0.0125/R^2$$

Thus spherical components in the outer Solar System cannot be determined by applying the Weak No-Zodi Principle to DIRBE data.

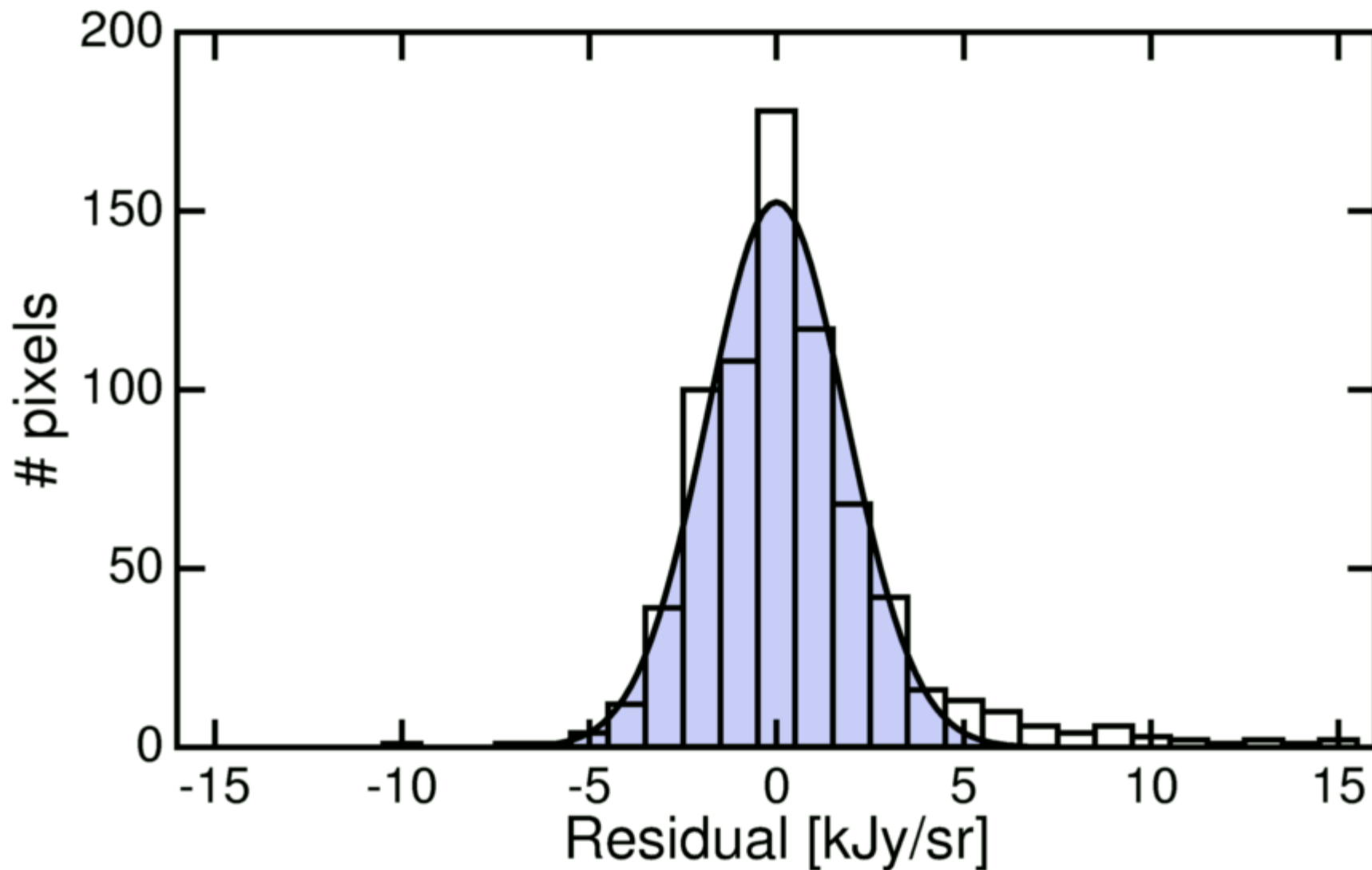
# Comparison of Zodi Models

Table 1: Intensities in the Lockman Hole  
Intensity in MJy/sr in the Lockman Hole

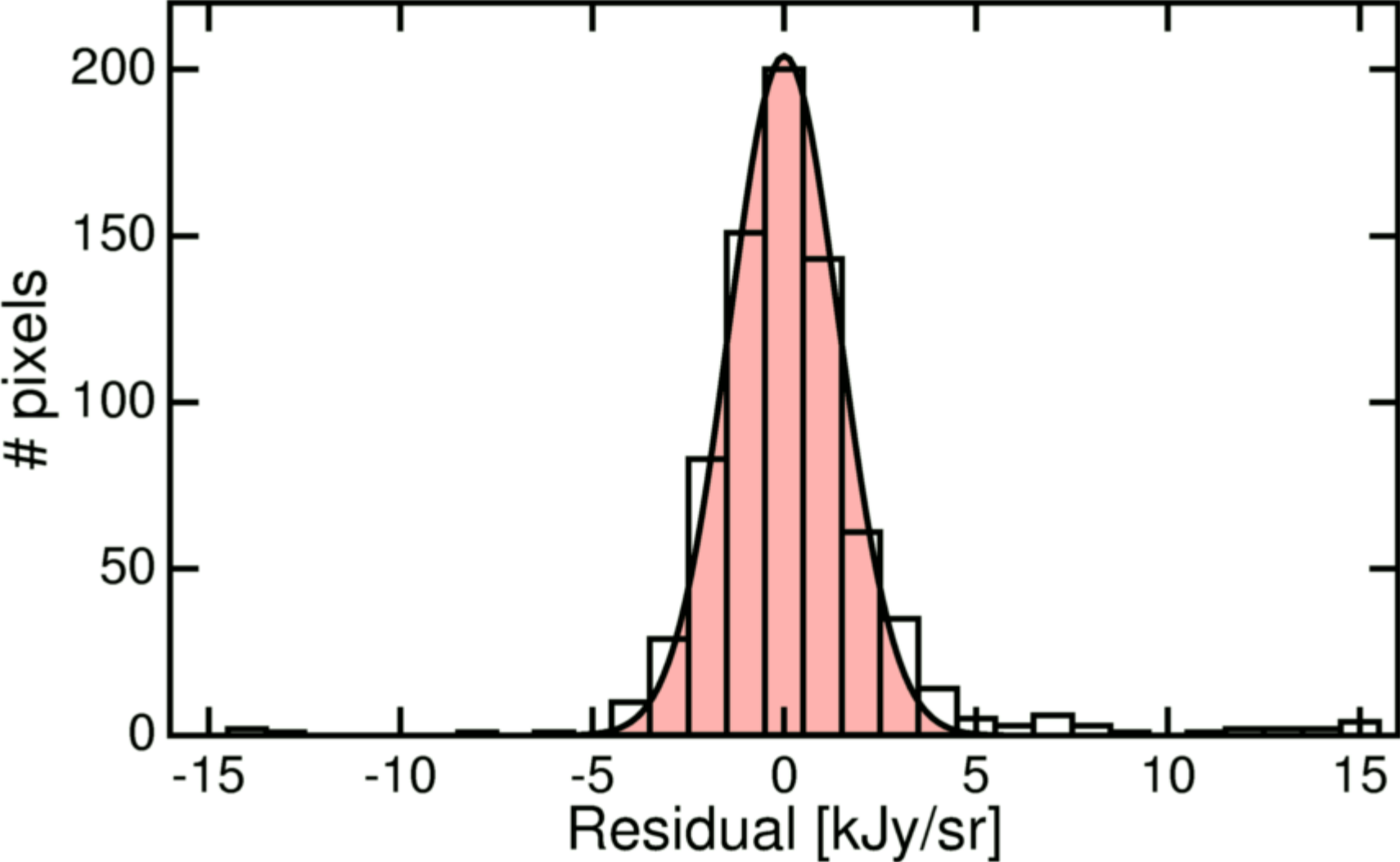
$\lambda$ [ $\mu\text{m}$ ]	GOOD1	GOOD2	GOOD3	FIZZ1	FIZZ2	FIZZ3	FIZZ3P	REALITY
1.25	0.1364	0.1407	0.1501	0.1696	0.1694	0.1806	0.1797	0.2228
2.2	0.0942	0.0971	0.1037	0.1171	0.1170	0.1247	0.1240	0.1492
3.5	0.0770	0.0788	0.0856	0.0897	0.0913	0.0976	0.0974	0.1149
5	0.4287	0.4491	0.4699	0.4993	0.4954	0.5097	0.5088	0.5389
12	13.4690	14.1695	14.7320	16.5817	16.2459	16.8580	16.8374	16.5239
25	24.6018	27.2783	30.4484	29.2170	28.4249	30.2288	30.2234	30.0306
60	6.8324	7.2153	7.4800	8.3259	8.0314	8.6145	8.6166	8.7382
100	2.4155	2.5468	2.6346	3.0647	2.9380	3.1912	3.1932	4.1884
140	1.1585	1.2219	1.3091	1.4817	1.4246	1.5838	1.5848	2.4480
240	0.3571	0.3769	0.4028	0.4622	0.4414	0.4923	0.4927	0.9459

