Extragalactic and Galactic Gamma-Rays and Neutrinos from Dark Matter Annihilation

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Discovery Time...

We are about to enter into an era of major discovery

Dark Matter: we need new particles to explain the content of the universe

Standard Model: we need new physics

Supersymmetry solves both problems!

The super-partners are distributed around 100 GeV to a few TeV

LHC: directly probes TeV scale

Fermi, IceCube are probing this scale indirectly through DM annihilation into photons, neutrinos [This talk]

Direct detection [XENON, CDMS, Cogent etc.] are also probing the new physics scale

So Far at the LHC

- Recent Higgs search results from Atlas and CMS indicate excess of events beyond background which is consistent with a Higgs mass of around 125 GeV
 - in the tight MSSM window: 115-135 GeV
- → squark mass (first generation) ~ gluino mass ≥ 1TeV
- →For heavy squark mass, gluino mass is ≥ 700 GeV
- → stop (squark) produced from gluinos, stop mass ≥ 400 GeV
- → stop (squark) produced directly, stop mass ≥ 180 GeV

Models for This Talk

1. mSUGRA/CMSSM : neutralino dark matter

- 4 parameters + sign: m_0 , $m_{1/2}$, A_0 , tan β and Sign(μ)
- 2. $SU(3)_{c}xSU(2)_{L}xU(1)_{Y}xU(1)_{B-L}$

Motivations for B-L models: B-L models are used (for several decades) to explain Neutrino mass

Right handed neutrino and corresponding sneutrino are included in this model

 $W = W_{\text{MSSM}} + y_D \mathbf{N}^c \mathbf{H}_u \mathbf{L} + f \mathbf{H}_2' \mathbf{N}^c \mathbf{N}^c + \mu' \mathbf{H}_1' \mathbf{H}_2',$

Right sneutrinos can be dark matter candidates

mSUGRA Parameter space



 The direct searches at the LHC,
 the Br(B_s→μμ) measurement from LHC, Tevatron and direct DM detection experiments are probing the parameter space

Focus point

Dutta, Mimura, Santoso arXiv:1107.3020

mSUGRA Parameter space...



Olive, Ellis, 1202.3262

With Higgs mass...

Final States from DM annihilation

Focus point:

Larger m_0 reduces $\mu \rightarrow$ larger Higgsino component in the neutralino

Annihilation dominantly produces W⁺ W⁻ final states

Stau neutralino Coannihilation:

Annihilation dominantly produces $b\overline{b}$, $\tau^+\tau^-$ final states However, due to the absence of t at present, the annihilation cross-section is smaller than the freeze out time

Sneutrino Annihilation in the B-L model:

Sneutrino annihilation can produce right handed neutrinos
Higgs(h)

γ-ray Intensity

Calculating γ-ray intensity

Particle side:

Annihilation Cross-section: σv (v: relative velocity), Dark matter mass m_{DM}, Photon spectrum per annihilation dN $\gamma(E\gamma)$ /dE γ (Include radiative emission by charged products and decays of unstable products.)

Astro side: Density Profile (NFW profile: $\rho_h(r) = \frac{\rho_s}{\frac{r}{rs}(1+\frac{r}{r_s})^2}$), $[r_s=R_{vir}/c]$

Mass function: dn(z)/dM (number densities of M mass halos: Sheth-Tormen), velocity variance profile

γ-ray Intensity....

Extragalactic

$$I_{\gamma,\mathrm{EG}}(E_{\gamma}) = \sigma v \int \frac{dz}{H(z)} W((1+z)E_{\gamma}, z) \langle \rho^2 \rangle(z).$$

Intensity Window function: $W(E_{\gamma}, z) = \frac{1}{8\pi m^2} \frac{1}{(1+z)^3} \frac{dN_{\gamma}}{dE_{\gamma}}(E_{\gamma})e^{-\tau(E_{\gamma}, z)},$

 $\tau(E_{\gamma},z)$ is the cosmic opacity to gamma rays

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Mean square matter density:
$$\langle \rho^2 \rangle(z) = \int dM \frac{dn}{dM}(M, z) \int d^3 \mathbf{r} \rho_h^2(r|M, z).$$

$$\rho_h(r|M, z) = \frac{\rho_s(M, z)}{[r/r_s(M, z)][1 + r/r_s(M, z)]^2} \xrightarrow{d_n \int d^3 r \rho_h^2} (\frac{\rho_s^2}{M_0}) = \mathbf{r} = 1$$
Sheth-Torman
$$\int_{10^{-4} - 10^{-2}}^{10^{-4} - 10^{-2}} \frac{10^{-2} (\rho_s^2 - 10^{-2})}{10^{-6} + 10^{-2} + 10^{-2} + 10^{-2}} = 0$$

γ-ray Intensity....

Galactic

$$I_{\gamma,G}(E_{\gamma},\psi) = \frac{\sigma v}{8\pi m^2} \frac{dN_{\gamma}}{dE_{\gamma}}(E_{\gamma})\hat{J}(\psi)$$

J-factor is the line of sight integration of the square dark matter density

$$\hat{I}(\psi) \equiv \int_0^{r_{\max}(\psi)} dr \left[\rho_h \left(\sqrt{r^2 - 2rR_{\odot}\cos\psi + R_{\odot}^2} | M_{\rm G}, 0 \right) \right]^2$$

$$r_{\max}(\psi) = R_{\odot}\cos\psi + \sqrt{R_{\mathrm{vir,G}}^2 - R_{\odot}^2\sin^2\psi}.$$

Dark Matter Annihilation

Dark matter annihilation cross-section:

 $\sigma v = a + b v^2$

a, b are constants

If S wave is suppressed then the cross-section is dominated by P wave \rightarrow b v²>>a