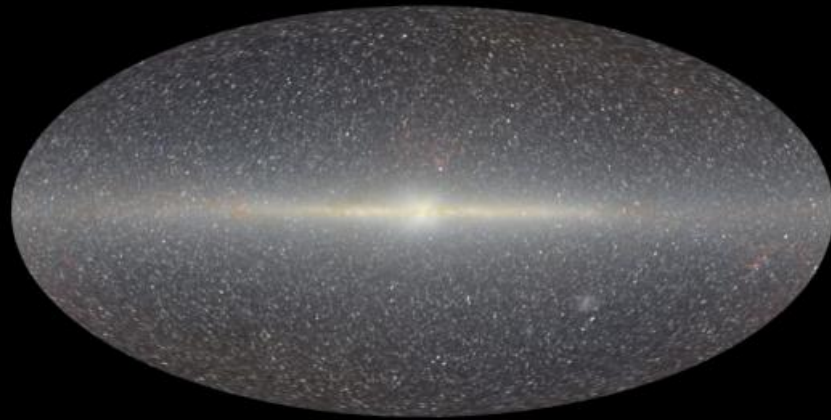
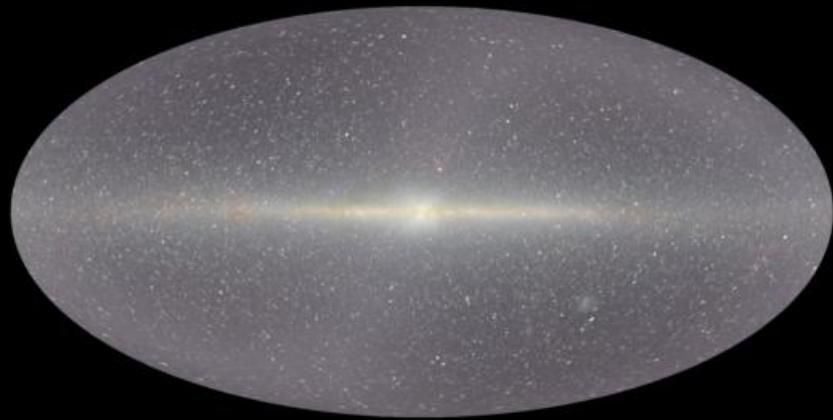
Biggest uncertainty is from the zodi models: change of 8 out of 90 kJy/sr at 3.5  $\mu\text{m}$  for different no-zodi principles.

# DIRBE-2MASS Residuals at K: $\sigma = 1.83$ kJy/sr

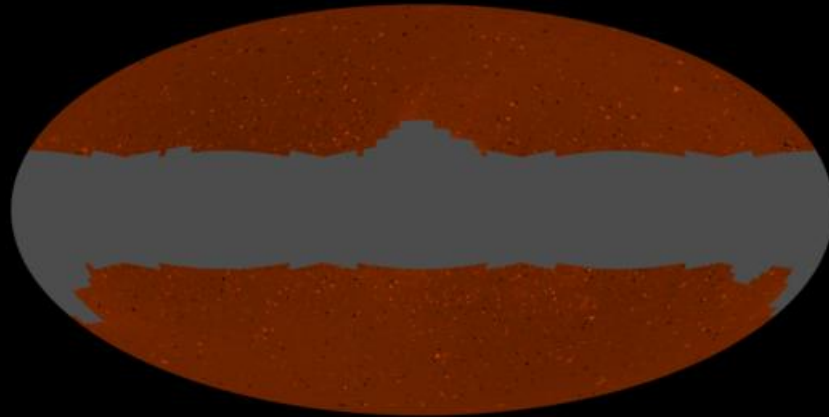


DIRBE-2MASS Residuals at L:  $\sigma = 1.43$  kJy/sr



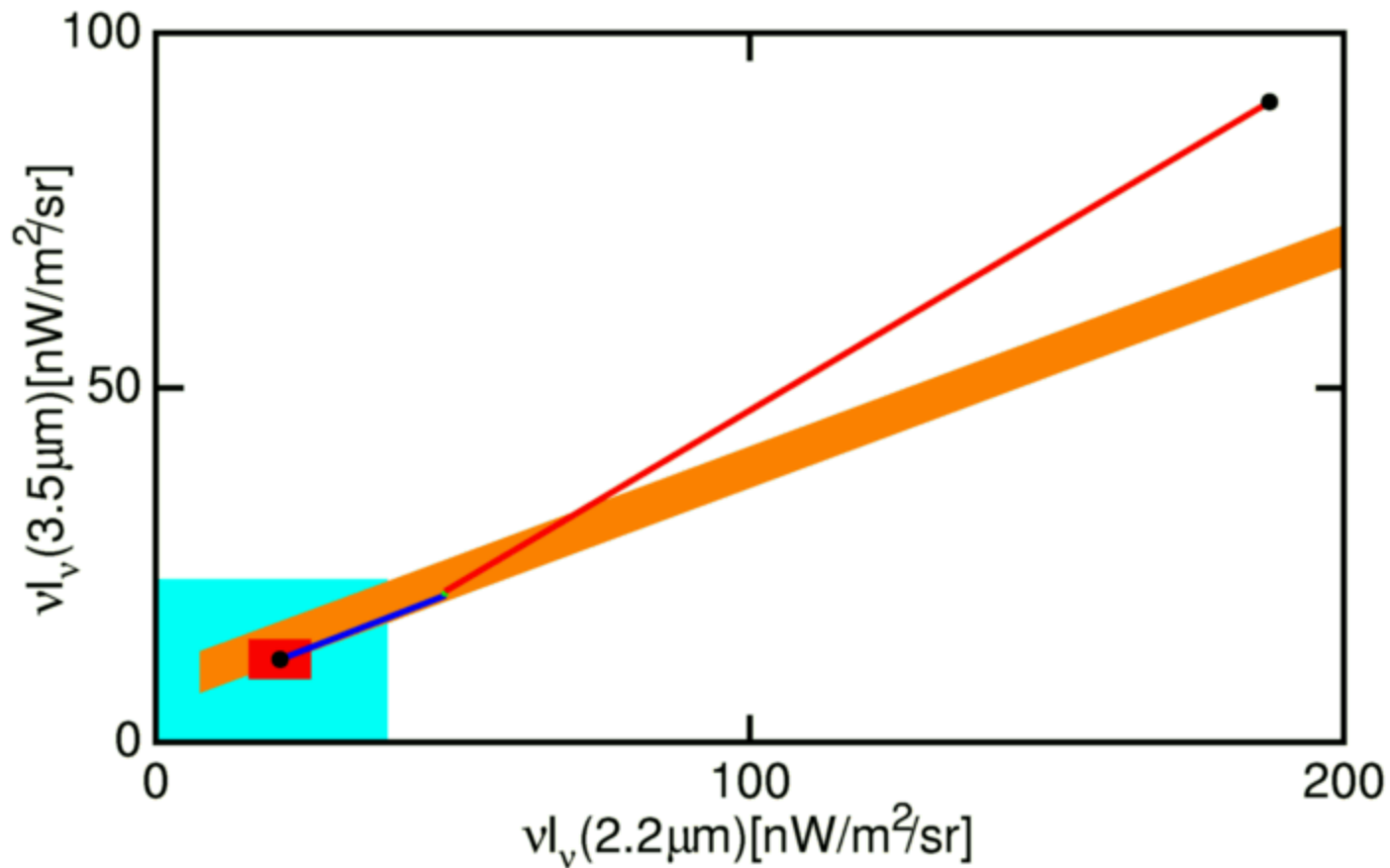


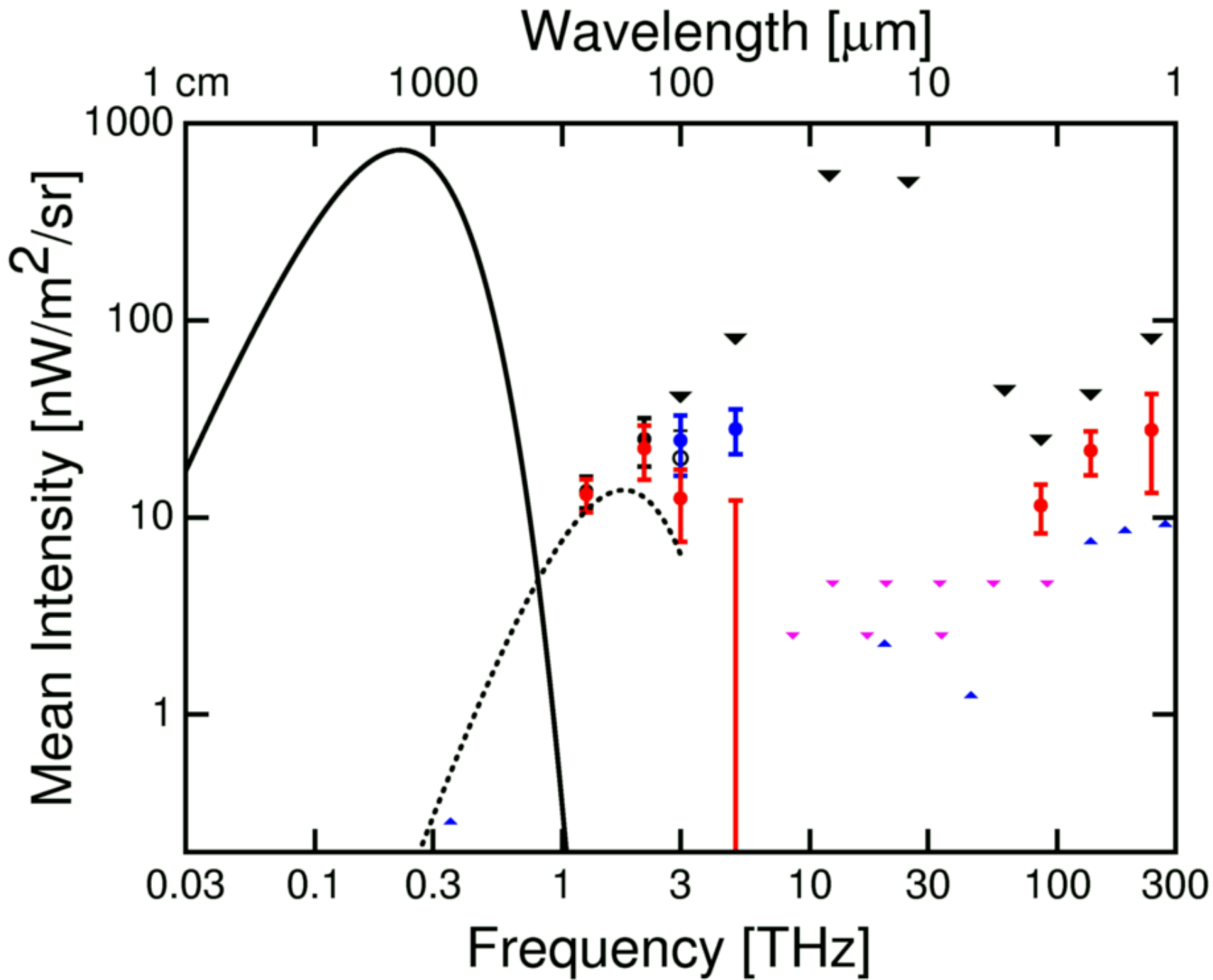
Dwek & Arendt  
(1998): L-0.5K (in  
Jy) is very smooth





# Near IR Decomposition





## Discrepancy between Counts & Measurements

- Rebecca Bernstein gets about  $2\times$  more optical extragalactic background light than one derives from the sum of the galaxy counts:

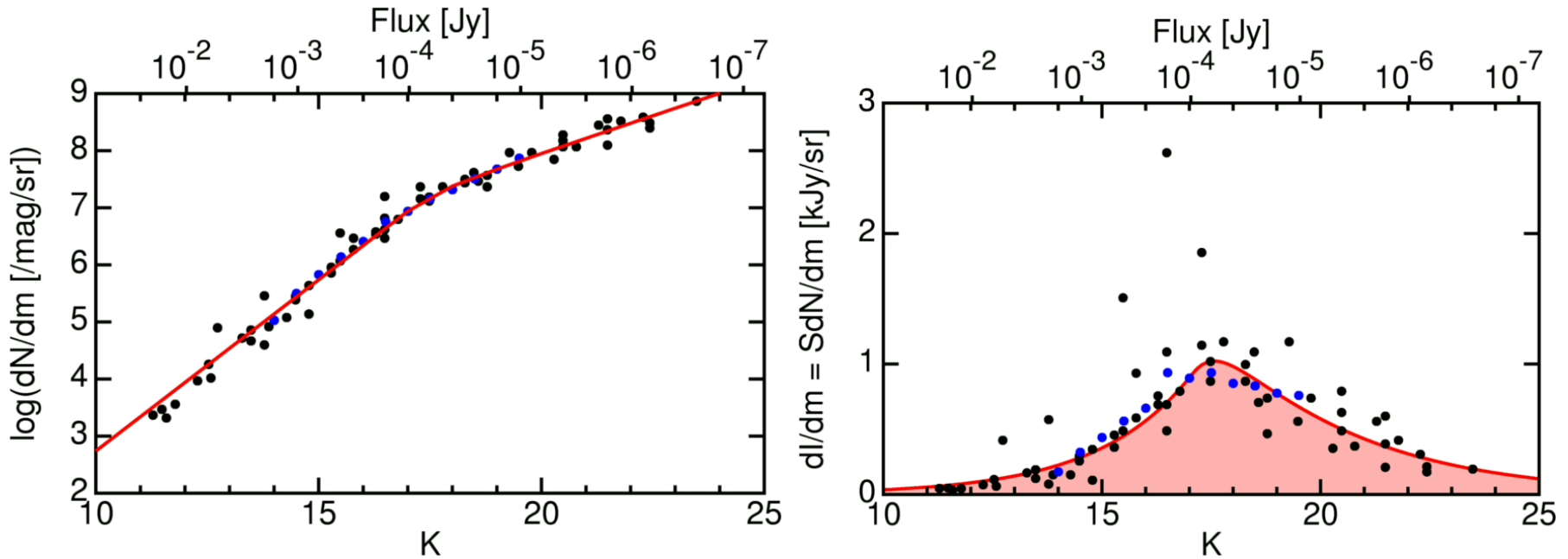
$$I_{counts} = \int S dN = \int S \left| \frac{\partial N}{\partial \ln S} \right| d \ln S$$

$$I_{obs} \approx 2 \times I_{counts}$$

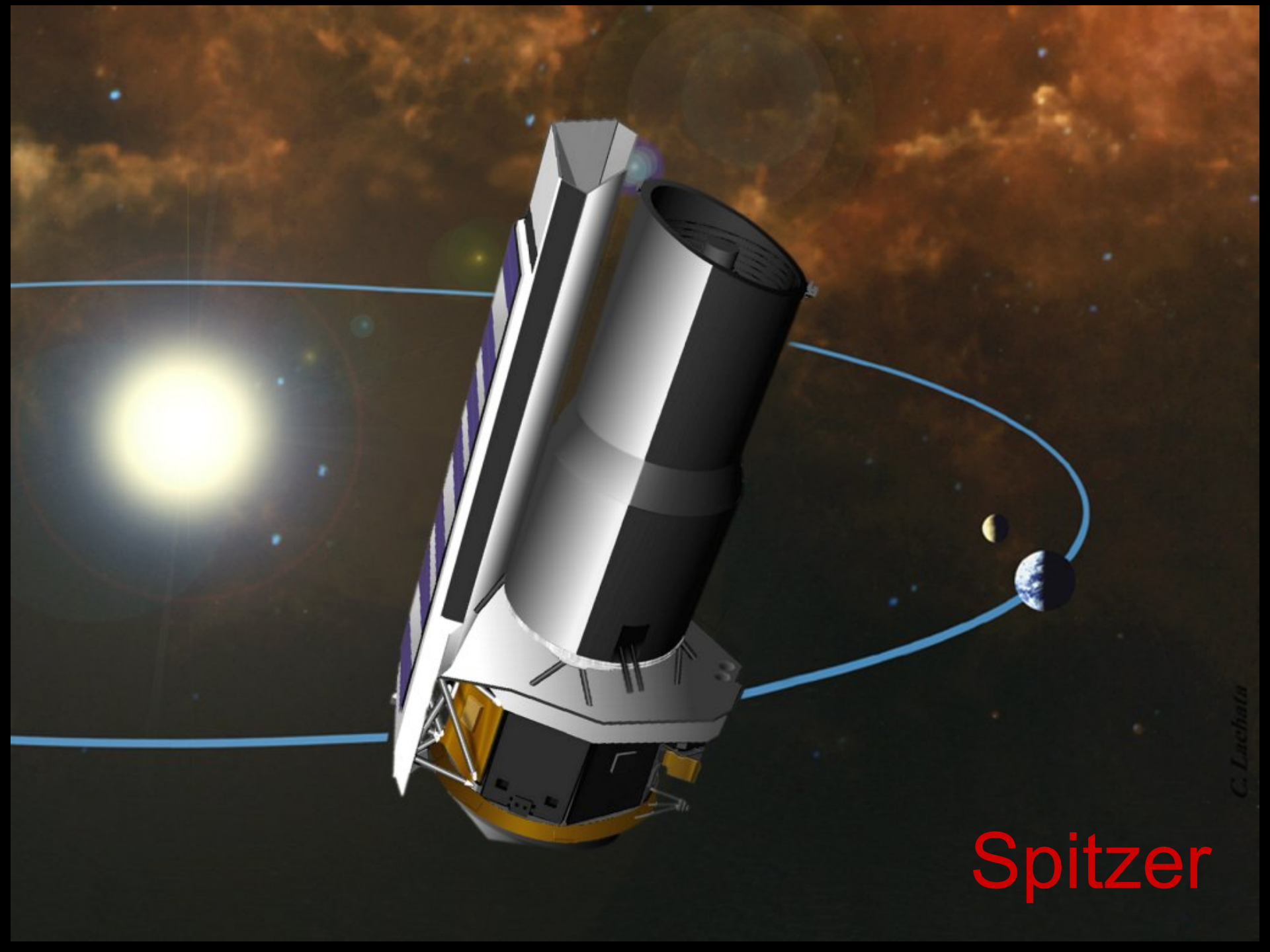
- Wright gets about  $(2 \pm 0.5)\times$  more near infrared cosmic background light than is expected from deep  $K$ -band counts.

Both measurement involve difficult and uncertain zodiacal light corrections but they use entirely different techniques.