 $\rightarrow \sigma v$ is much smaller today compared to the freeze-out time

Thermal Relic Density: At freezeout, $\sigma v = 3x10^{-26} \text{ cm}^3/\text{s}$ High b/a lowers the cross-section at small v

In order to get large annihilation cross-section with large b/a, we need to go to non-thermal scenarios where enhanced annihilation cross-section may be needed to explain the dark matter content B.D., L. Leblond, K. Sinha, Phys.Rev. D80 (2009) 035014

γ-ray (Extragalactic)

At a position r, integrate $\sigma v(r)$ over the local velocity distribution to find the mean $[\sigma v]_h(r)=a + \lambda b \sigma^2_{uh}(r)$

a, b are fixed to satisfy $\sigma v=a +b v^2_{freeze out}$

Use halo velocity profile to find universal halo annihilation cross-section profile

$$\langle I_{\gamma} \rangle \left(E_{\gamma} \right) = [\sigma v]_0 \int \frac{\mathrm{d}z}{H(z)} W((1+z)E_{\gamma},z) \left\langle \rho^2 \left(1 + \frac{\lambda b}{a} \sigma_u^2 \right) \right\rangle(z)$$

Where:

$$= \left\langle \rho^2 \left(1 + \frac{\lambda b}{a} \sigma_u^2 \right) \right\rangle(z) = \int \mathrm{d}^3 \mathbf{r} \mathrm{d}M \frac{\mathrm{d}n}{\mathrm{d}M}(M, z) \rho_h^2(r|M, z) \left[1 + \frac{\lambda b}{a} \sigma_{uh}^2(r|M, z) \right]$$

S. Campbell, B.D., E. Komatsu, Phys. Rev. D 82, 095007 (2010)

γ-ray (Extragalactic)



γ-ray (Extragalactic)

It is possible to have scenarios where b/a is very large



S. Campbell, B.D., E. Komatsu, Phys. Rev. D 82, 095007 (2010)

Galactic and Extragalactic y–Ray



Focus Point: $tan\beta=10$; $M_{DM}=150$ GeV

Annihilation is primarily into W+W⁻ pair, σ v=1.9 10⁻²⁶ cm³/s

 Extragalactic signal is higher at lower energy due to cosmological redshifting

•The relative importance of galactic and extragalactic signal depends on different choices of parameters R. Allahverdi, S. Campbell, B.D. Phys. Rev. D 85, 035004 (2012)

Galactic and Extragalactic y–Ray



Galactic γ -ray intensity diverges as the line of sight approaches galactic center

However, it will be difficult to observe the dark matter annihilation around that region due to astrophysical contamination

Substructures can increase the extragalactic signal considerably

Galactic and Extragalactic y–Ray

We have not included the effects of substructures

Substructure can increase the annihilation rate by a factor of 100 or so depending on minimum halo mass size

Based on simulations, substructures would increase the galactic signal by a factor of few

Dark Matter Annihilation into v's

In MSSM, type models the neutrinos appear from W, b and τ final states from the DM decays

In MSSMx U(1)_{B-L} models: DM particle Sneutrino_R (\tilde{N}) annihilation can produce neutrino final states

• These models may contain small amount of photons $\tilde{N} > N^{c}$



IceCube and Fermi jointly can probe these models

Galactic and Extragalactic v



•The W decays produce a prompt component

•The prompt feature is washed out in the extragalactic spectrum due to redshifting

• Both galactic and extragalactic are contributing in this simplistic scenario R. Allahverdi, S. Campbell, B.D. Phys. Rev. D 85, 035004 (2012)

v All Sky vs. the Galactic Center



- The galactic signal is much stronger from the galactic center (assuming NFW cusp).
- Similar results for γ -rays.

• Subhalos and uncertainties in the minimum halo scale, halo concentrations, and distribution at the core/cusp need to be appropriately quantified.

γ Rays vs v





Gamma-ray intensity from annihilating 150 GeV dark matter for ψ >18°.

•A model producing W+W- is indistinguishable from a model annihilating to $b\overline{b}$. •All-sky $\nu + \overline{\nu}$ event rate per detector mass.

•Leptonic W or τ decays produce prompt neutrinos, which are absent from b decays.

•The neutrino signal breaks the γ -ray degeneracy between W and b producing annihilations.

Annihilation to Neutrinos



All-sky event rates for 150 GeV sneutrino dark matter that annihilates to two 135 GeV right-handed neutrinos (each flavor equally represented), each of which decays to a light neutrino and 120 GeV standard model Higgs particle

Annihilation into v's

- The secondary neutrinos produced from the Higgs decay result in a broad, soft spectrum, whereas the neutrinos produced directly from *N^c* decays produce a narrower peak at lower energies on the order of the mass difference between the *N^c* and the Higgs
- Due to the Higgs decays, there is also a gamma-ray component to the signal
 - If Higgs mass is small—negligible compared to the right sneutrino (DM) mass
- The spectrum of the produced light neutrinos is at the energy of the right sneutrino
- This simple scenario results in a prominent neutrino line feature

Annihilation to Neutrinos



R. Allahverdi, S. Campbell, B.D. Phys. Rev. D 85, 035004 (2012) 24

v final states and IceCube



R. Allahverdi, S. Bornhauser, B. D., K. Richardson-Mcdaniel, Phys.Rev. D80 (2009) 055026

Conclusion

Simultaneous observation of gamma-rays and neutrinos allows for more constrained conclusions about models

The signal contains both galactic and extragalactic component

Final state intensity depends on the annihilation cross-section, density profiles of the cores and halos substructures

Neutrinos may be more suited than the gamma rays for observing a signal from the galactic center