# 2.2 $\mu\text{m}$ Galaxy Counts

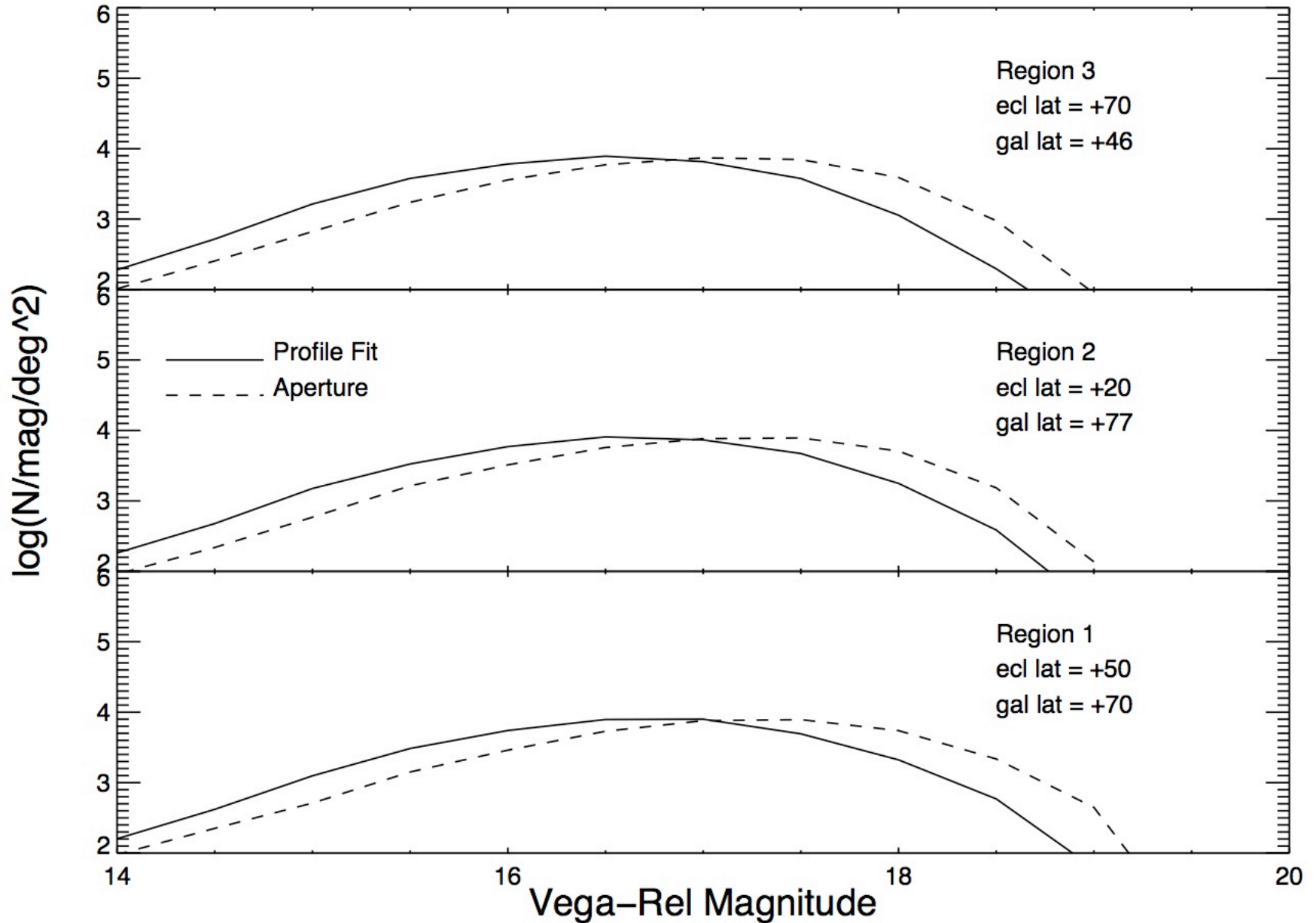


- K counts from Figure 1 of Madau & Pozzetti, MNRAS, 312, L9-L15 (2000)
- CADIS counts from Huang et al astro-ph/0101269
- Integral under fit gives 6.3 kJy/sr or 8.6 nW/m<sup>2</sup>/sr



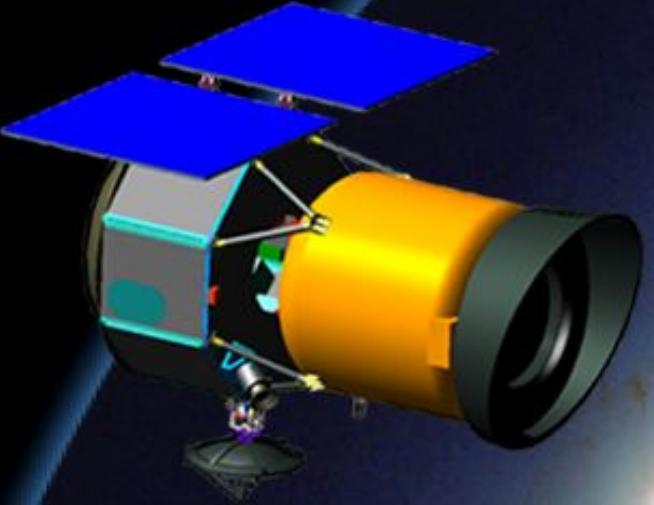
Spitzer

# Aperture fluxes systematically low





# WIDE-FIELD INFRARED SURVEY EXPLORER



I am the PI on WISE, an all-sky survey in 4 bands from 3.4 to 22  $\mu\text{m}$ . WISE found the closest stars to the Sun, the most luminous galaxies in the Universe, and also measured 100's of millions of galaxies that contribute to the CIRB.

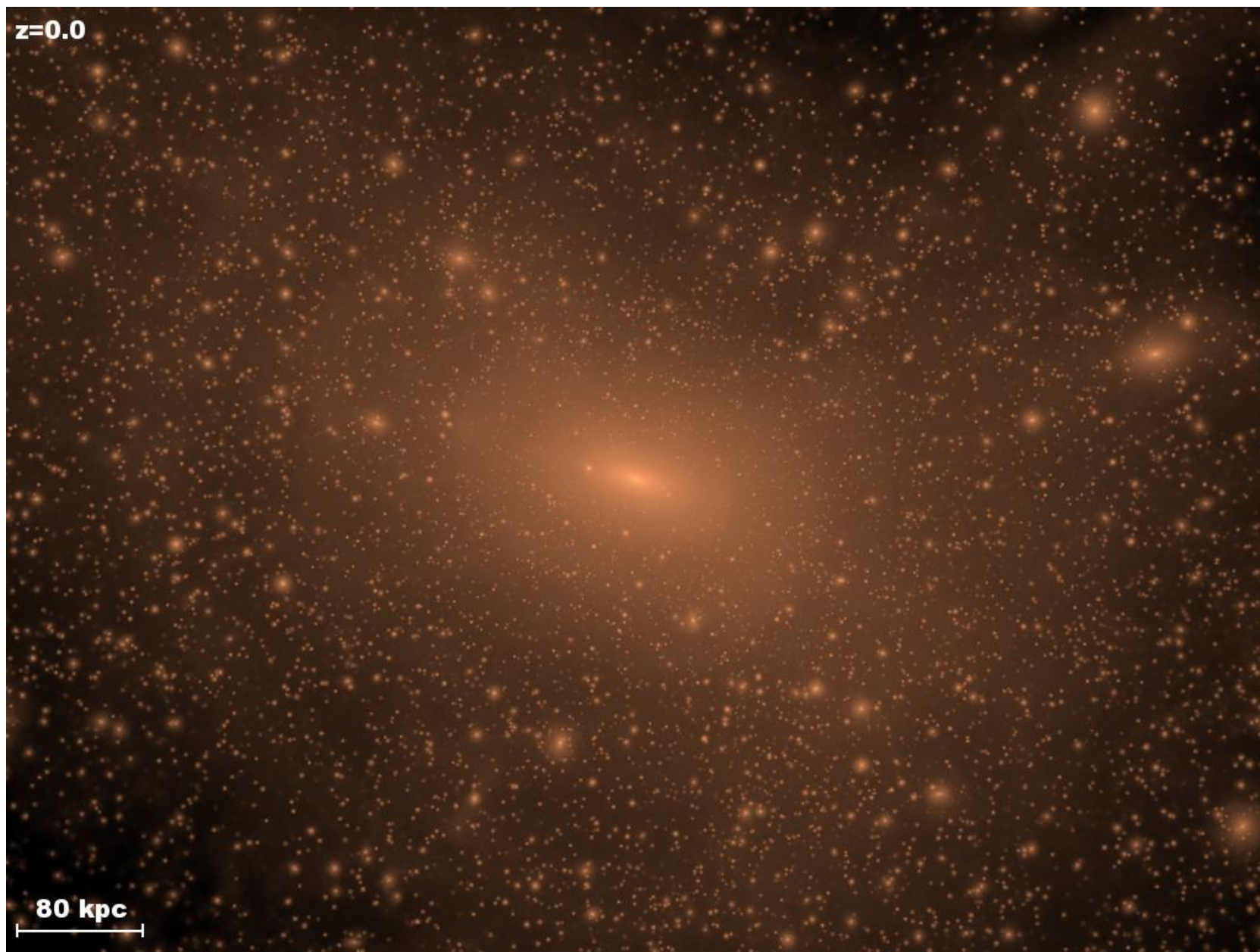
WISE launched 14 Dec 2009, and released an all-sky survey on 14 Mar 2012.







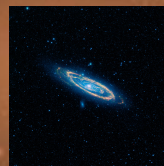
# M31-sized halo from Via Lactea



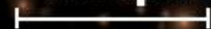


**z=0.0**

# WISE image of M31 to scale

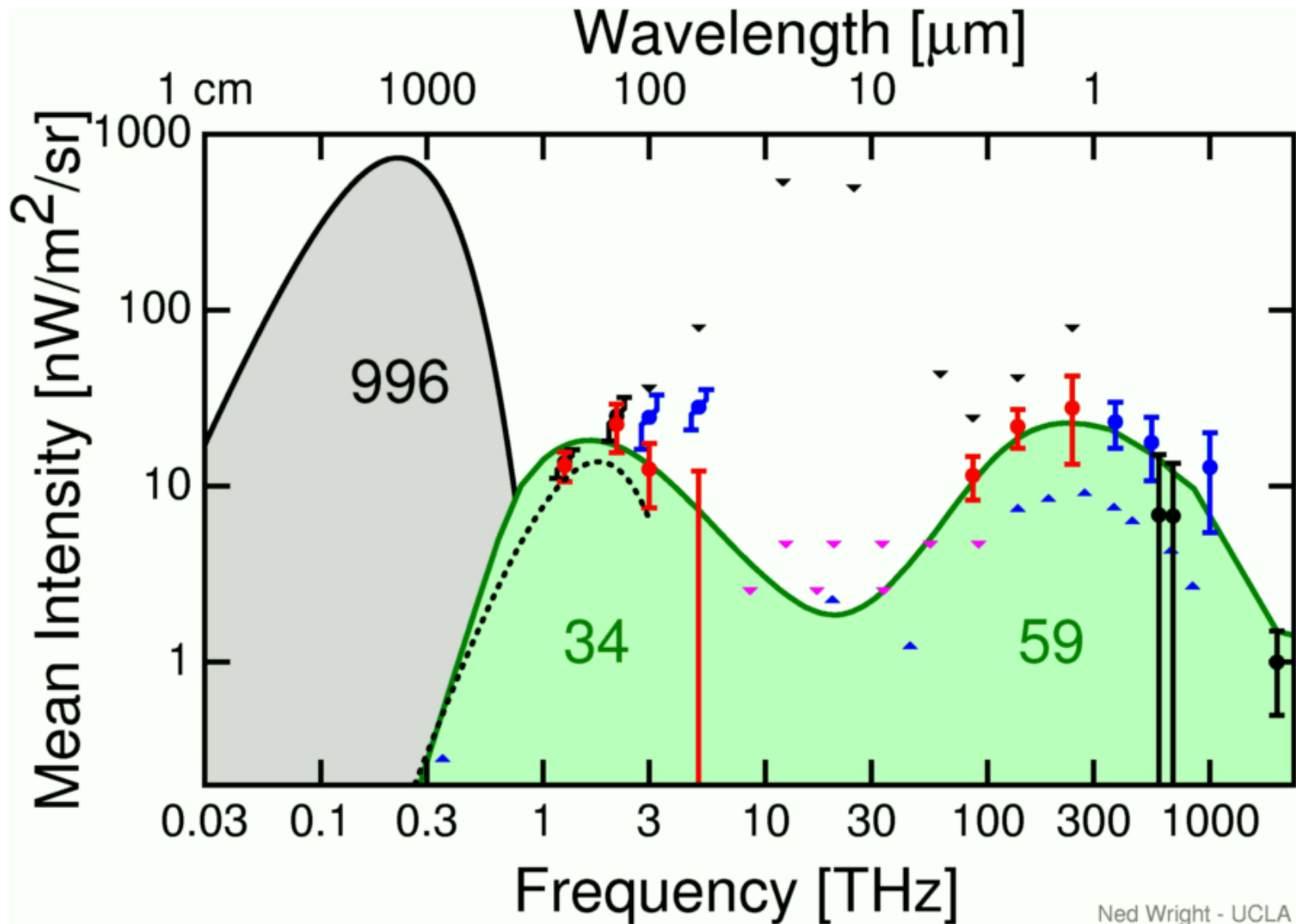


**80 kpc**



There is no blank sky  
between galaxies

# Cosmic Optical & IR Background



# $\gamma$ -Ray Connection

- The reaction  $\gamma_1 + \gamma_2 \rightarrow e^+ + e^-$  has a threshold of  $E_1 E_2 > (m_e c^2)^2$ .
- The peak cross-section of  $1.7 \times 10^{-25} \text{ cm}^2$  occurs at twice the threshold energy.
- 1 MJy/sr corresponds to a photon density of  $0.63 \text{ cm}^{-3} \text{ oct}^{-1}$ .
- Expect absorption of  $\gamma$ -rays over distance of 3 Mpc for 1 MJy/sr, or  $450/\lambda[\mu\text{m}]$  Mpc for  $20 \text{ nW/m}^2/\text{sr}$ , at  $E \approx 400\lambda[\mu\text{m}] \text{ GeV}$ .



# VERITAS



- Operating at Mt Hopkins base camp
- Four 12 meter diameter segmented mirror telescopes
- Competitive with HESS in Namibia



# Ablaze in the distance

- Blazar H1426+428 has a flux of  $4 \times 10^{-12}$  erg/cm<sup>2</sup>/sec at  $E > 1$  TeV.
- Redshift  $z = 0.129$
- 96%  $\gamma$ -ray absorption at 1 TeV for my CIRB
- Many other blazars now known
- But see Kusenko Tuesday morning about the “proton tunnel”

# Conclusions

- The CIRB has been detected in both the far IR and the near IR windows through the interplanetary dust, but measurements between 5-60  $\mu\text{m}$  are impossible from 1 AU
- Bolometric OIR background is about 100 nW/m<sup>2</sup>/sr
- Ratio of optical plus near IR to the far IR is about 2:1
- Biggest uncertainty is the zodiacal light, & DESIRE, LZM or ZEBRA would help.

# Discussion: $\Delta X$

- For UVO “Madau” curve, fuel burn over current energy density ratio is  $f/U_0 = 1.9$
- Current CIRB bolometric energy density is about  $100 \text{ nW/m}^2/\text{sr}$
- Therefore  $\Delta X = -0.033$
- Madau curve with Rowan-Robinson add-on at high  $z$  burns more fuel at high redshift, so  $f/U_0 = 2.3$ ,  $\Delta X = -0.04$
- At  $1/3$  solar from cluster gas,  $\Delta X = -0.02$
- Do we need more baryons [CMB], more AGN, or less CIRB [zodi]?

# White Dwarf Helium Reservoir?

- Oppenheimer *et al.* (2001) claim 3% of local halo in old WDs

$$\Delta X = -0.04h \frac{f_{WD}}{0.03} \frac{M_H}{5 \times 10^{11}} \frac{\ell_{ONIR}}{5.6h \times 10^8} \frac{3 \times 10^{10}}{L_{MW}}$$

- BUT Richer on 2 Apr 2001 withdrew the claimed detection of faint, high proper motion stars in the HDF [astro-ph/9908270]
- The Oppenheimer *et al.* objects do not have a halo velocity distribution, and can not be part of a spherical halo.

## REFERENCES

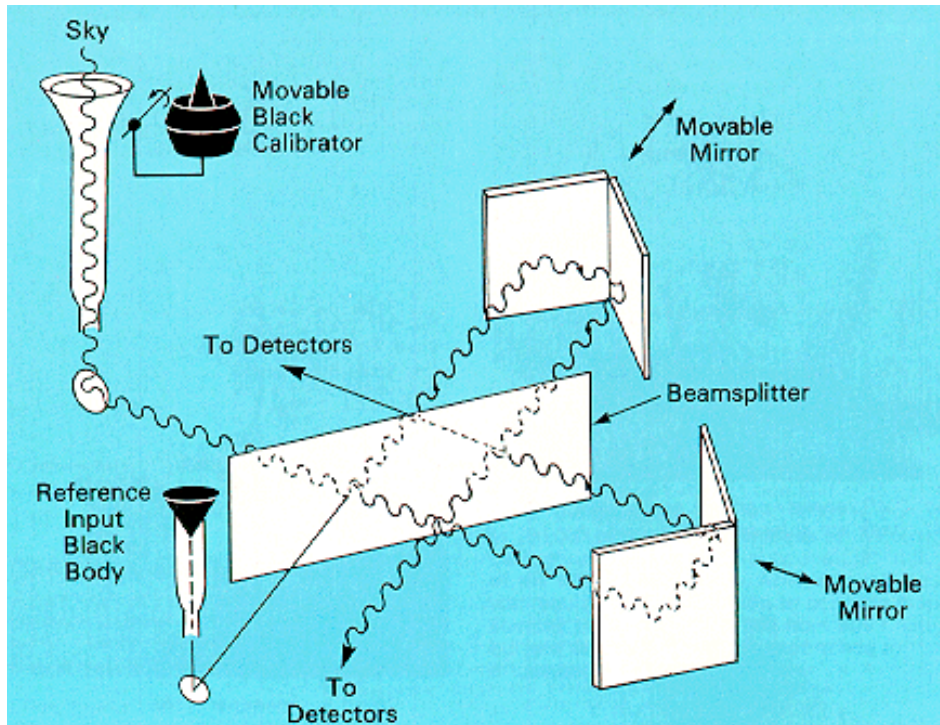
### COSMOLOGY TUTORIAL:

<http://www.astro.ucla.edu/~wright/cosmolog.htm>

<http://www.astro.ucla.edu/~wright/CIBR>

- DIRBE: <http://www.astro.ucla.edu/~wright/DIRBE/>
- SIRTf: <http://sirtf.caltech.edu/>
- <http://xxx.lanl.gov/abs/astro-ph/9909428> (Dark Spot)
- <http://xxx.lanl.gov/abs/astro-ph/9912523> (Histogram)
- <http://xxx.lanl.gov/abs/astro-ph/0004192> (2MASS)

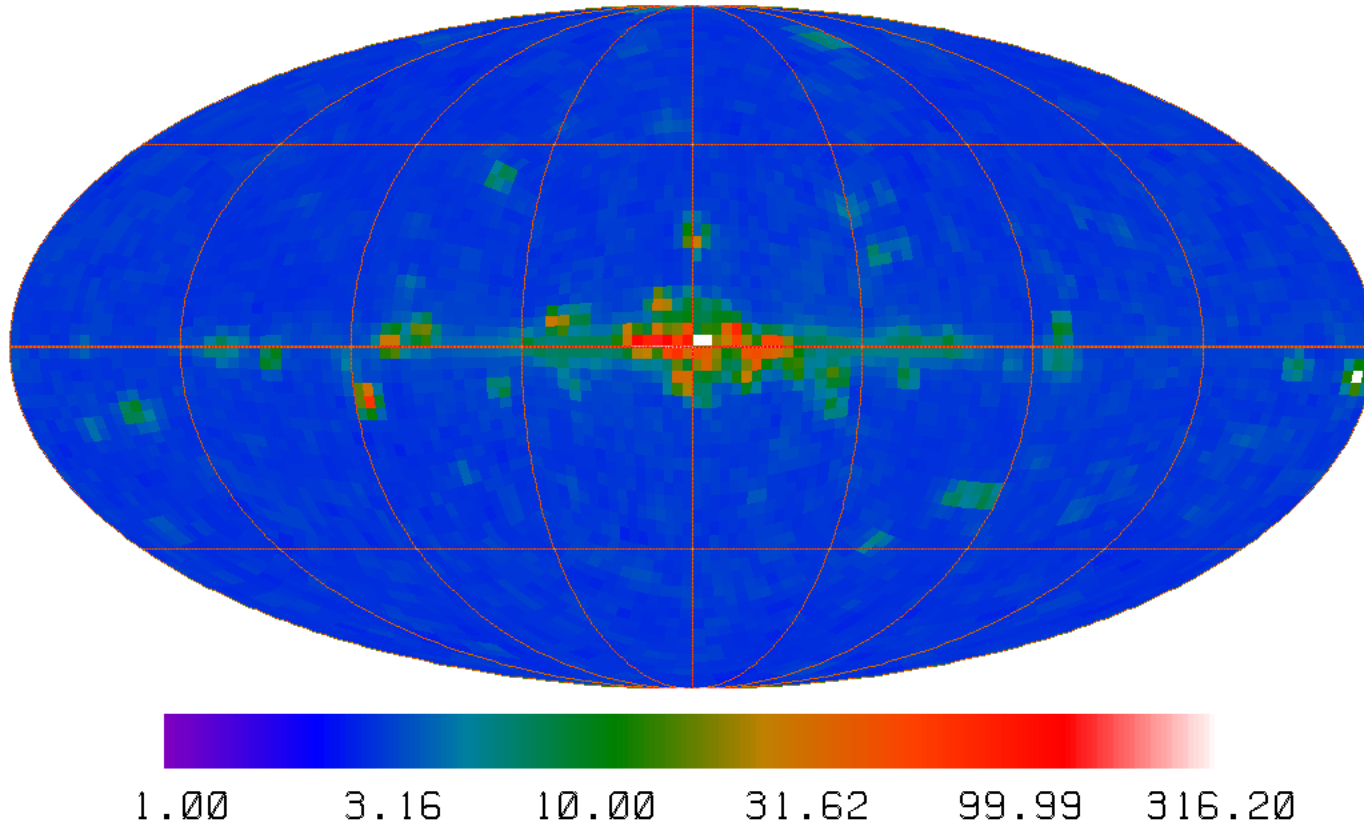
# FIRAS measured the CMB spectrum



- Far InfraRed Absolute Spectrophotometer
- A differential polarizing Michelson interferometer
- Zero output when the two inputs are equal
- One input is either the sky or a very good blackbody, other is a pretty good blackbody

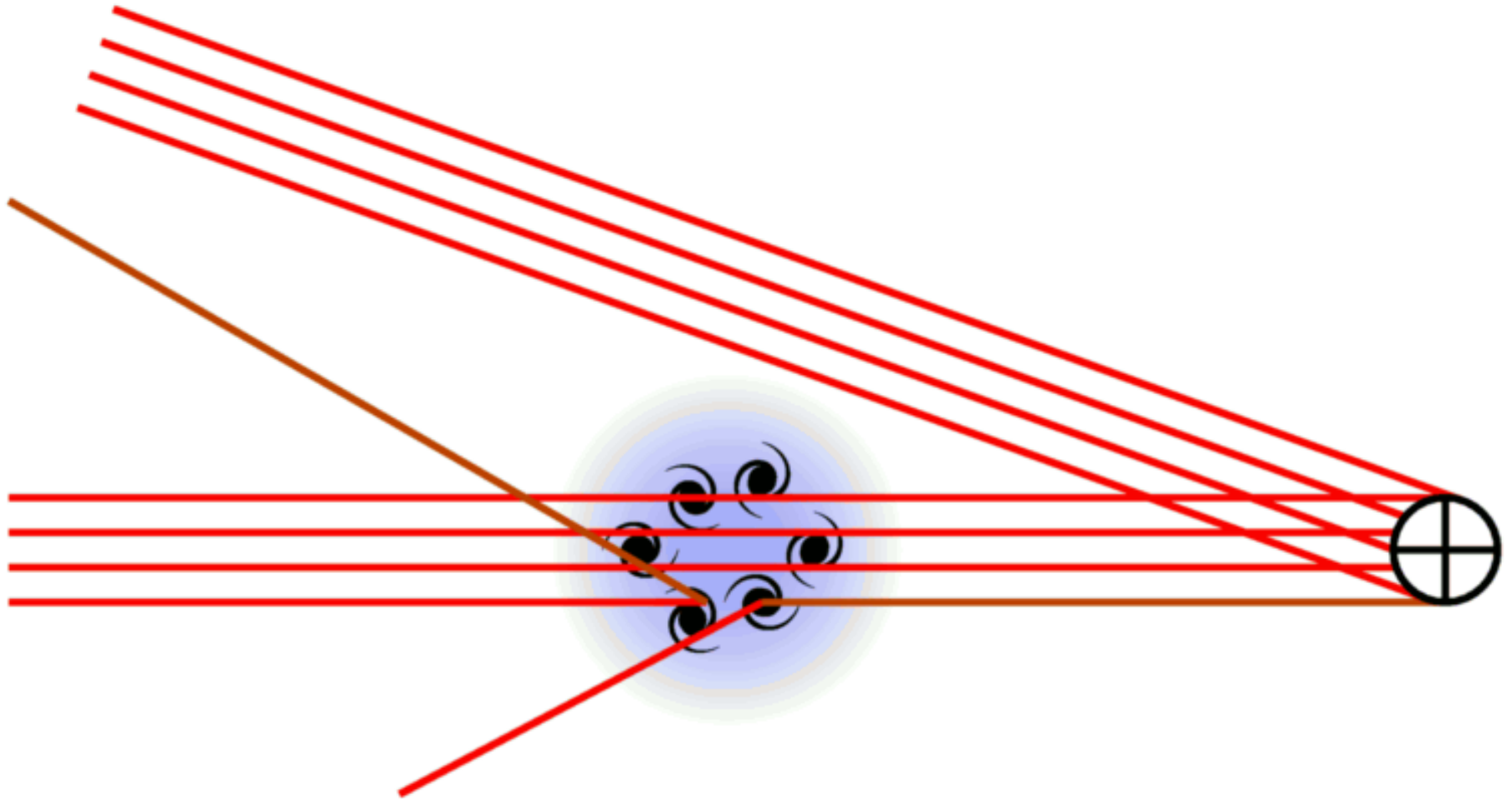
# X-Ray Background

HEAO A2 X-RAY FLUX



Hot diffuse plasma gives right spectrum but violates FIRAS limit on  $y$ . Source thought to be highly obscured AGNs.

# Sunyaev-Zeldovich Effect

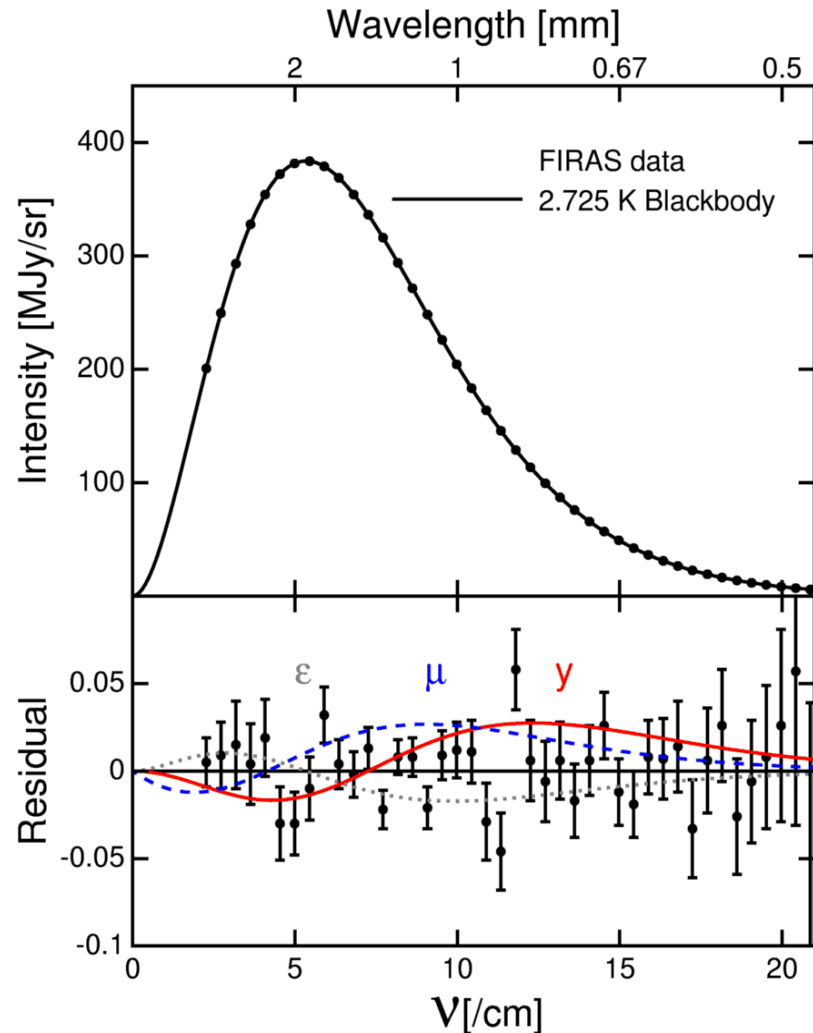


Hot electrons increase high  $\nu'$  s, decrease low  $\nu'$  s



# FIRAS Final Spectrum

- SZ Effect  $\propto$   
 $y = N_e \sigma_T k T_e / m_e c^2$   
 $< 15 \times 10^{-6}$
- Bose-Einstein  
 $\mu < 9 \times 10^{-5}$
- Energy from hot electrons into CMB  $< 60$  parts per million



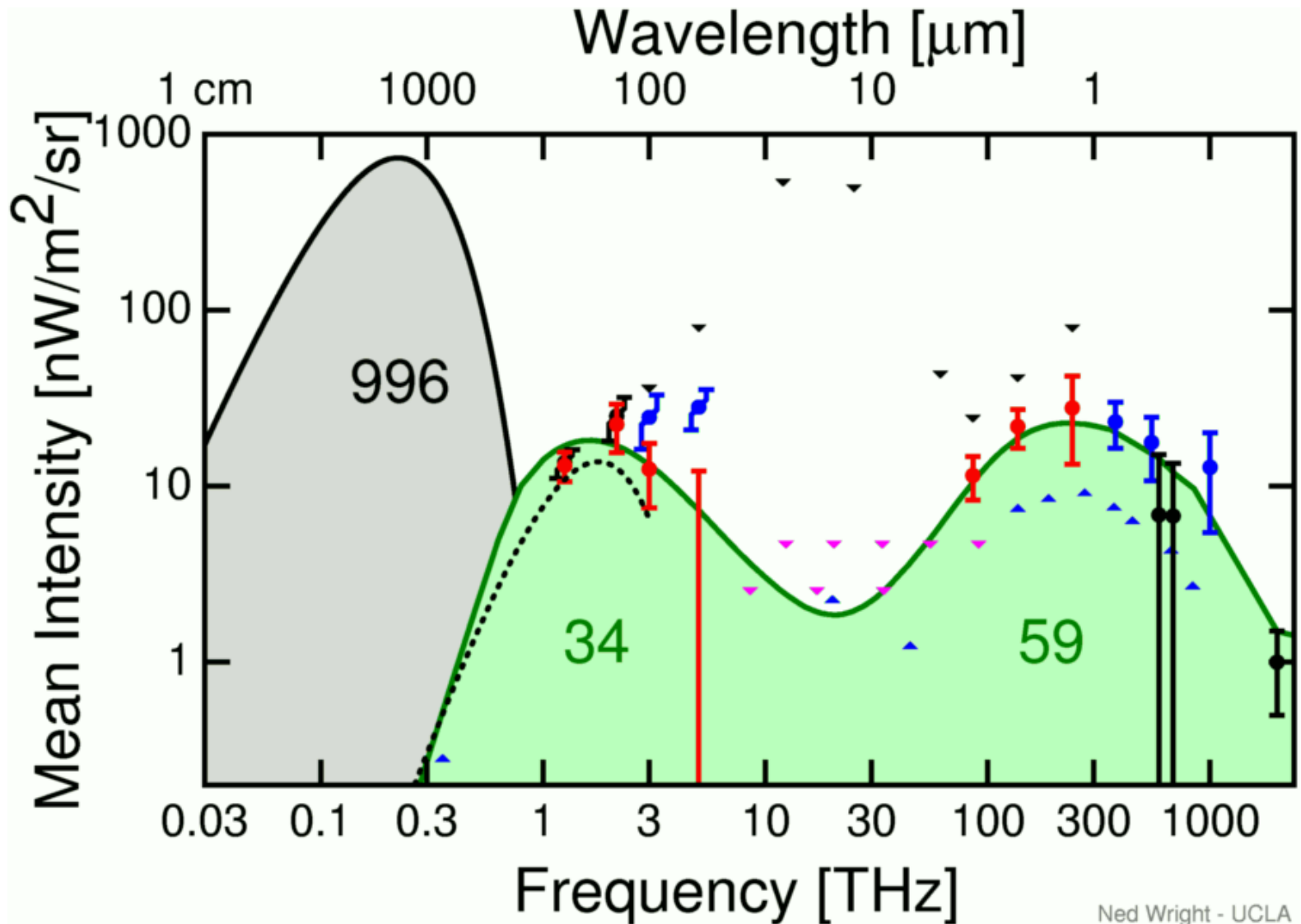
# FIRAS Far IR Background

- Any proposed FIR background must be compatible with the limits on the CMB distortion.
- This means the FIRB is either small or similar to the galactic spectrum.
- The FIRAS fit to the  $N_{\text{H}} = 0$  intercept of the high frequency channel, [astro-ph/9803021]  
$$I_{\nu} = 1.3 \times 10^{-5} (\nu/100)^{0.64} B_{\nu}(18.5 \text{ K}),$$
fails with  $\Delta\chi^2 = 22.6$  (4.75  $\sigma$ )

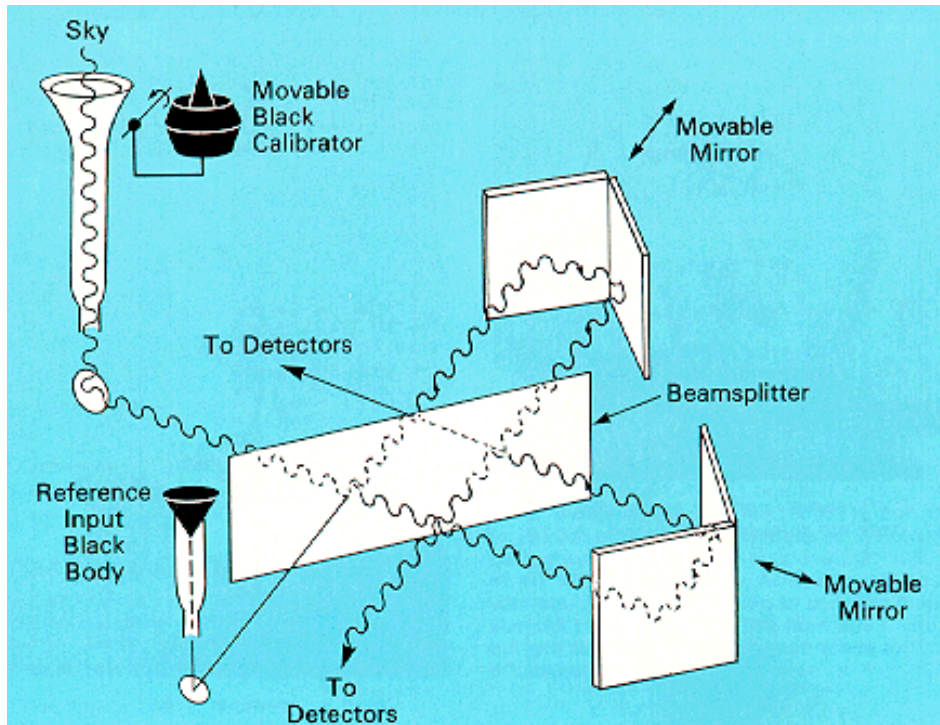
# No simple limit at 850 $\mu\text{m}$

- FIRAS fit fails with  $\Delta\chi^2 = 22.6$  but  
 $I_{850} = 143 \text{ kJy/sr}$
- Lagache *et al.* fit [astro-ph/9901059] is marginal with  $\Delta\chi^2 = 5.8$  but  
 $I_{850} = 115 \text{ kJy/sr}$
- The modified scaled Primack model on the next slide is fine with  $\Delta\chi^2 = 2.2$  but  
 $I_{850} = 195 \text{ kJy/sr}$

# Cosmic Optical & IR Background



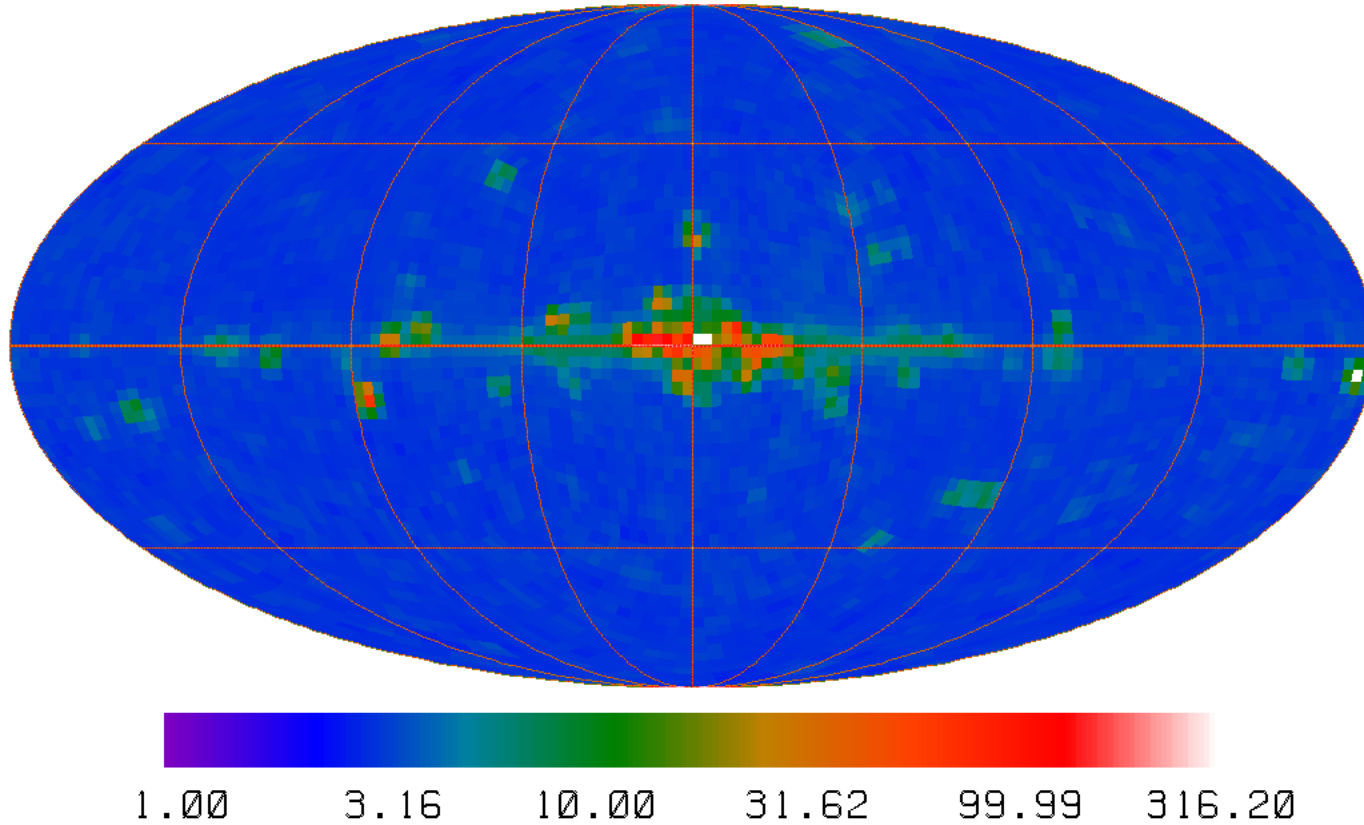
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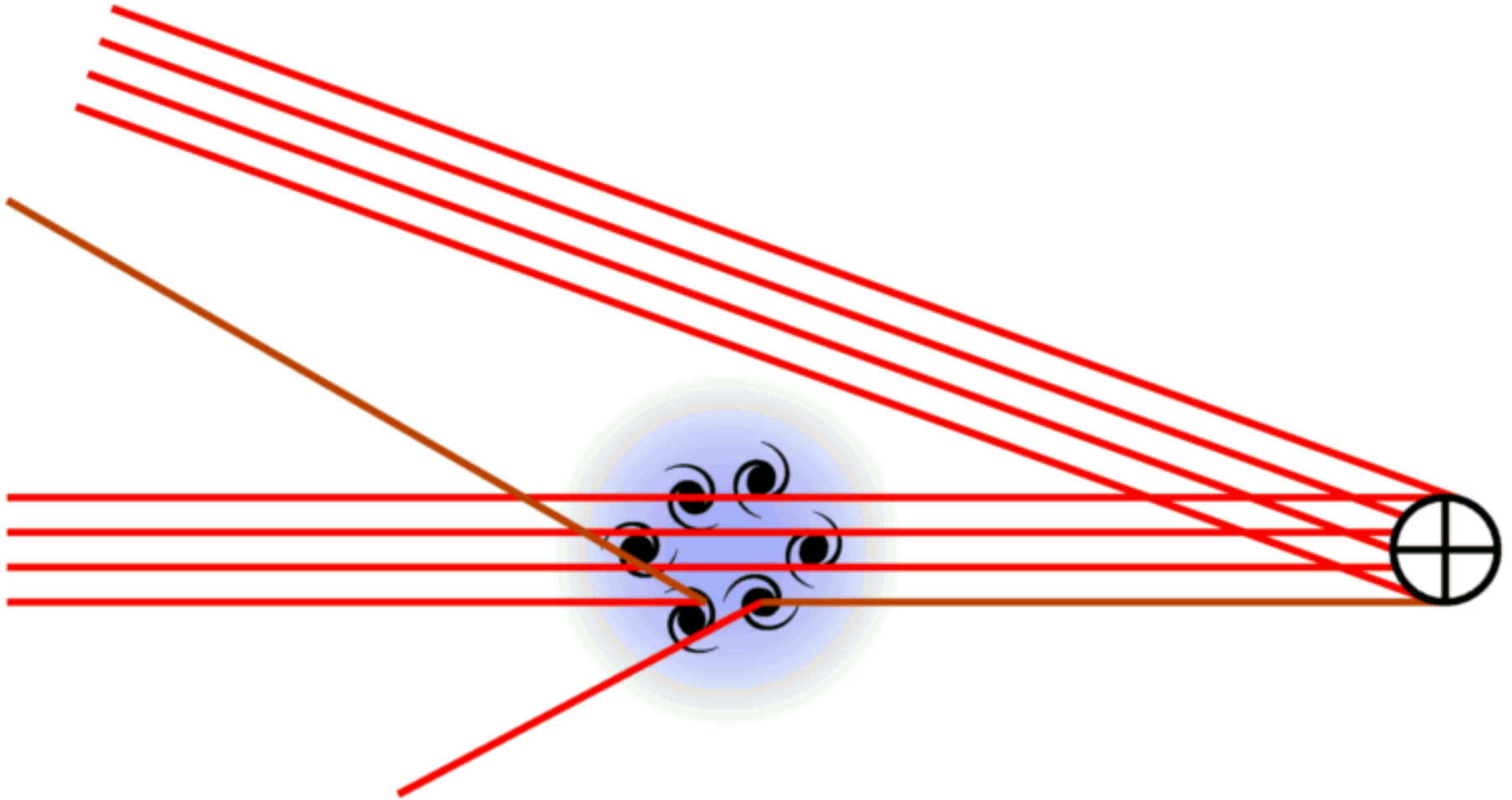
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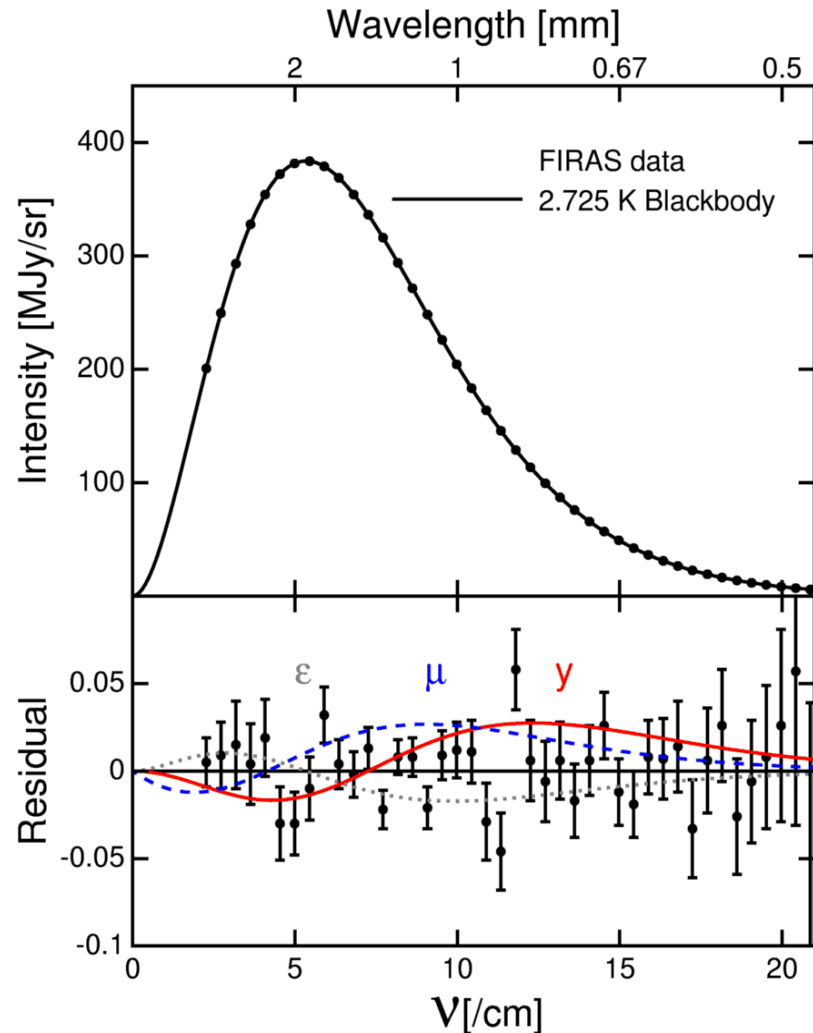


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